

The Grainger College of Engineering
Materials Research Laboratory

AMC2023

Advanced Materials
Characterization
Workshop



Optical Characterization Methods

Part I

Julio A. N. T. Soares

Materials Research Laboratory
University of Illinois at Urbana-Champaign

The Grainger College of Engineering
Materials Research Laboratory

UNIVERSITY OF ILLINOIS URBANA-CHAMPAIGN

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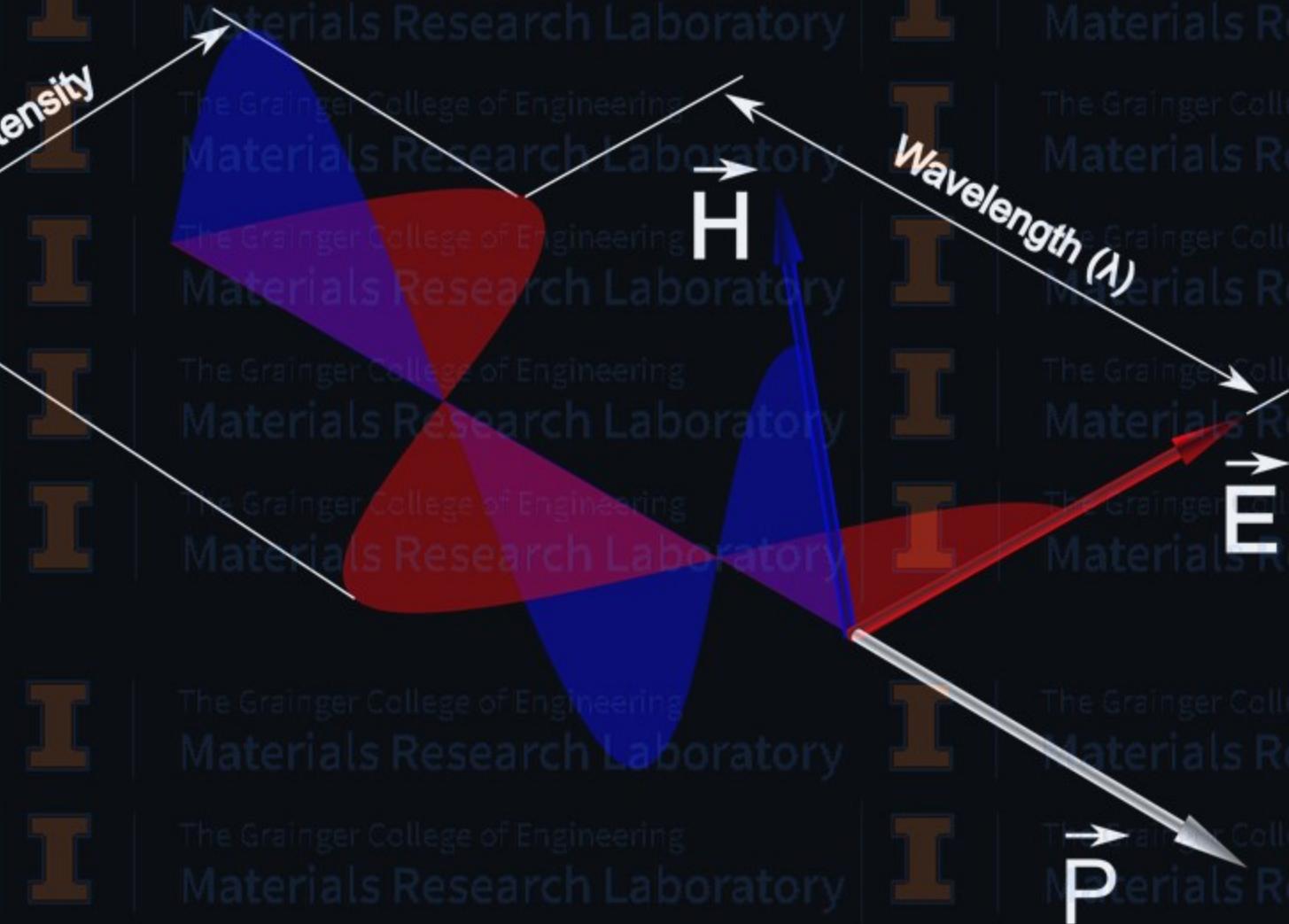
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Optical Characterization

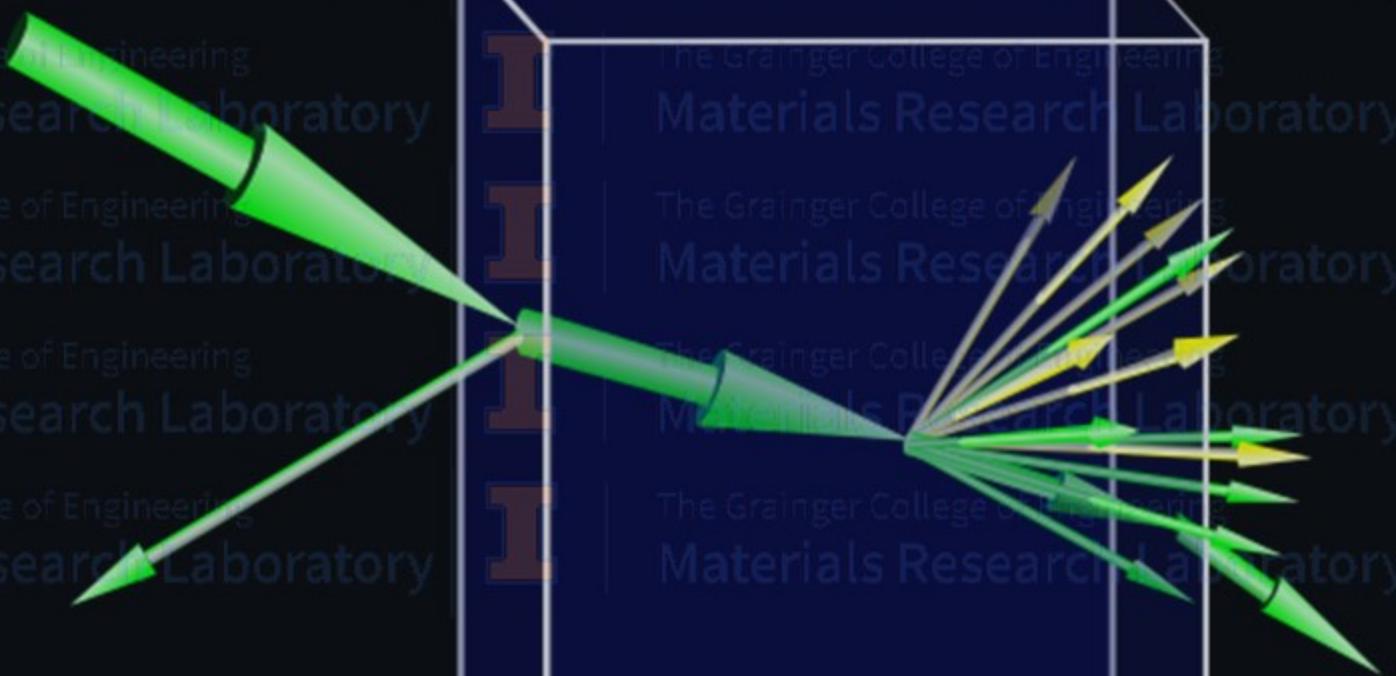
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Light properties
Light-matter interaction
Instrumentation and methods
Application examples
Strengths and limitations
Complementary techniques

- Direction of propagation
- Electric field direction or polarization
- Photon energy or wavelength
- Intensity



- Transmission
- Reflection
- Absorption
- Emission
- Scattering
- Refraction

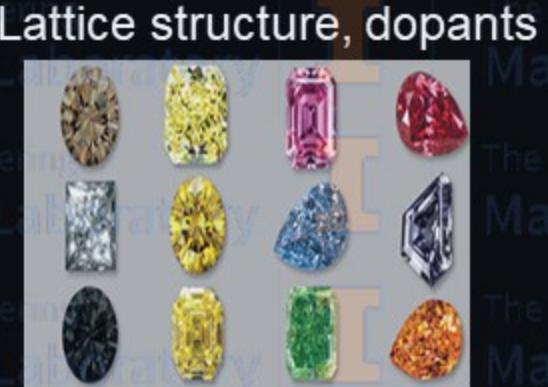


- Non-linear effects
 - SFG
 - SHG
 - DFG
 - Multi-photon absorption

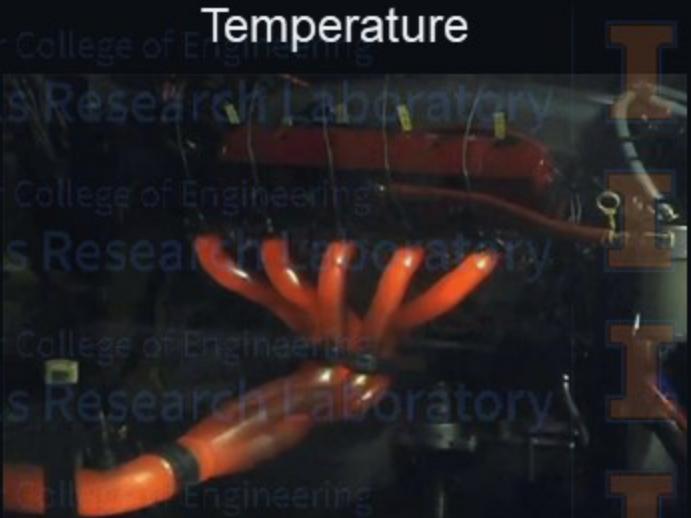
Light interactions with matter



Size



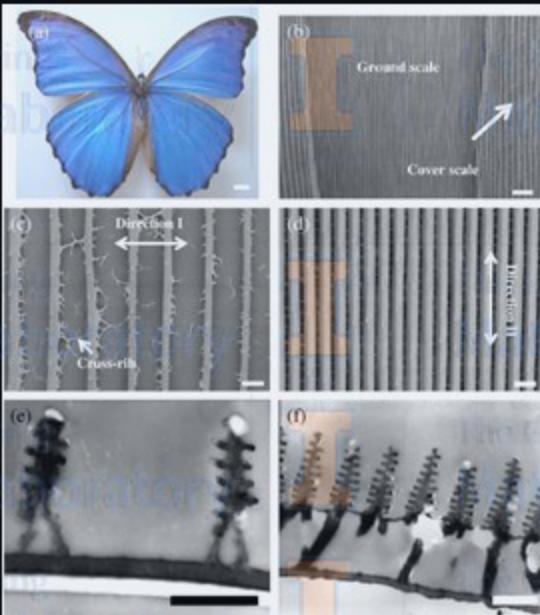
Lattice structure, dopants



Temperature



Thickness



Microstructure



Concentration

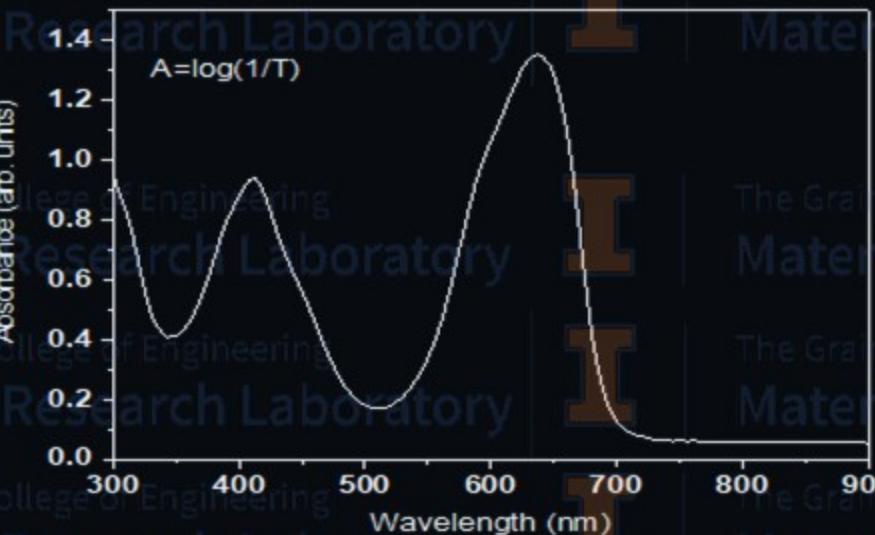
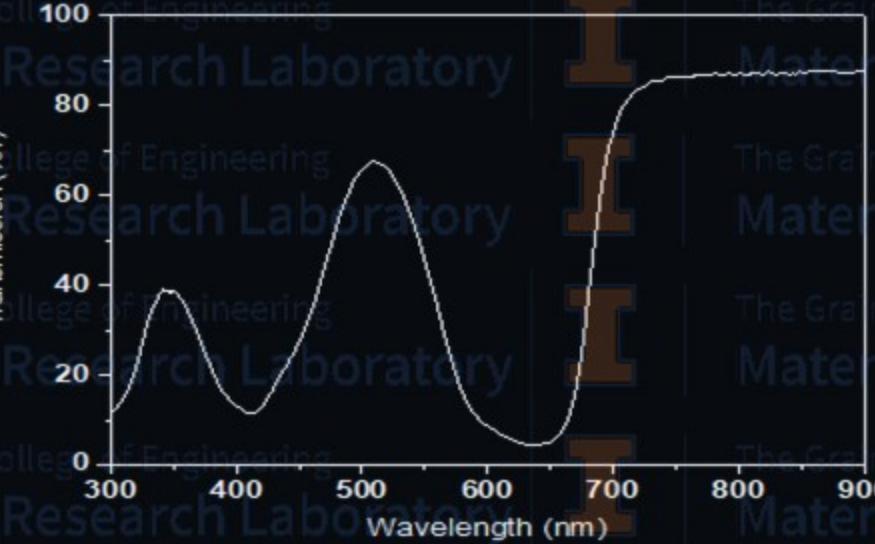
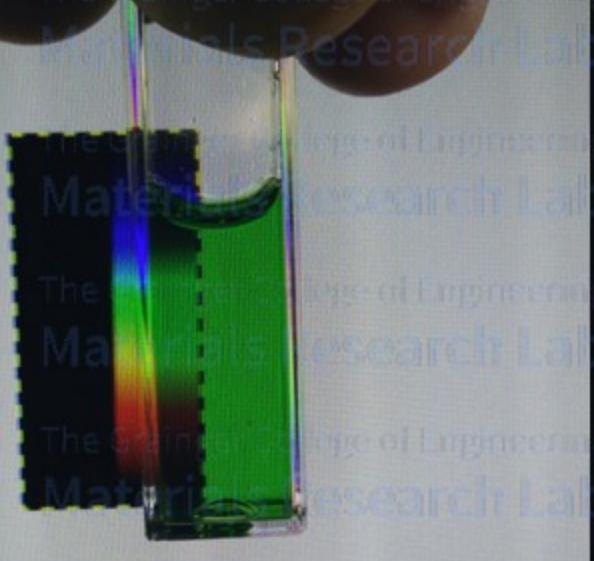


Composition



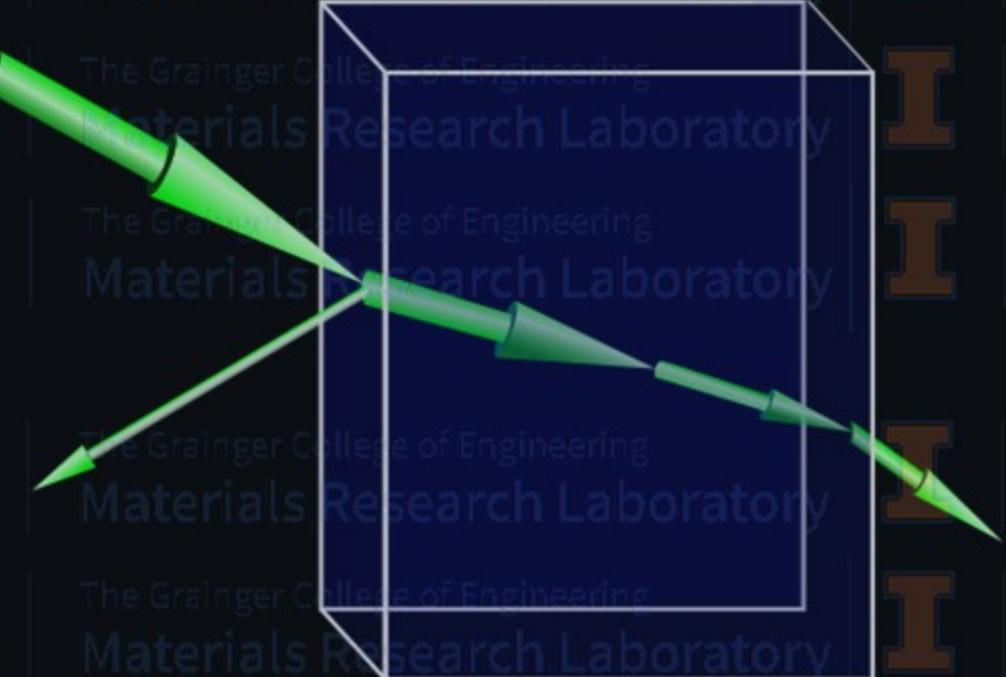
photon energy





I What is measured:

The transmitted and reflected light intensity as a function of the incident photon energy, which depends on the material's electronic, atomic, chemical and morphological structure.



I UV-Vis-NIR

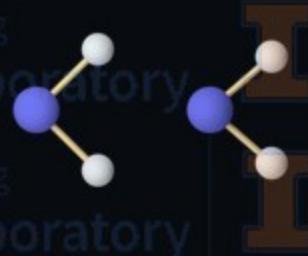
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A red sphere with a yellow starburst pattern in the center, representing a light source or detector.

I IR

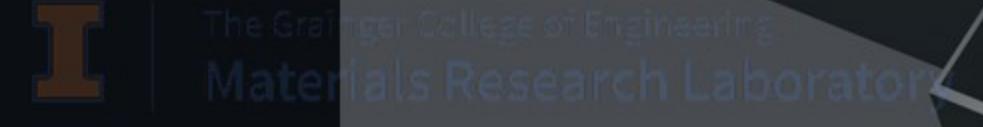
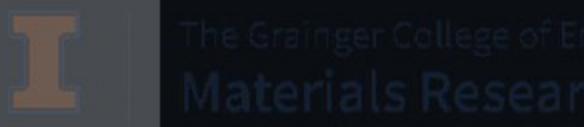
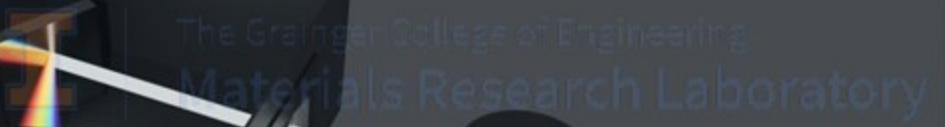
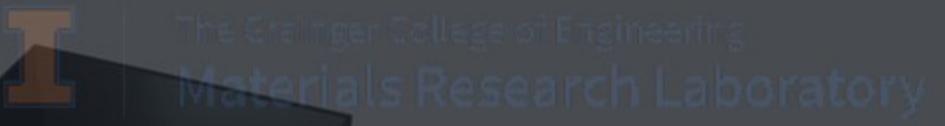
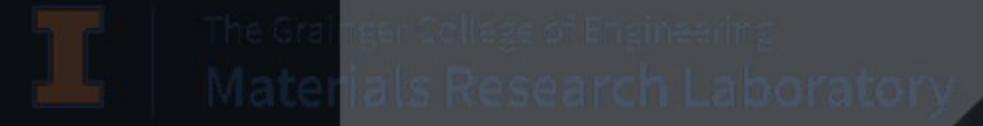
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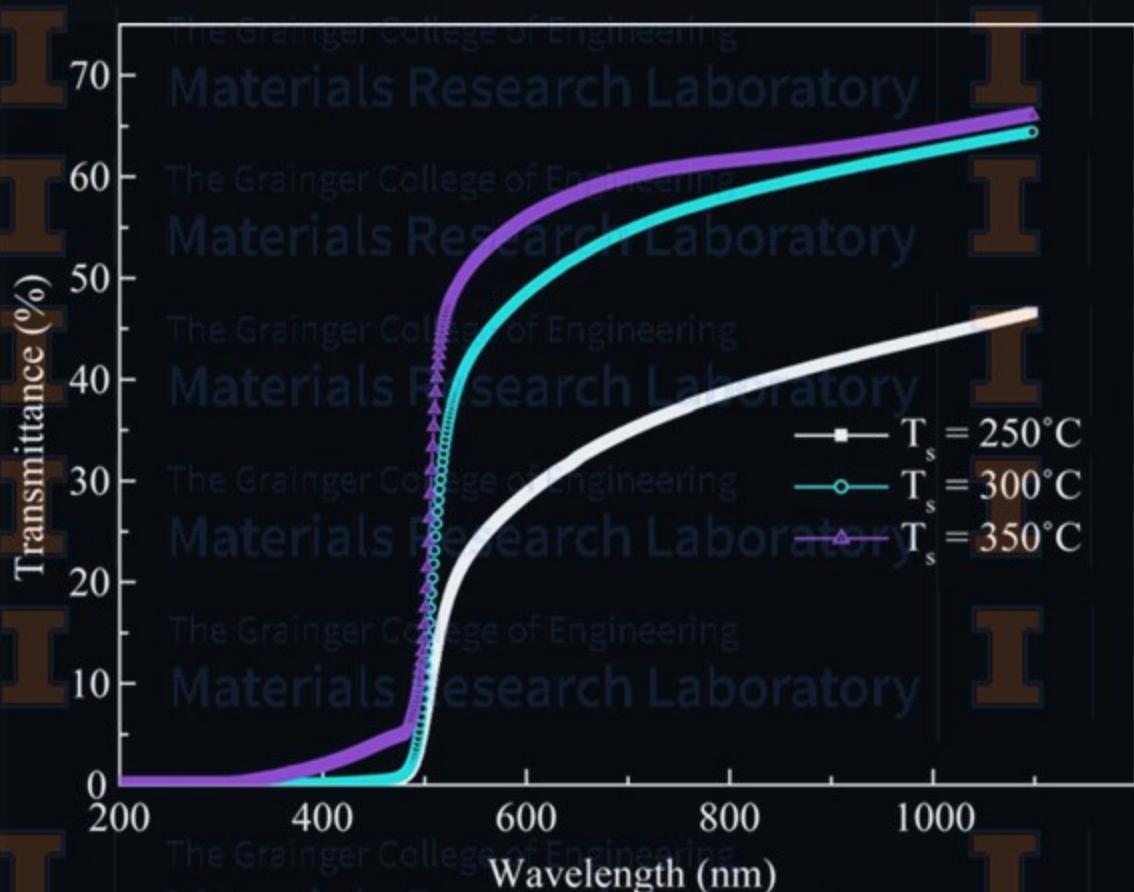


Two molecular models. The first shows a single molecule with two blue spheres connected by a yellow line. The second shows a dimer molecule with two blue spheres connected by two yellow lines, each attached to a white sphere.

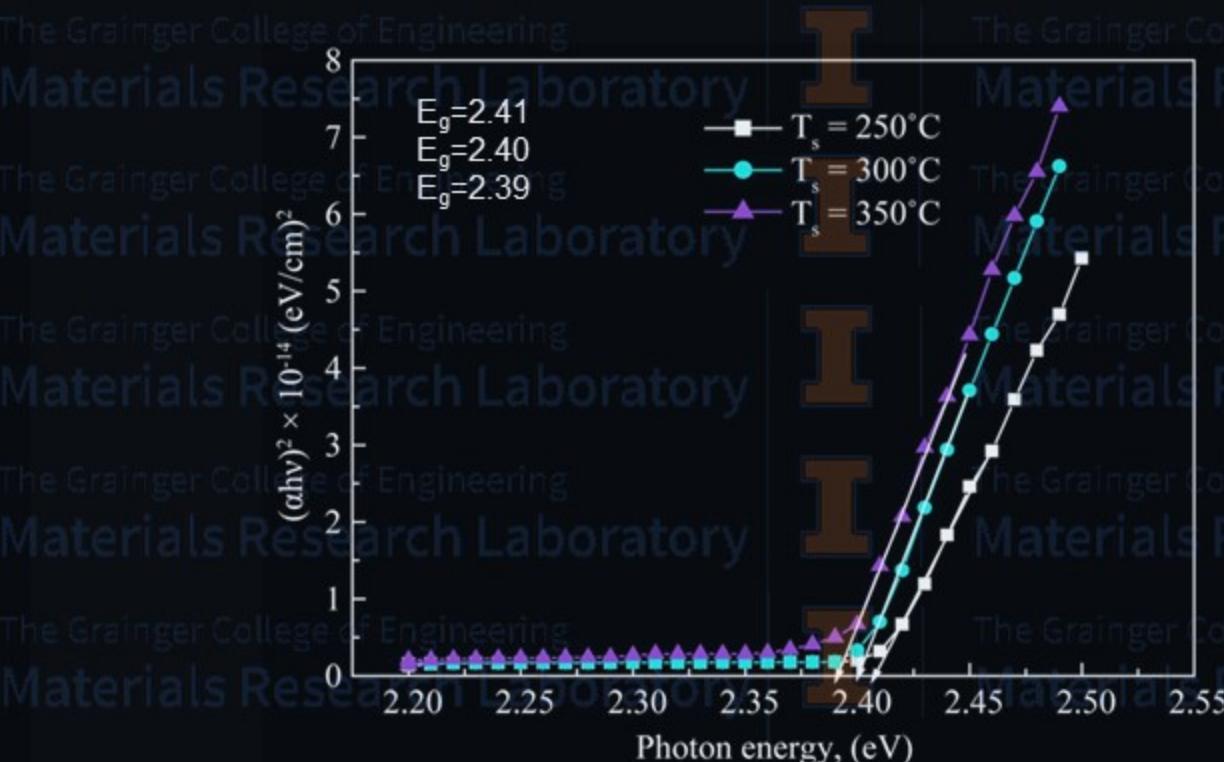


Spectrophotometry (UV-VIS-NIR)





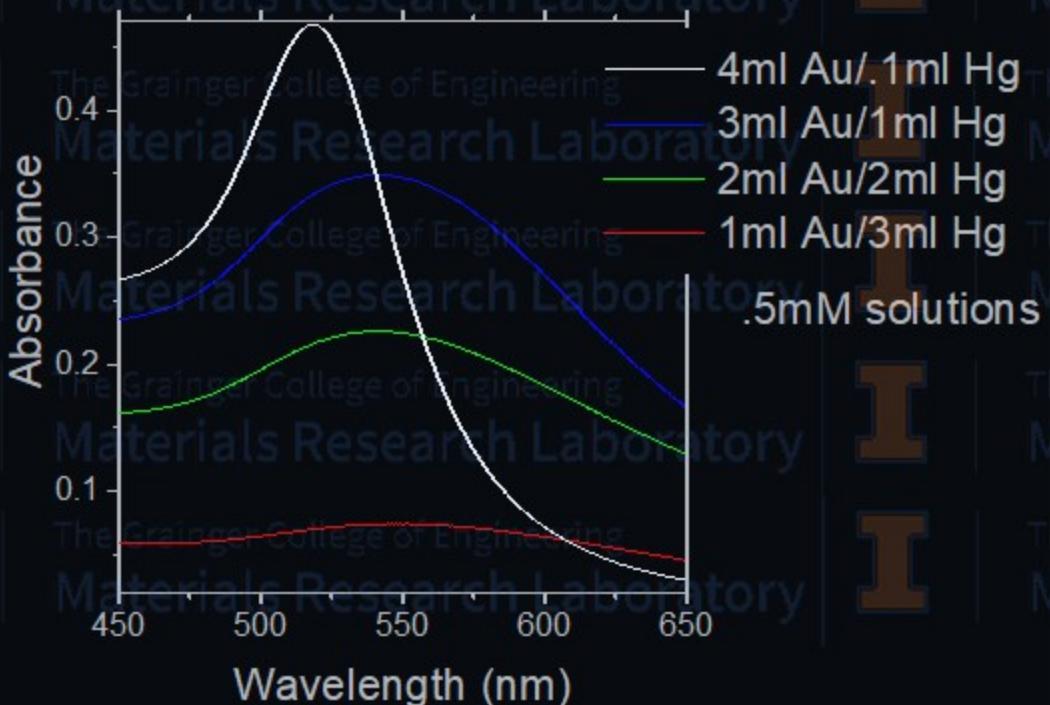
Optical band gap determination of CdS thin films as a function of growth substrate temperatures



Tauc's relation:

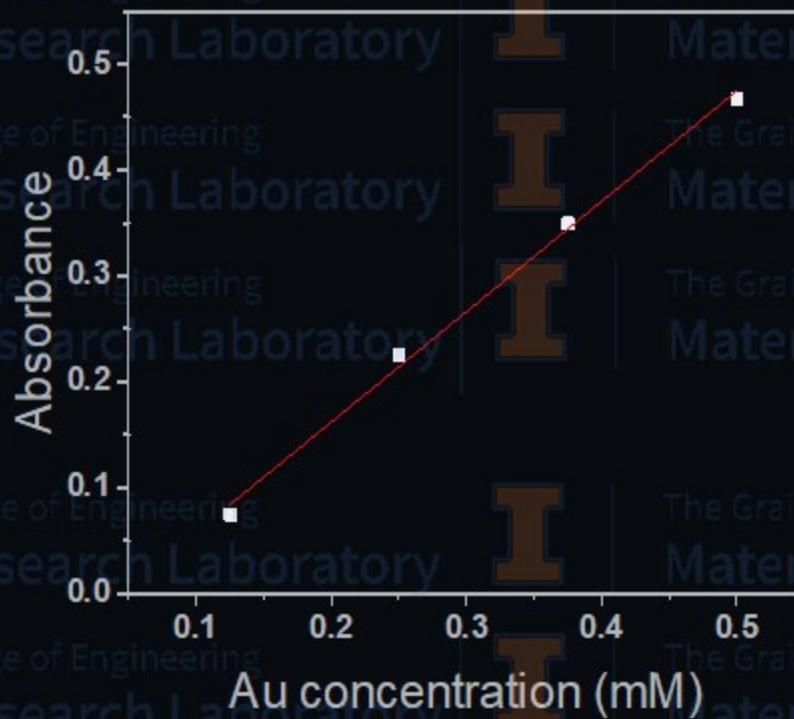
$$ahv = A(hv - E_g)^m$$

m = 0.5 for direct and 2 for indirect allowed transitions.



Beer-Lambert Law

$$Abs = K \ell c = a \ell$$
$$Abs = \log (1/T)$$



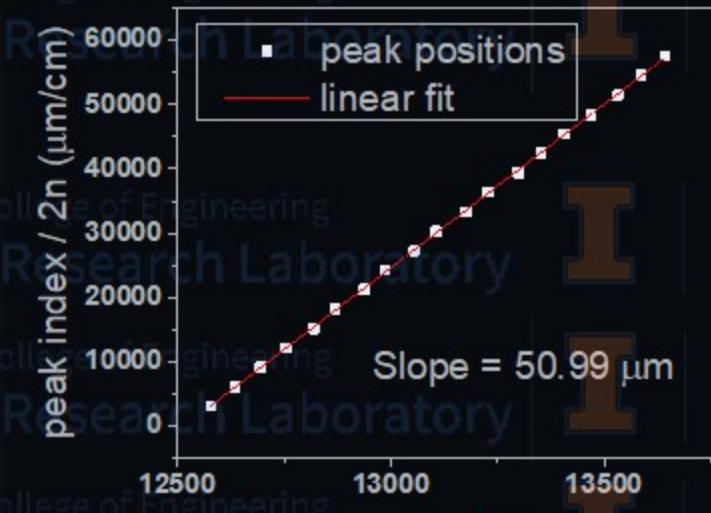
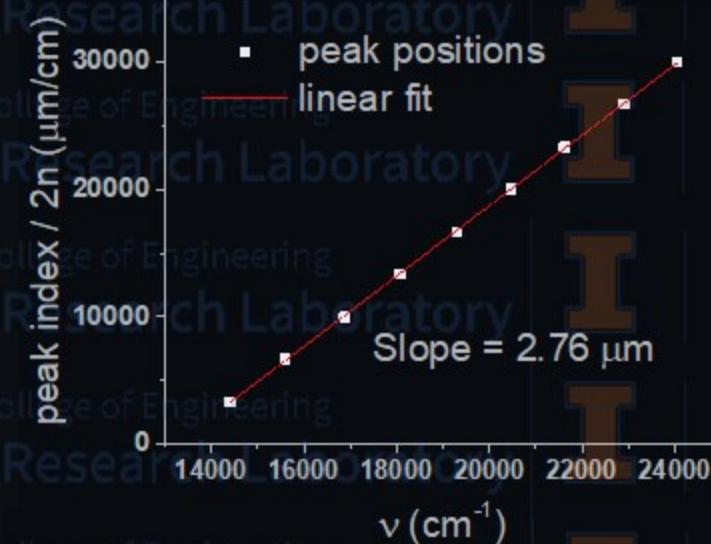
Using absorbance to determine Au/Hg concentration in water solutions

Spectrophotometry (UV-VIS-NIR)



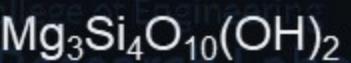
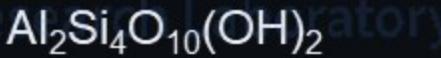
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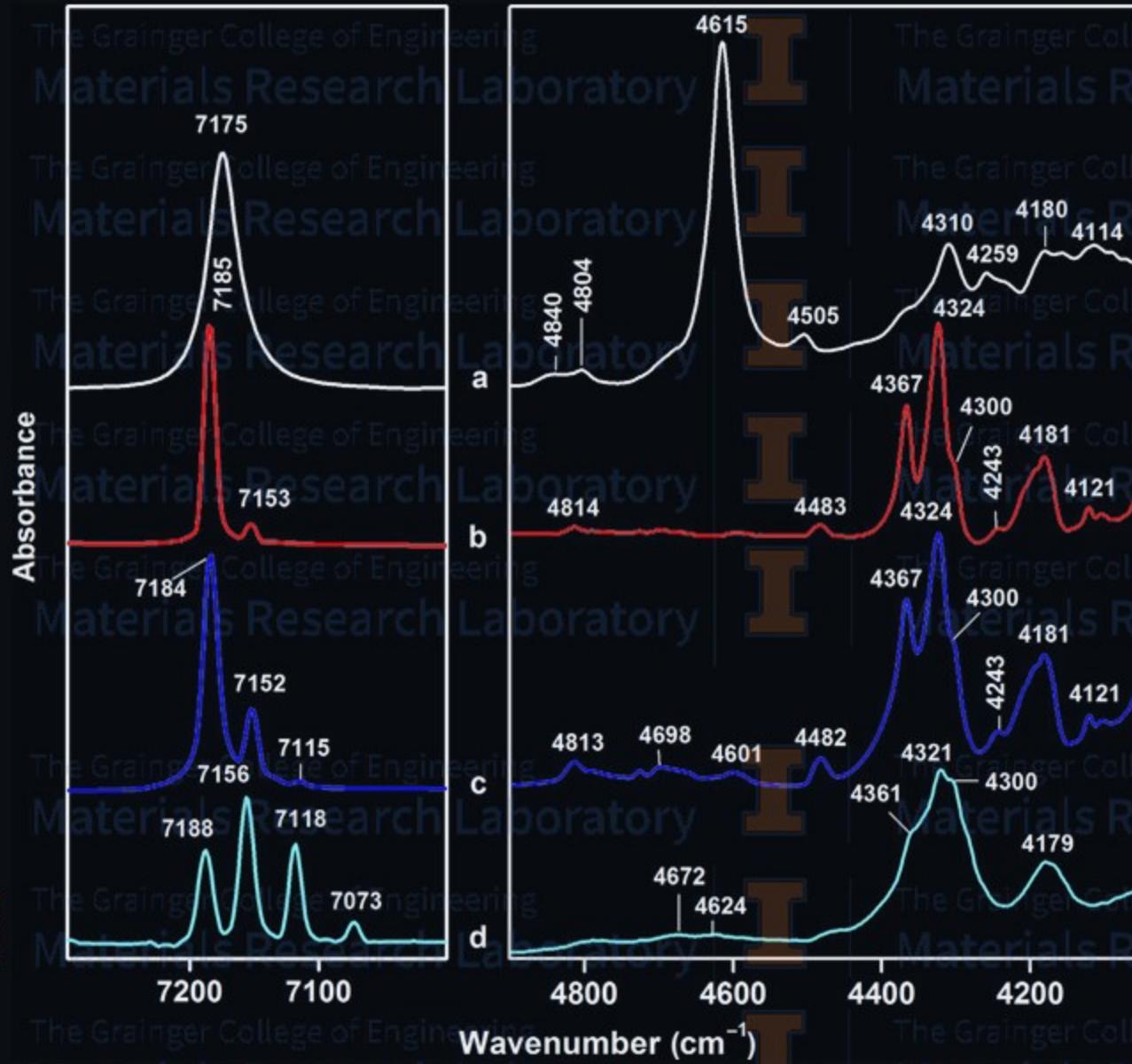
Using transmission interference fringes to determine thickness

Spectrophotometry (UV-VIS-NIR)



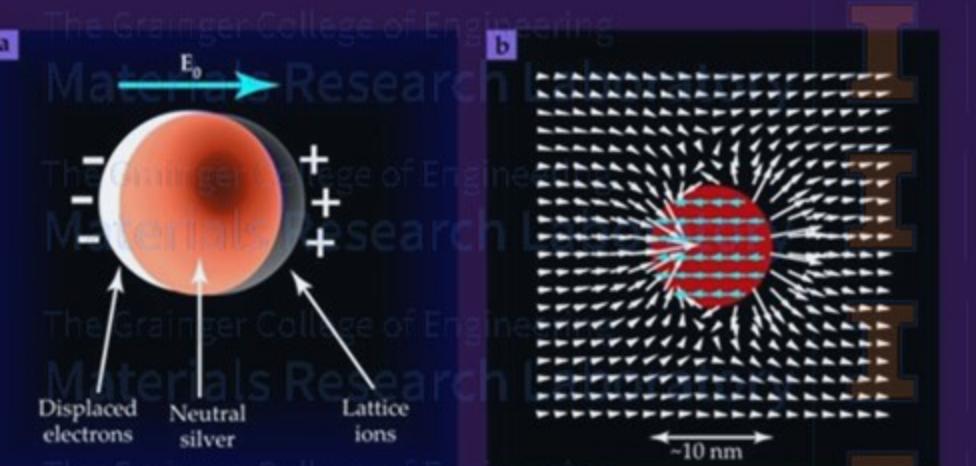
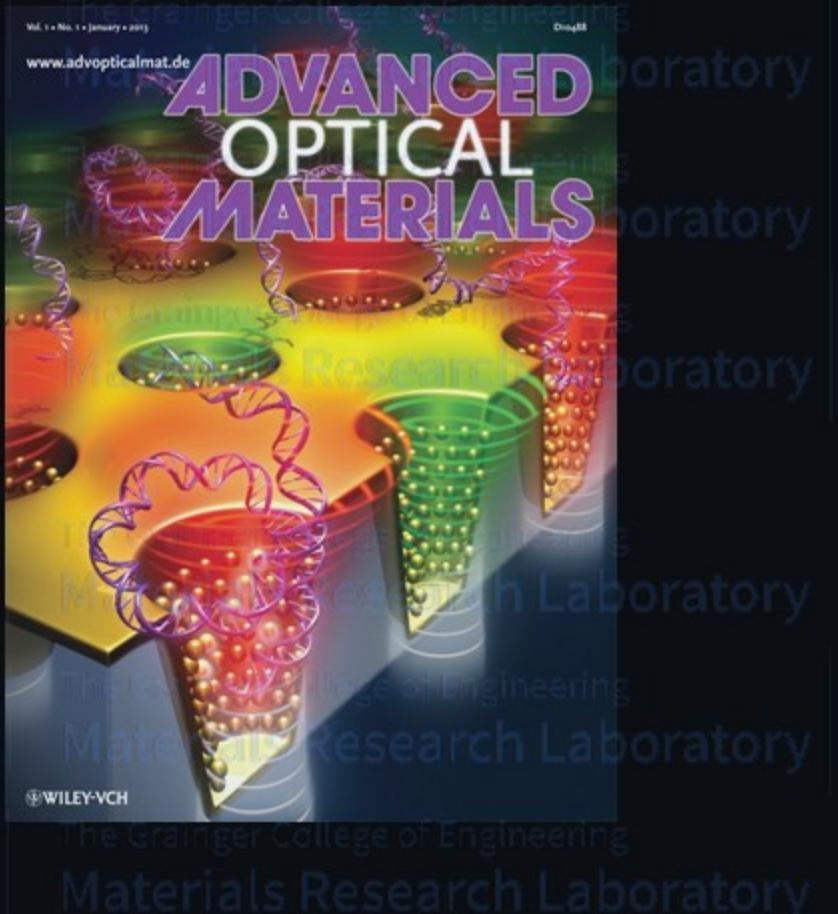
Images from Wikipedia and eurotalc.com

Spectra from *Developments in Clay Science*, Vol. 8, Ch. 5



Excitations in materials

- Plasmons



Phys. Today, 64, 39 (2011)

Dielectric
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Metal

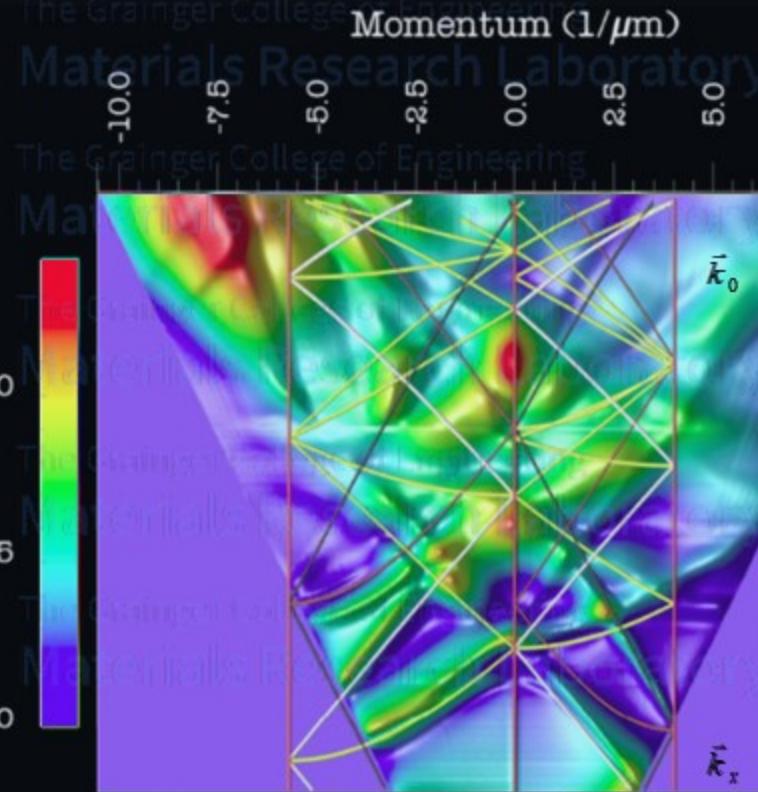
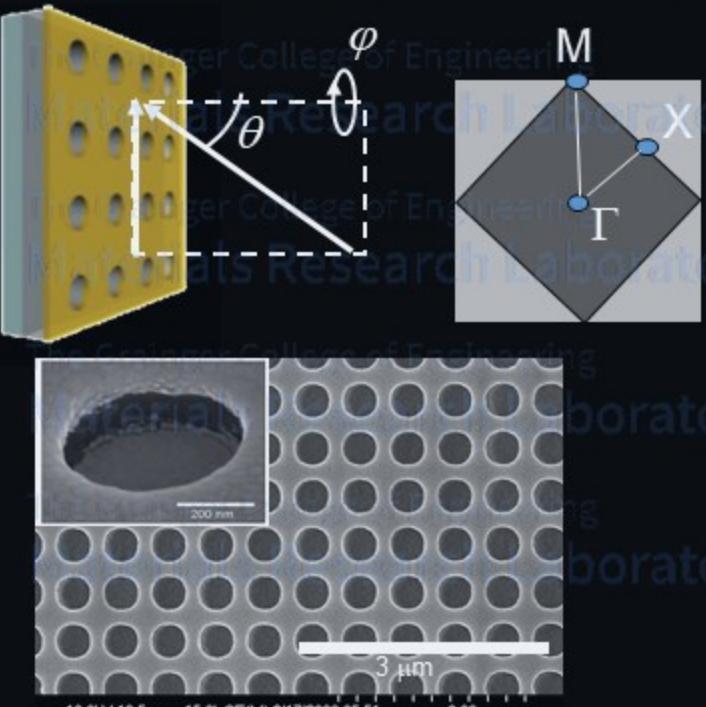
www.juluribk.com

Plasmons are quanta of collective motion of charge-carriers in a gas with respect of an oppositely charged background. They play a significant role on transmission and reflection of light.

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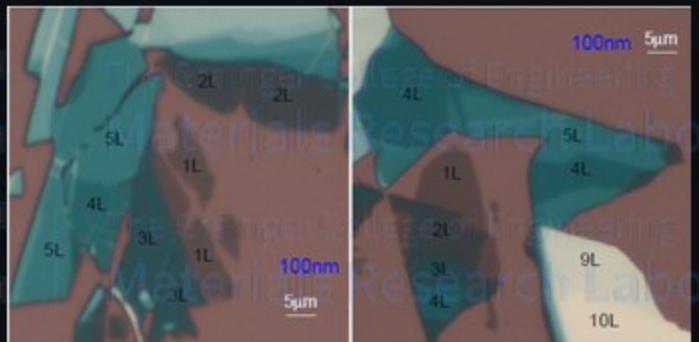
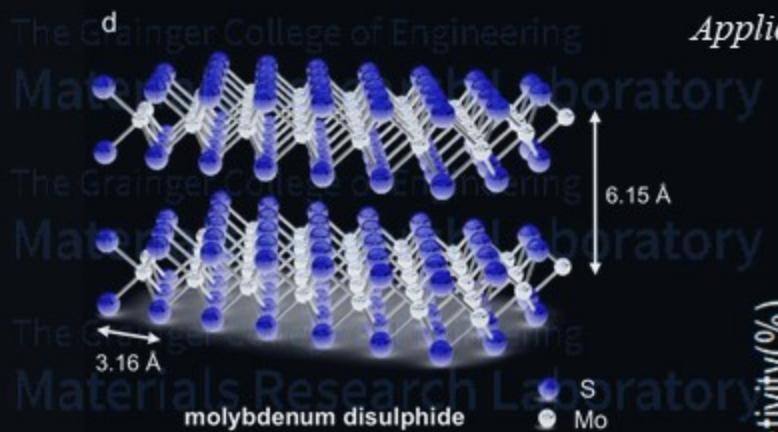
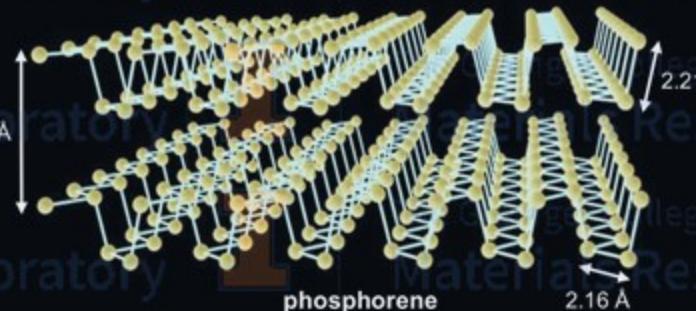
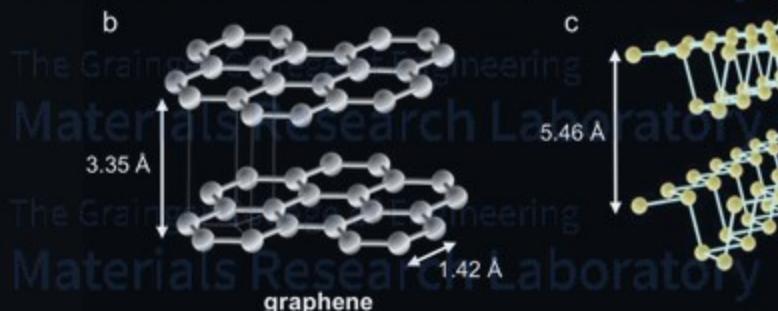
Spectrophotometry (UV-VIS-NIR)



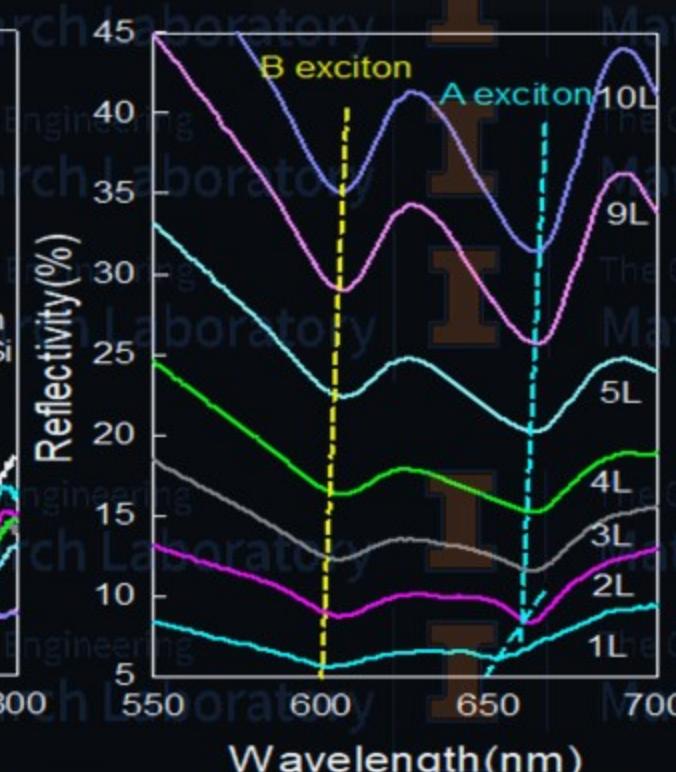
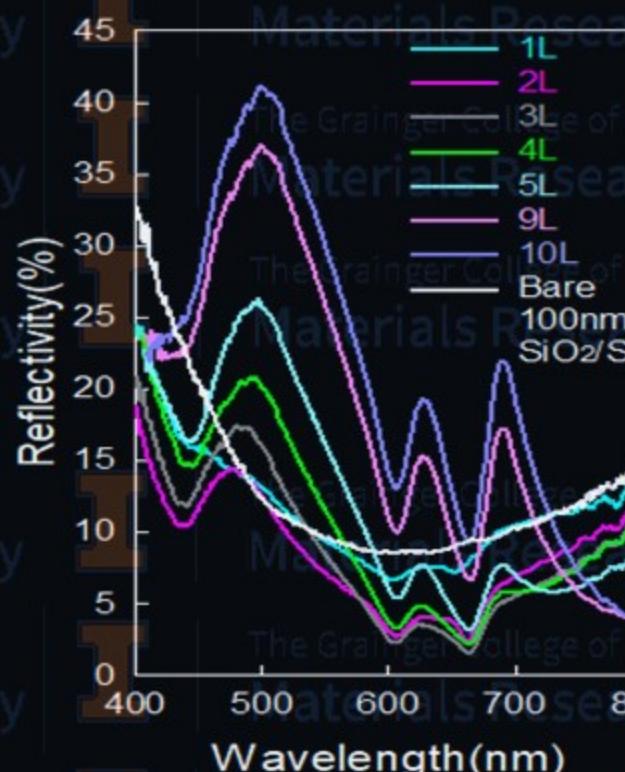
Optics Express 13, 5669 (2005)

Plasmonic crystal Brillouin zone from the transmission spectra measured for many different angles of incidence.

Spectrophotometry (UV-VIS-NIR)



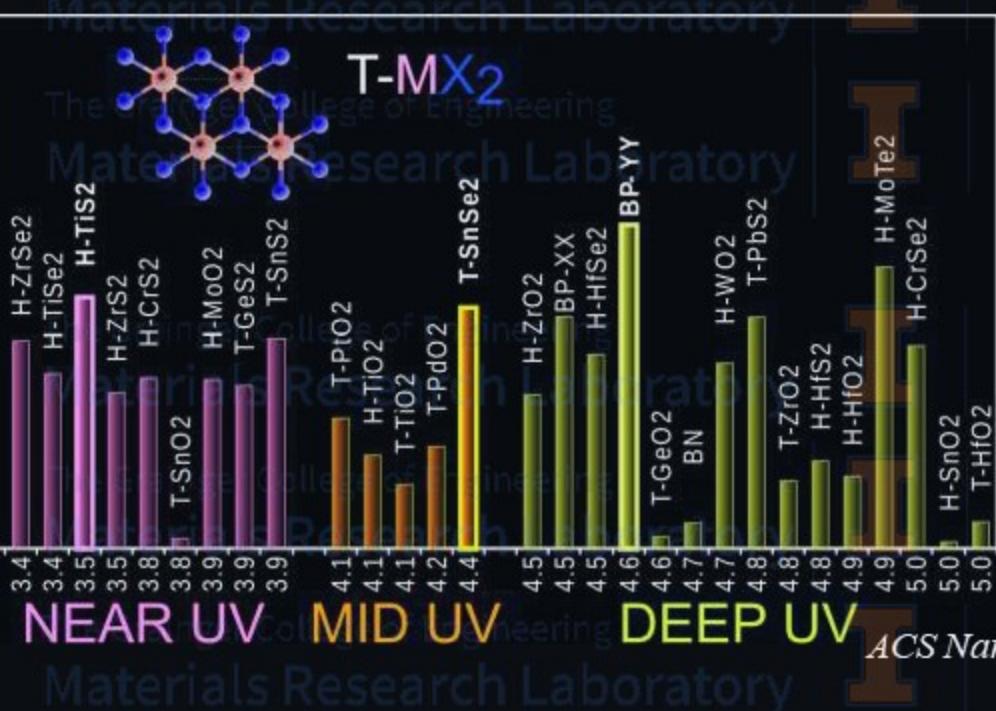
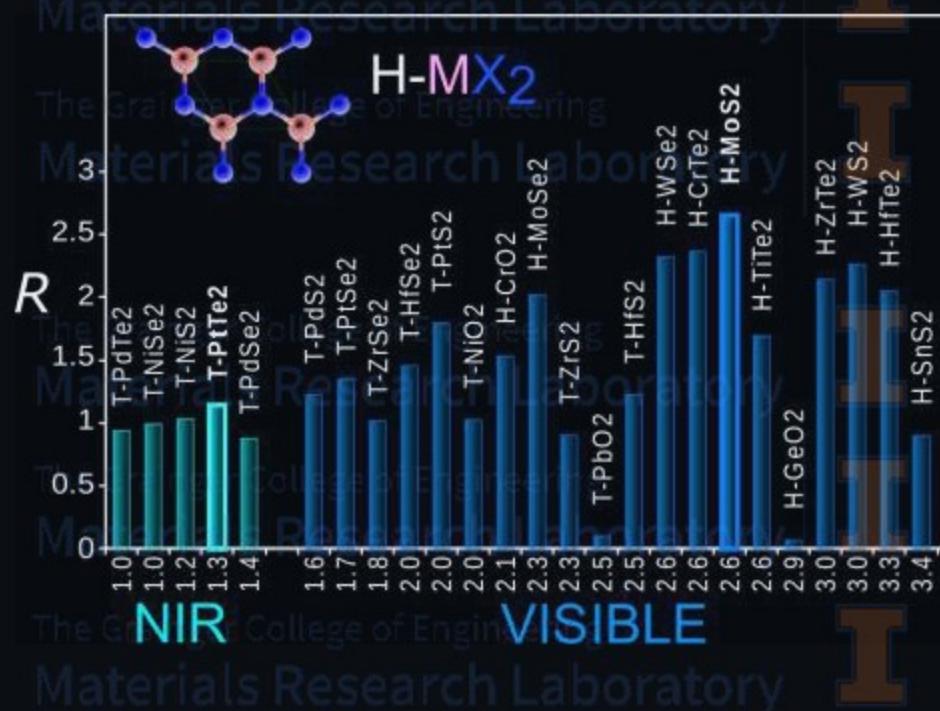
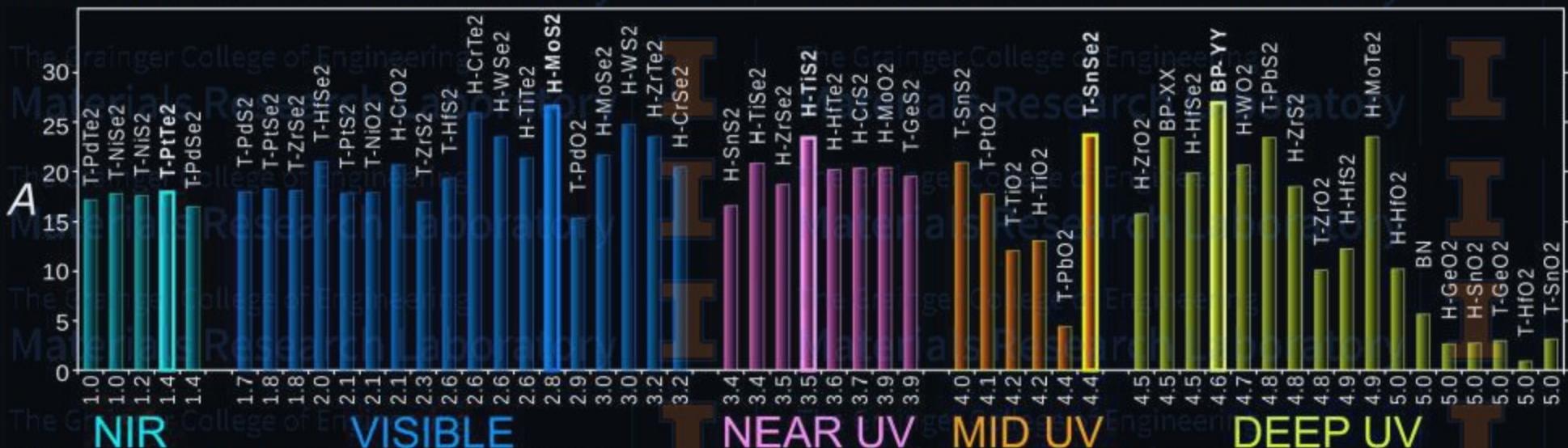
Applied Materials Today 8, 68 (2017)

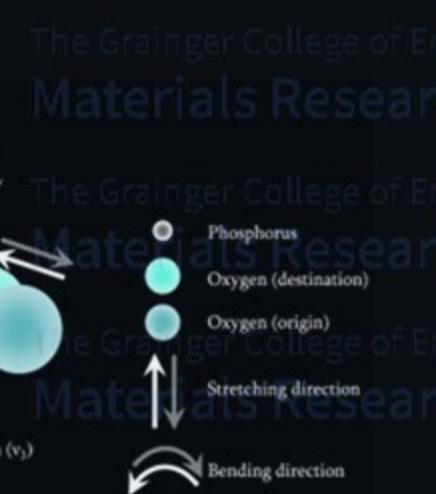
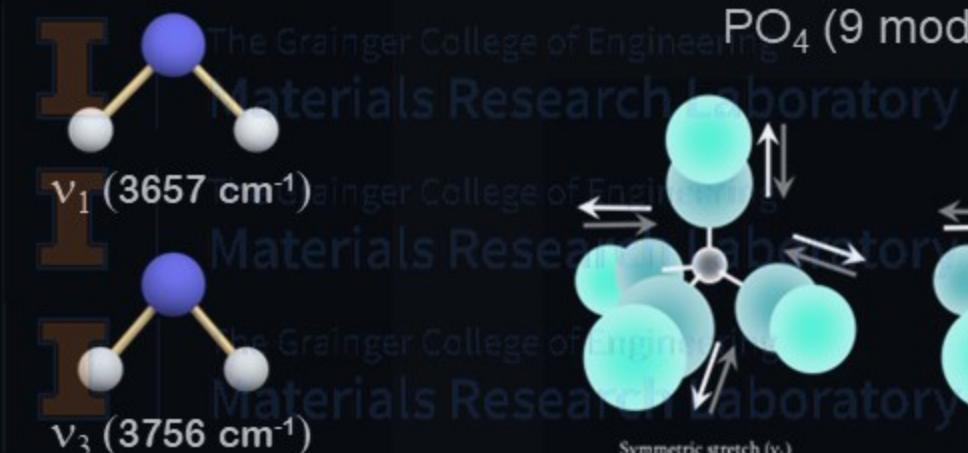
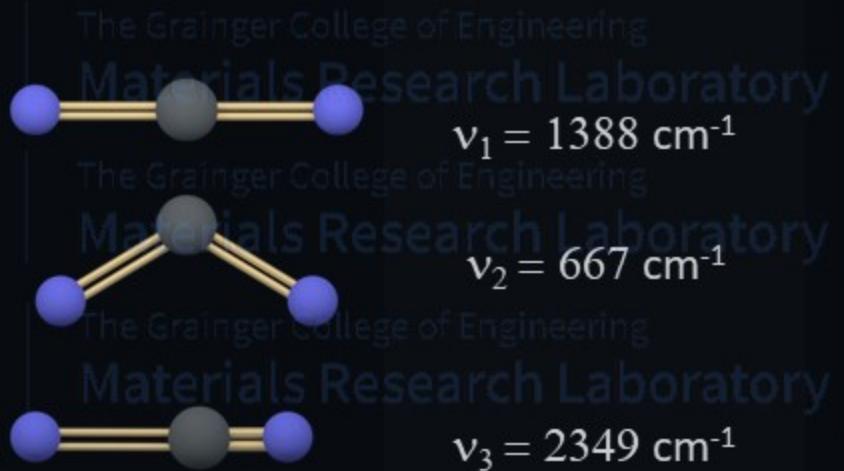


Optical Materials Express, 332858 (2018)

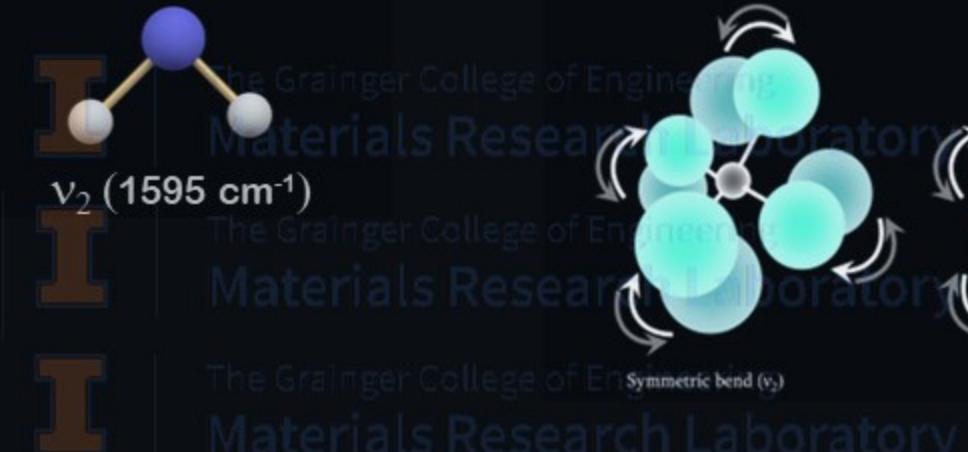


Spectrophotometry (UV-VIS-NIR)

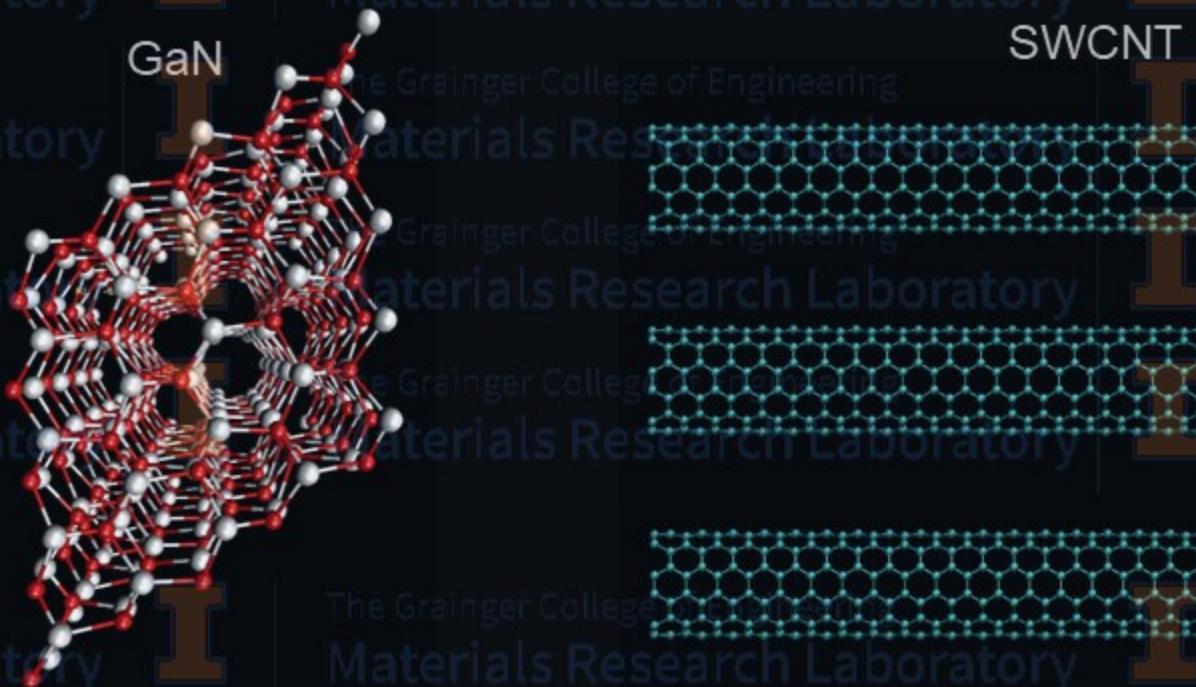




Number of modes:
3N-6 for non-linear molecules
3N-5 for linear molecules



Normal vibrational modes in solids:



<http://www.phonon.fc.pl>

Fourier Transform IR spectroscopy (FTIR)

IR active vibrations

The intensity of a vibrational absorption depends on the change of the transition dipole momentum caused by that vibration, so a vibration mode j will be “IR active” only when the vibration causes a change in the dipole momentum of the molecule, i.e. $\Delta\mu \neq 0$

IR active



$\Delta\mu \neq 0$

IR active



$\Delta\mu = 0$

IR active

IR active



$\Delta\mu \neq 0$



$\Delta\mu = 0$



$\Delta\mu \neq 0$

IR active



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$\Delta\mu \neq 0$

IR active



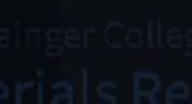
$\Delta\mu = 0$

IR active



$\Delta\mu \neq 0$

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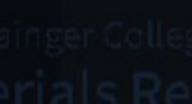
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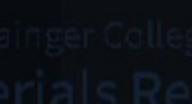
$\Delta\mu \neq 0$

IR active



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IR active



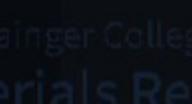
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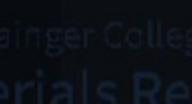
$\Delta\mu = 0$

IR active



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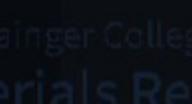
$\Delta\mu = 0$

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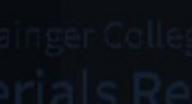
$\Delta\mu = 0$

IR active



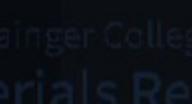
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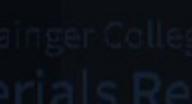
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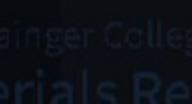
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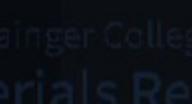
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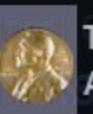


$\Delta\mu \neq 0$

IR active



Fourier Transform IR spectroscopy (FTIR)



The Nobel Prize in Physics 1907

Albert A. Michelson

"for his optical precision instruments and the
spectroscopic and metrological investigations
carried out with their aid"

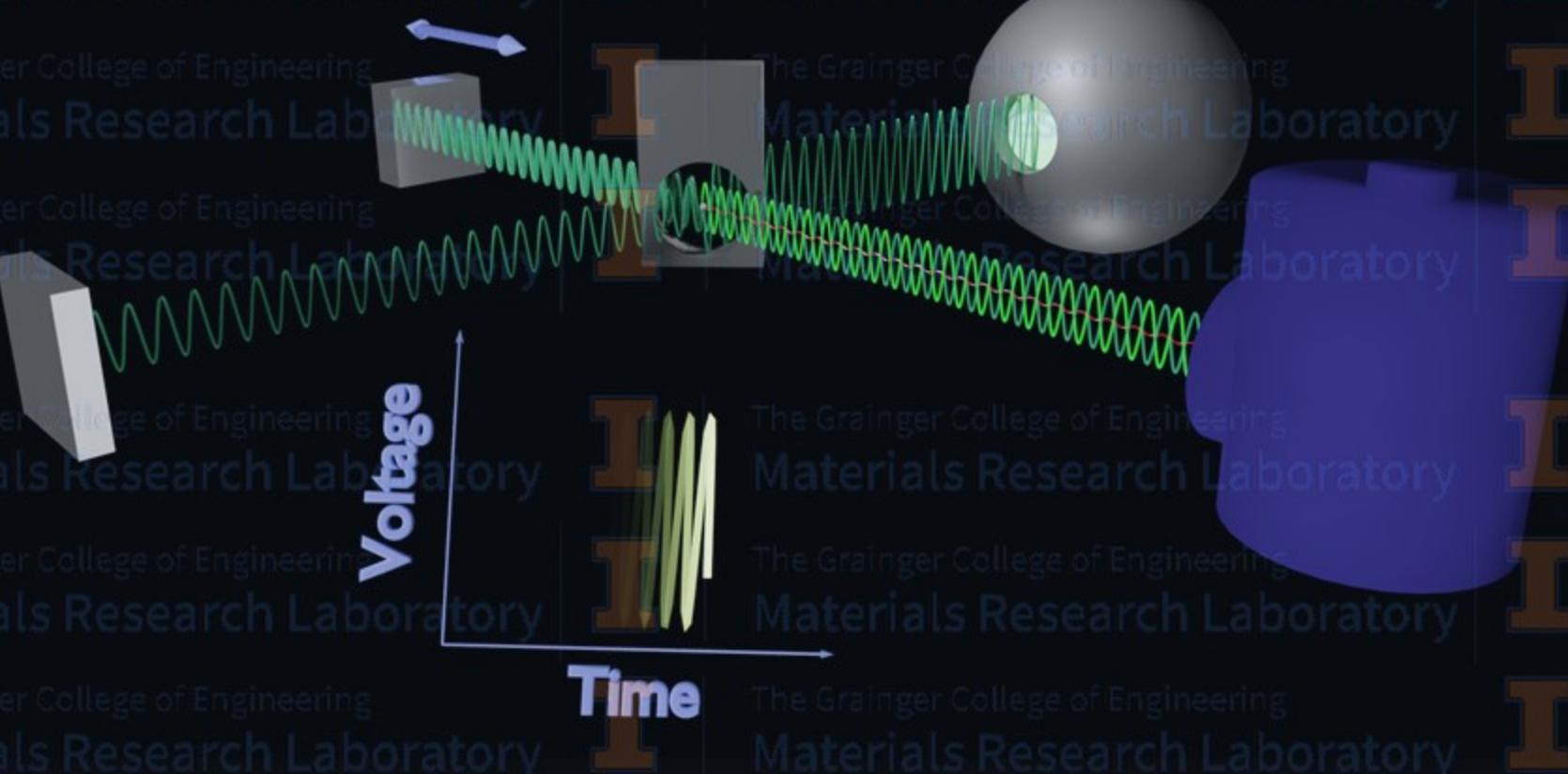
The Nobel Foundation

The Grainger College of Engineering Instrumentation:

The FTIR uses a Michelson interferometer with a moving mirror,
in place of a diffraction grating or prism.

$$\Delta L = n\lambda \Rightarrow \text{constructive interference}$$

$$\Delta L = (n+1/2) \lambda \Rightarrow \text{destructive interference}$$



Fourier Transform IR spectroscopy (FTIR)



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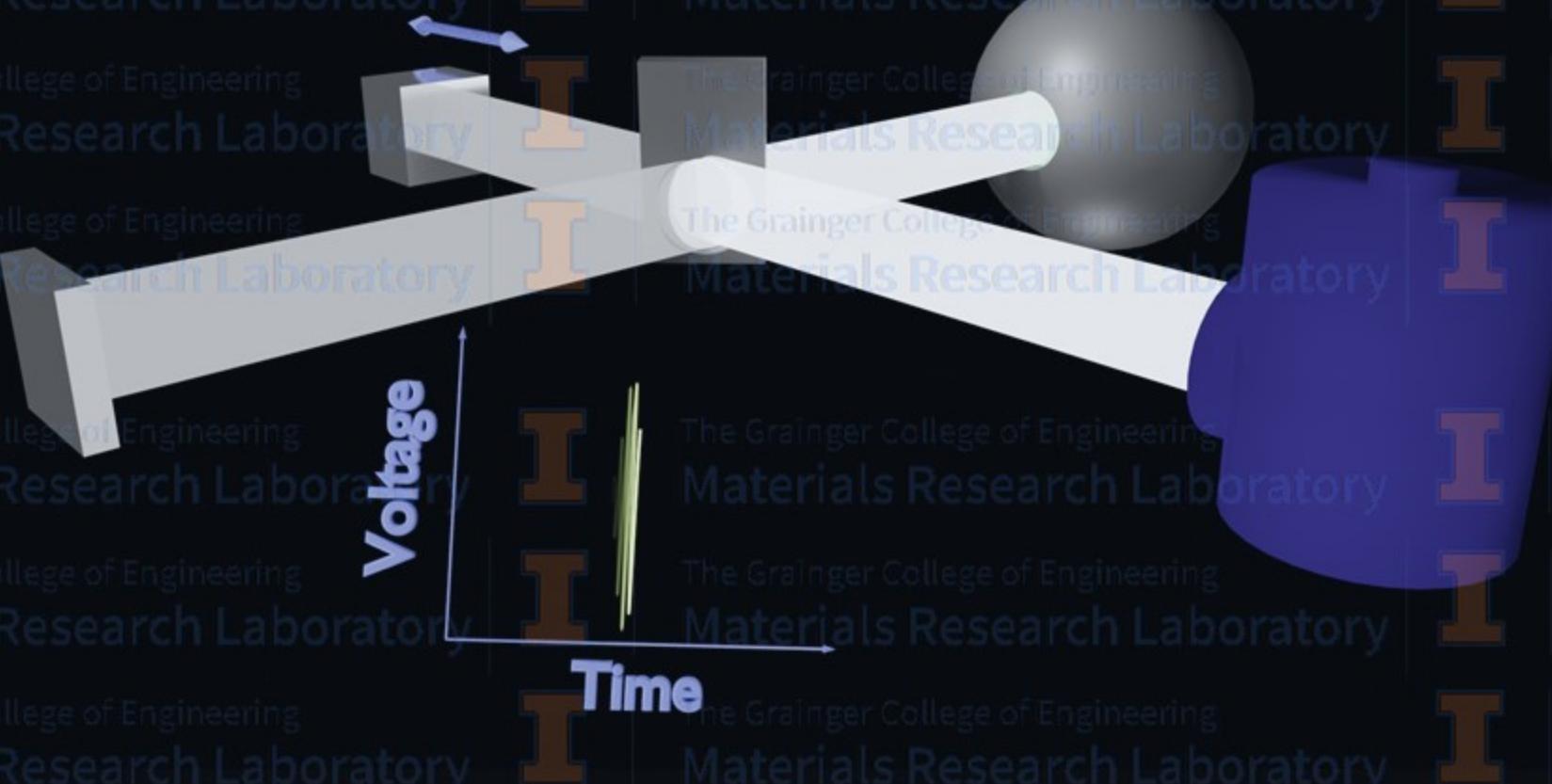
The Nobel Foundation

The Grainger College of Engineering Instrumentation:

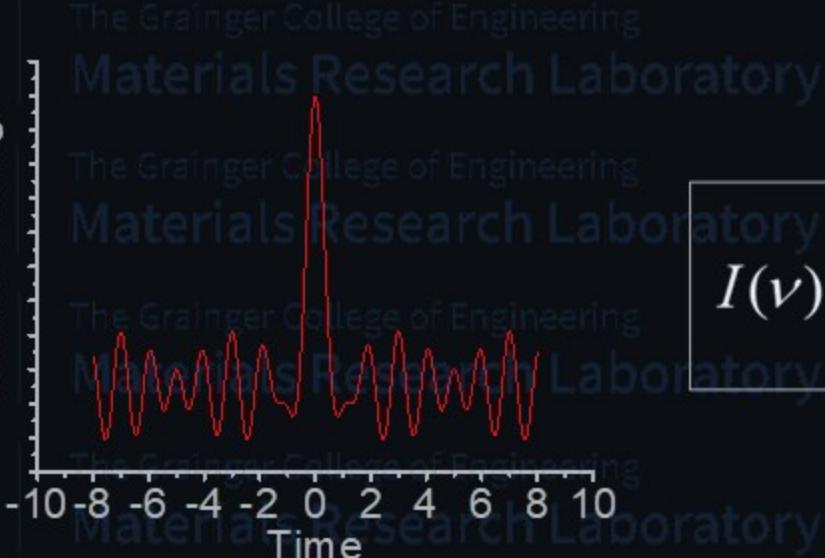
The FTIR uses a Michelson interferometer with a moving mirror,
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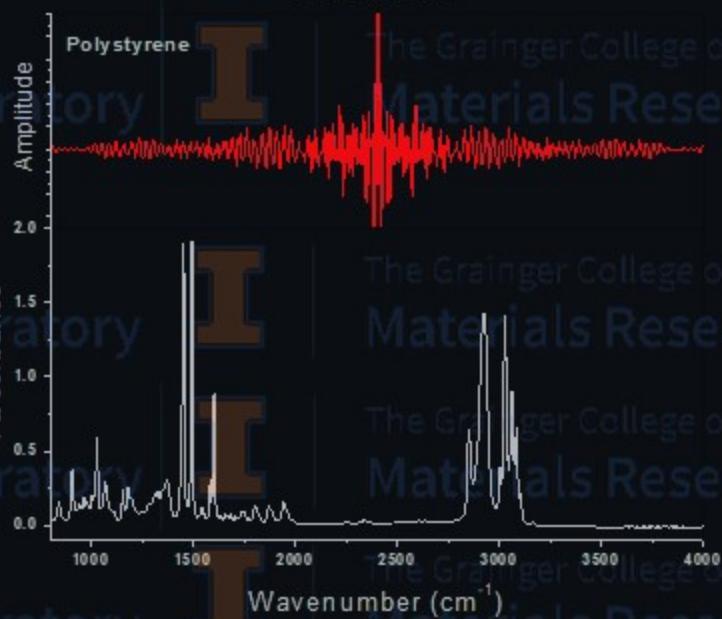
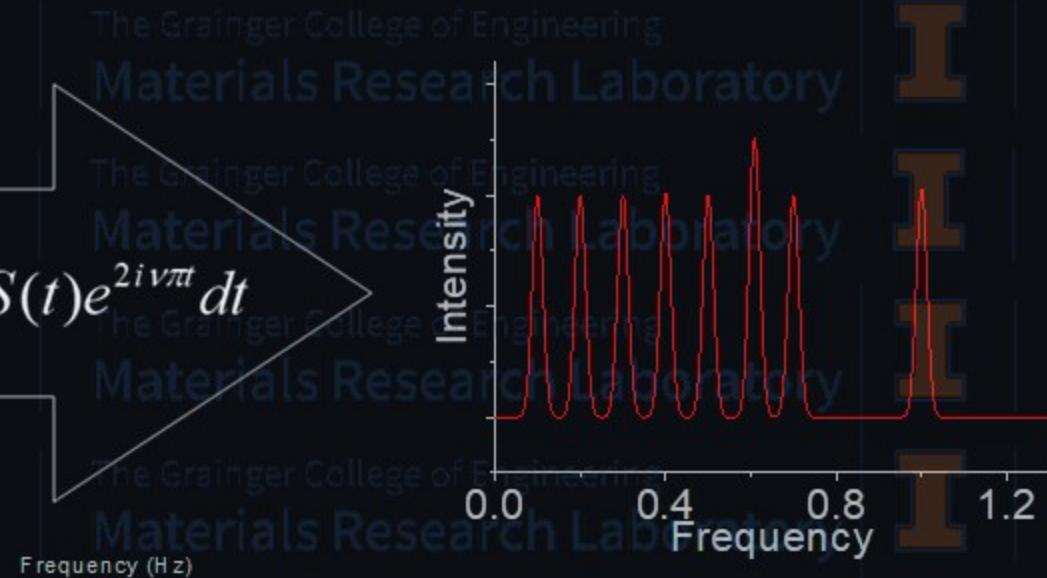
$$\Delta L = (n+1/2) \lambda \Rightarrow \text{destructive interference}$$



Fourier Transform IR spectroscopy (FTIR)



$$I(v) = \int S(t) e^{2i v \pi t} dt$$

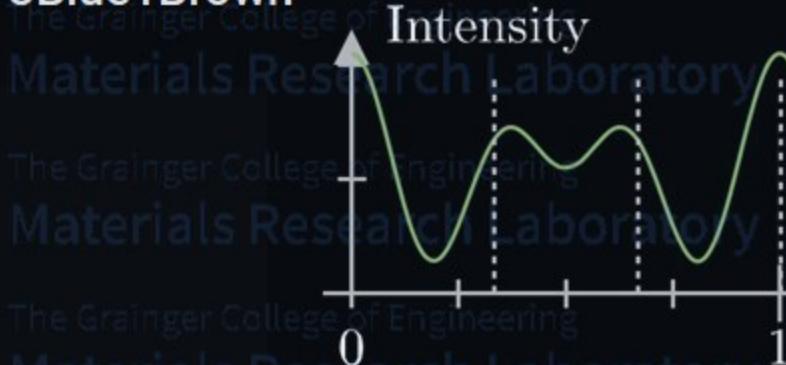


But what is the Fourier Transform? A visual introduction.

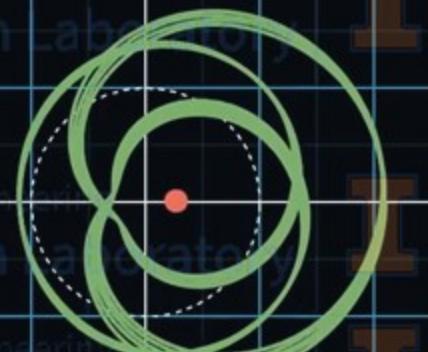


3Blue1Brown

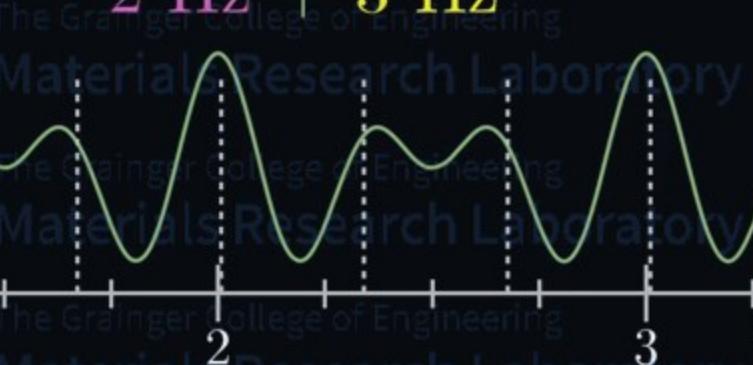
Intensity



2.99 cycles/second

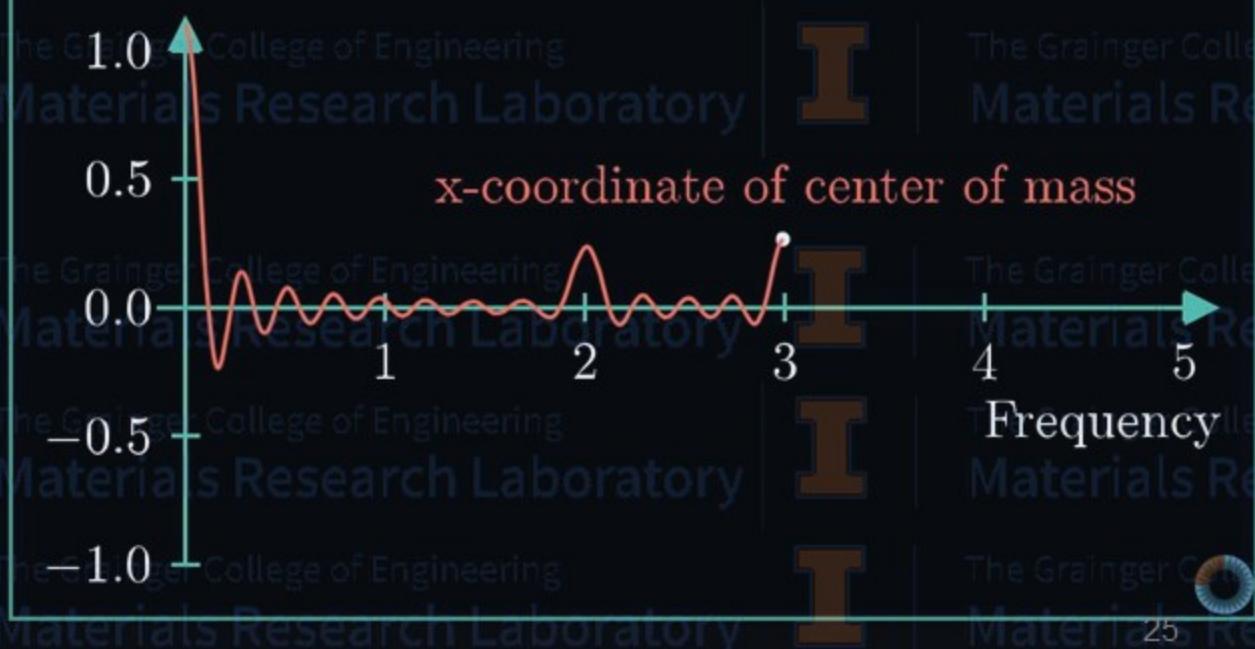


2 Hz + 3 Hz

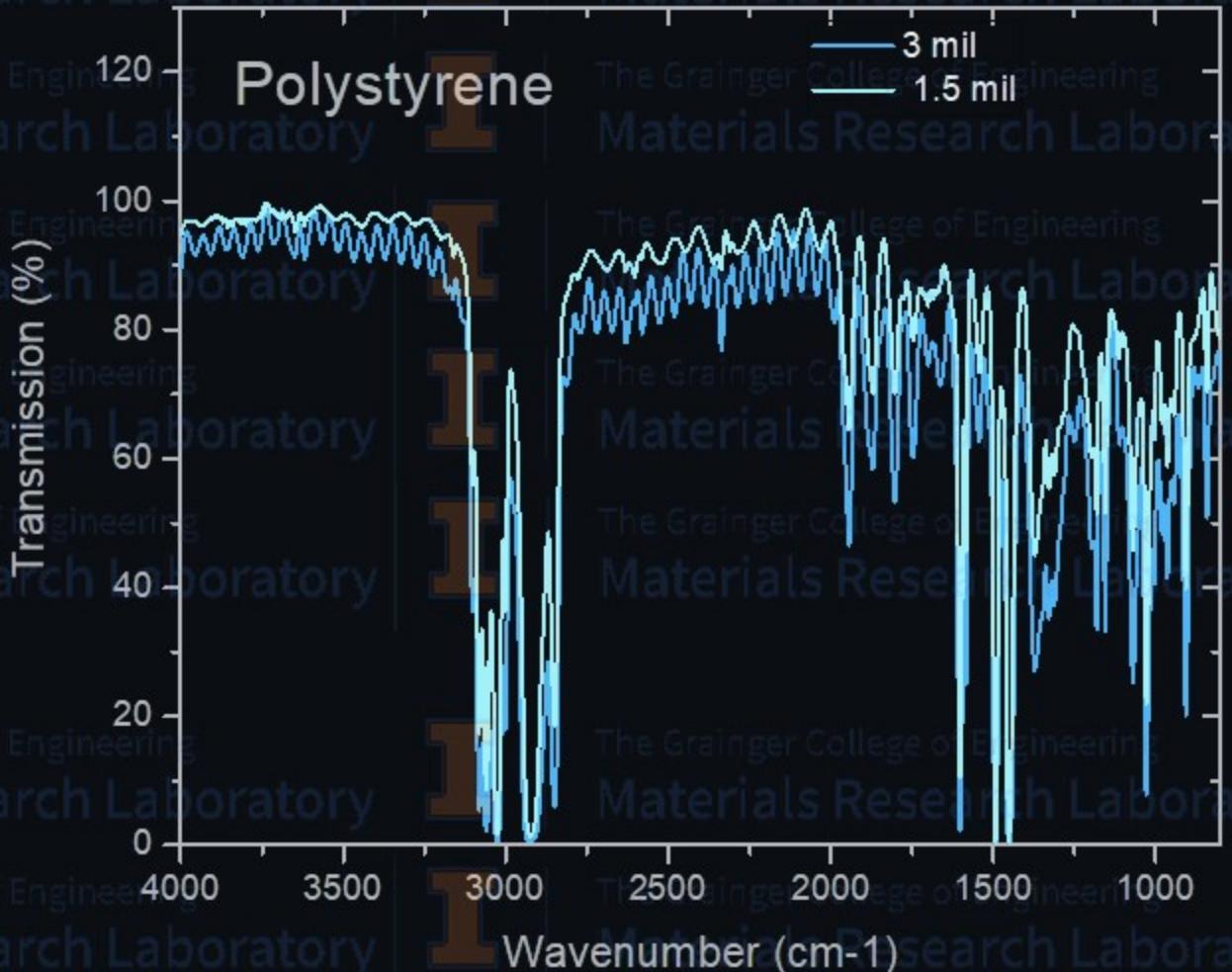


2 Hz + 3 Hz

x-coordinate of center of mass



Fourier Transform IR spectroscopy (FTIR)

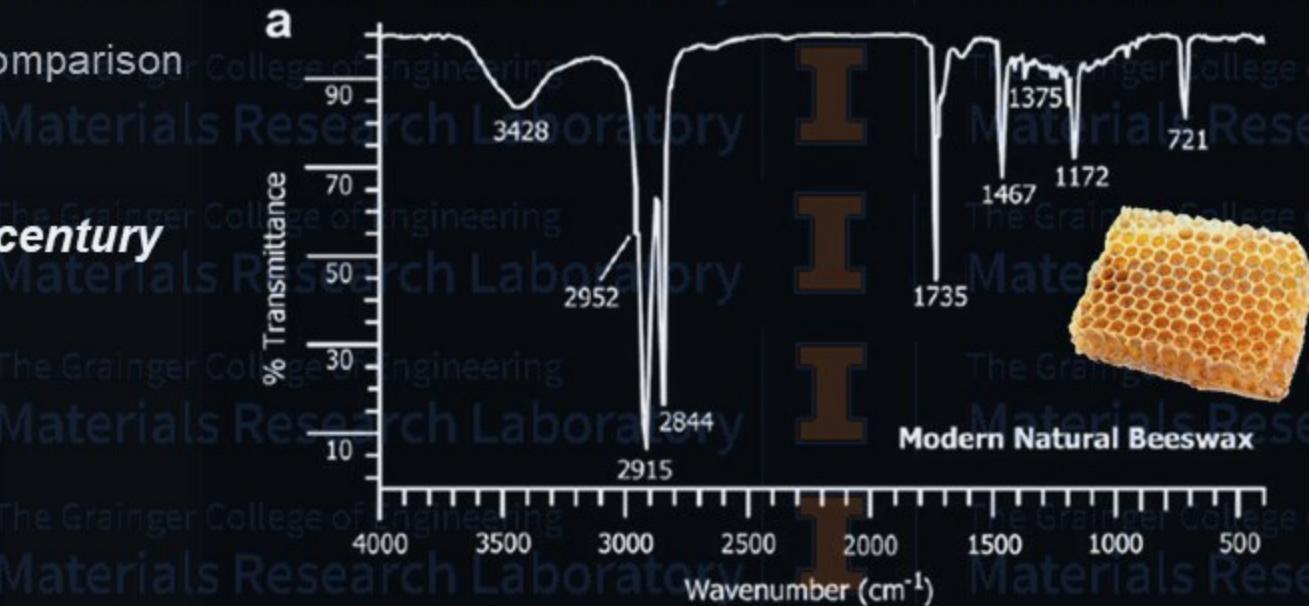
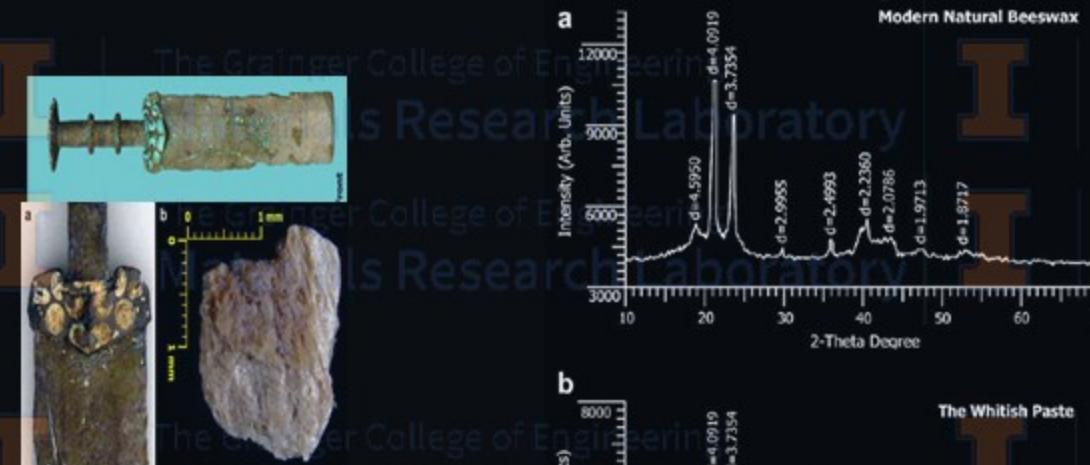


Fourier Transform IR spectroscopy (FTIR)

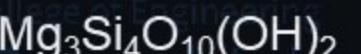
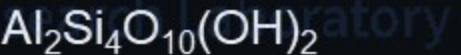
FTIR can be used to identify components in a mixture by comparison with reference spectra.

Discovery of beeswax as binding agent on a 6th-century BC Chinese turquoise-inlaid bronze sword

Wugan Luo, Tao Li, Changsui Wang, Fengchun Huang

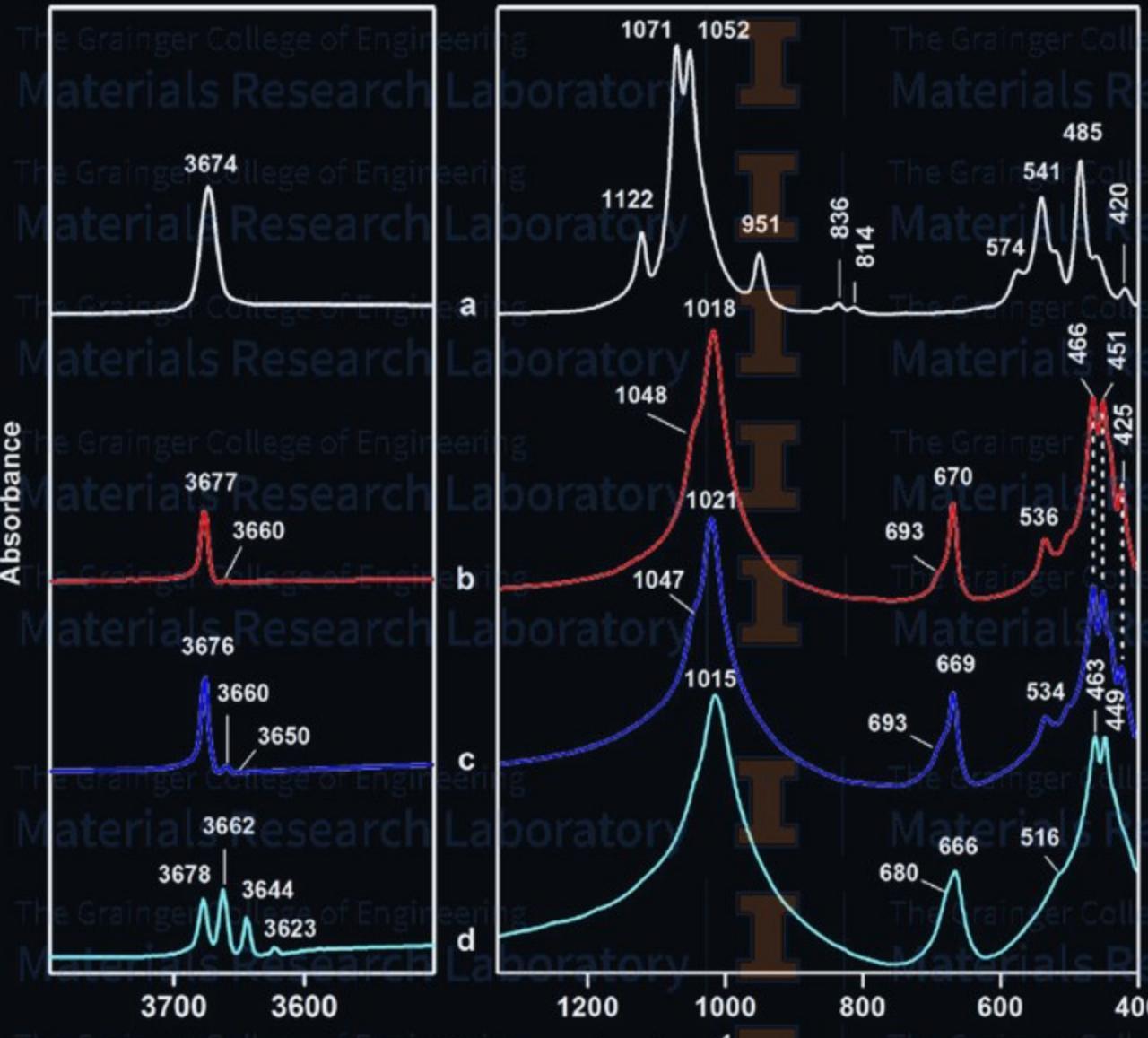


Fourier Transform IR spectroscopy (FTIR)

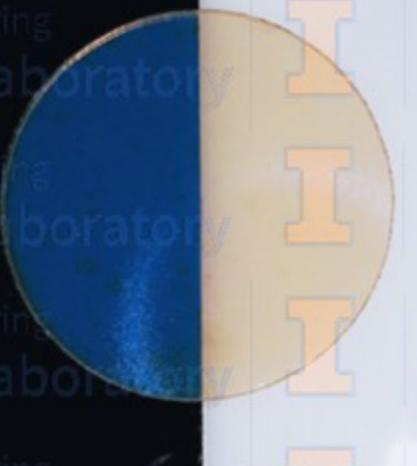


Images from Wikipedia and eurotalc.com

Spectra from *Developments in Clay Science*, Vol. 8, Ch. 5



Spectrophotometry (UV-VIS-NIR) and FTIR



Strengths:

- Very little or simple sample preparation.
- Simplicity of use and data interpretation.
- Short acquisition time, for most cases.
- Non destructive.
- Broad range of photon energies.
- High sensitivity (~ 0.1 wt% typical for FTIR).

Complementary techniques:

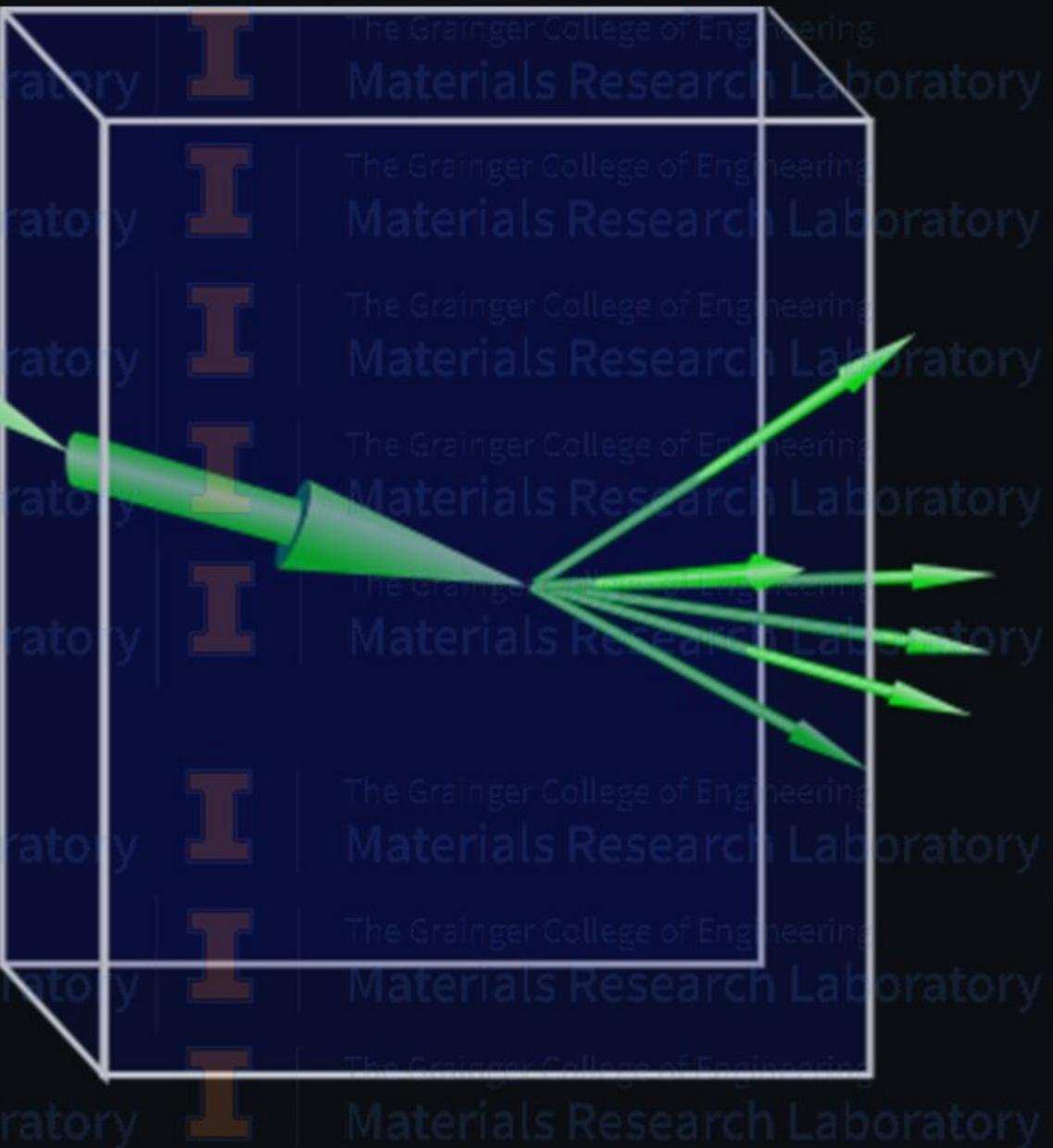
Raman, Electron Energy Loss Spectroscopy (EELS), Extended X-ray Absorption Fine Structure (EXAFS), XPS, Auger, SIMS, XRD, SFG.

Limitations:

- Reference sample is often needed for quantitative analysis.
- Many contributions to the spectrum are small and can be buried in the background.
- Usually, unambiguous chemical identification requires the use of complementary techniques.
- Limited spatial resolution.



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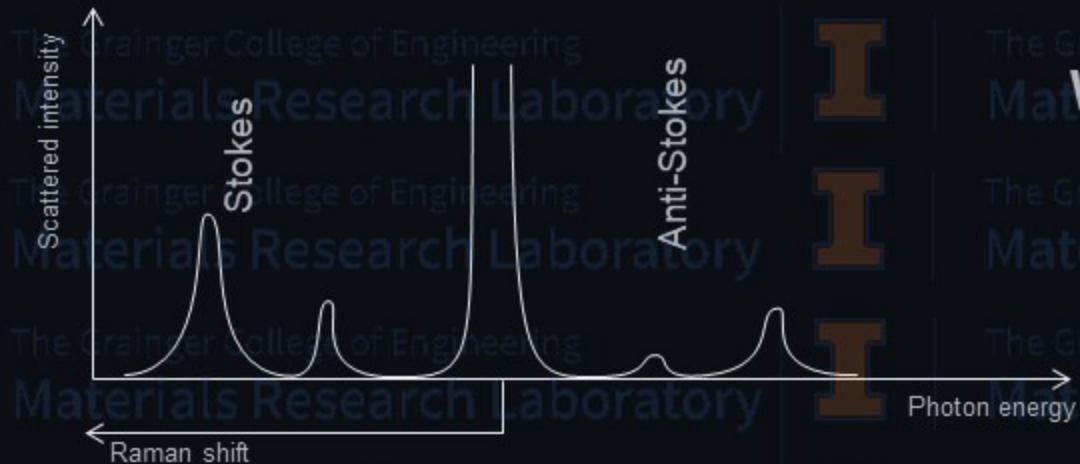
Sir Chandrasekhara Venkata Raman



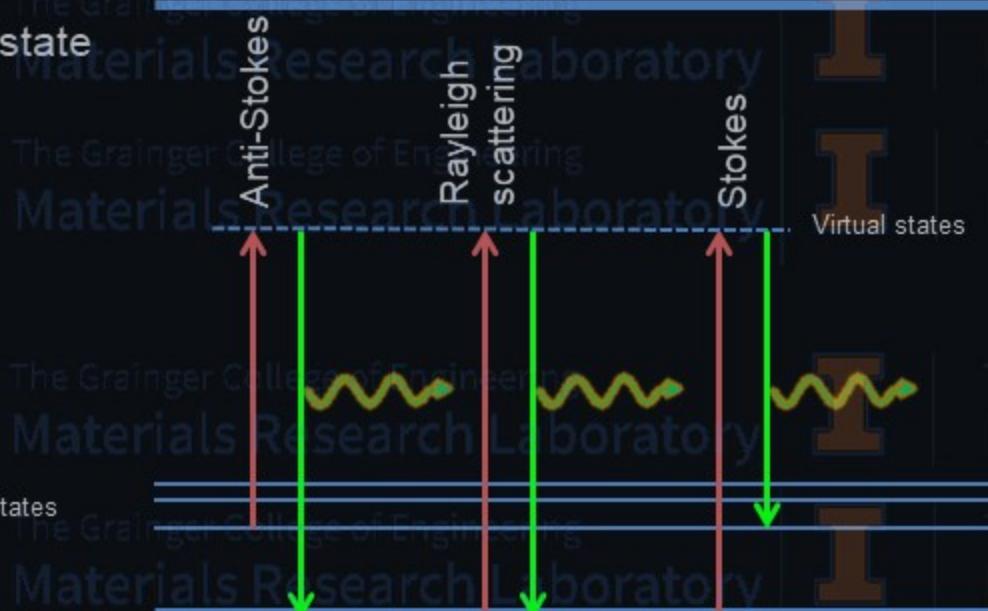
The Nobel Foundation

Sir Kariamanikkam Srinivasa Krishnan

Co-discoverer of Raman scattering, for which his mentor C. V. Raman was awarded the 1930 Nobel Prize in Physics



Excited state



Ground state

What is measured:

The light inelastically scattered by the material.

Basic principle:

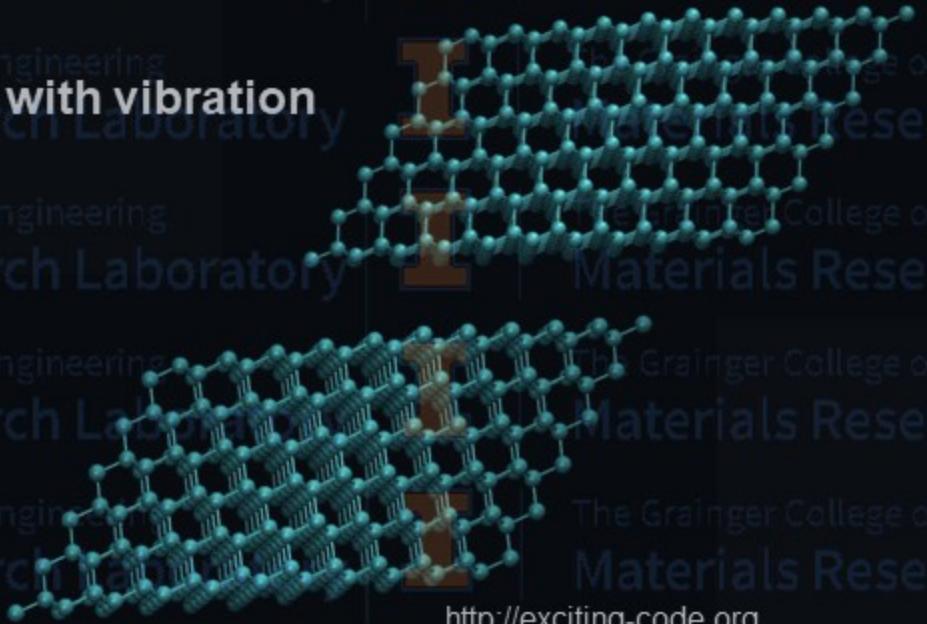
The impinging light couples with the lattice vibrations (phonons) of the material, and a small portion of it is inelastically scattered. The difference between the energy of the scattered light and the incident beam is the energy absorbed or released by the phonons.



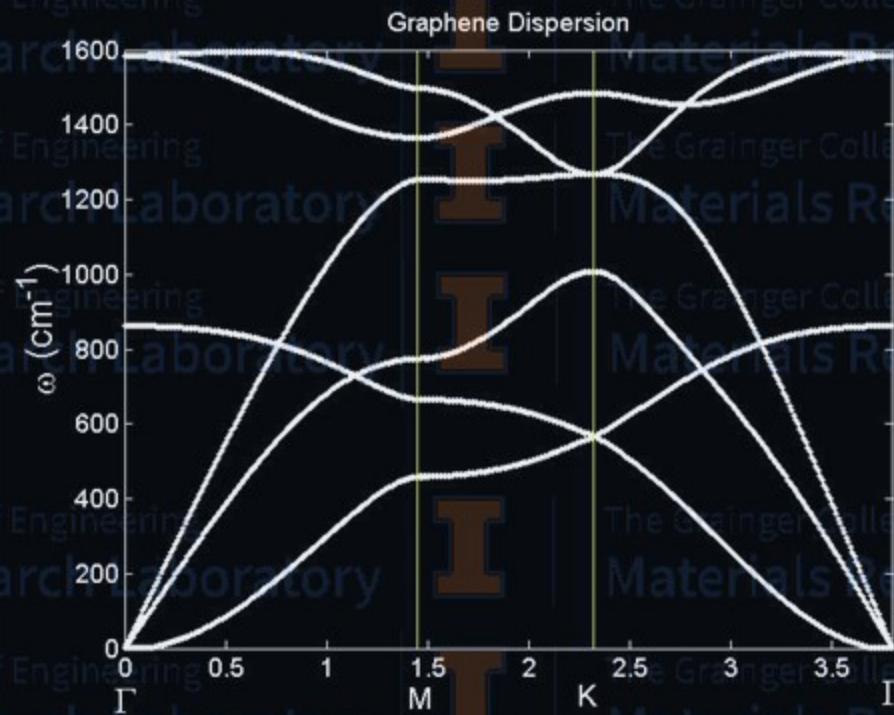
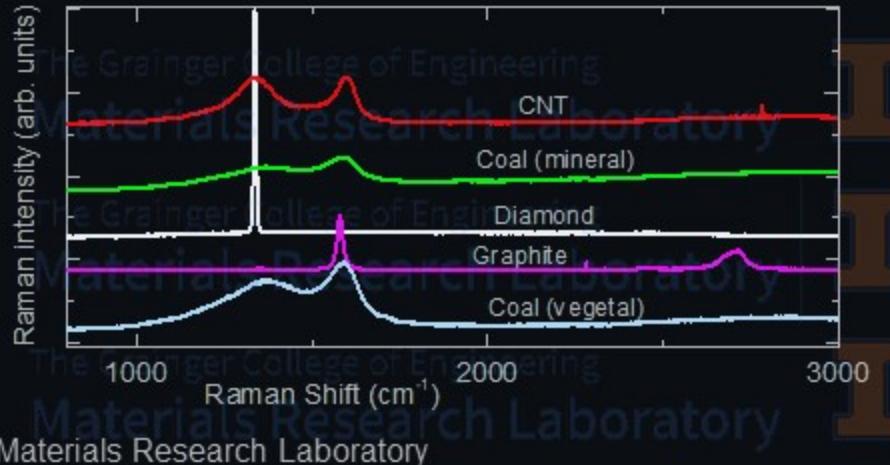
Resonance Raman {
IR}

Impinging light couples with vibration modes of the material:

- Phonons
- Molecular vibrations



<http://exciting-code.org>

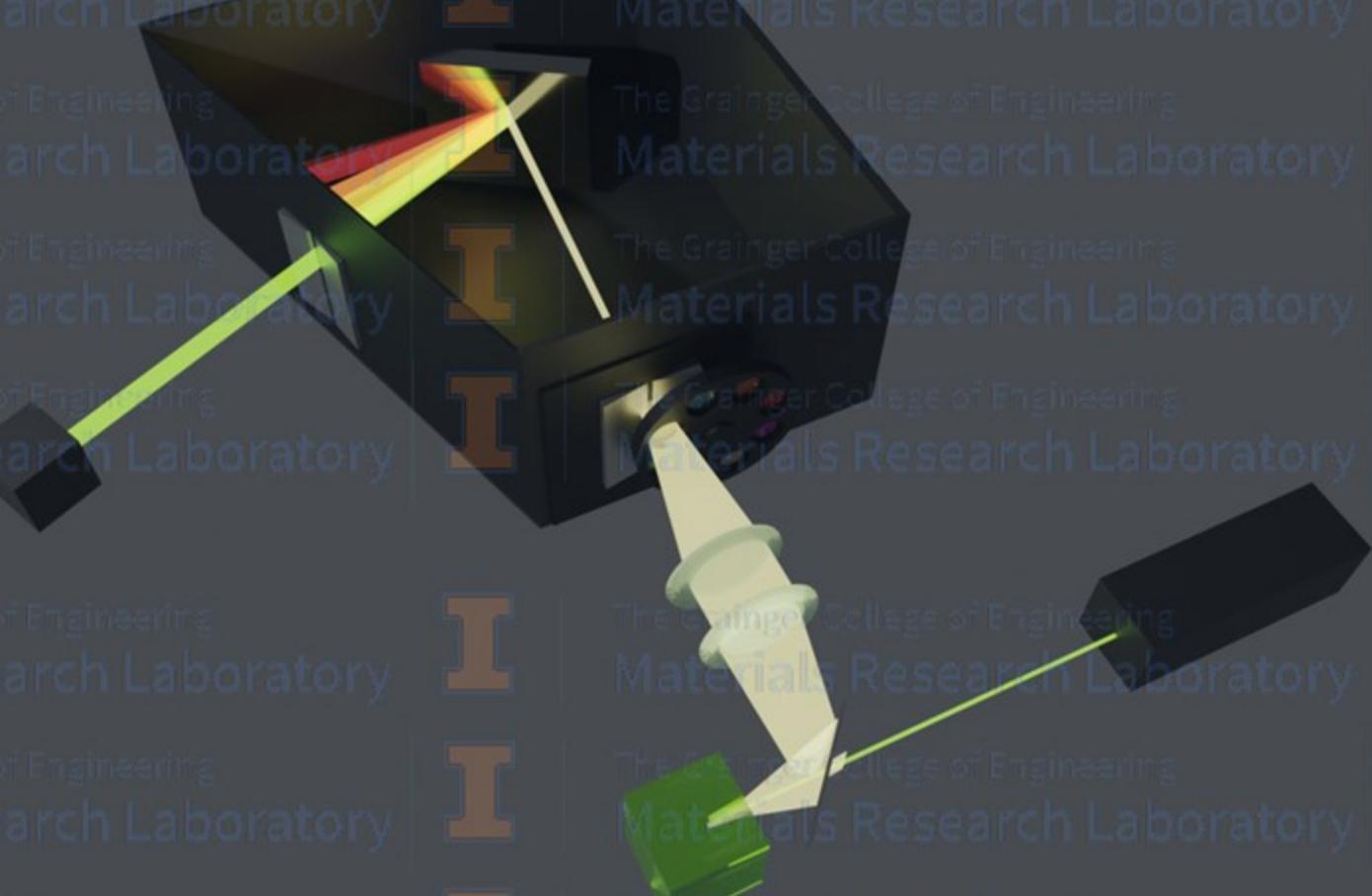


<http://flex.phys.tohoku.ac.jp/~pourya>

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Instrumentation: Materials Research Laboratory

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Raman spectroscopy
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p = 0
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Materials Research

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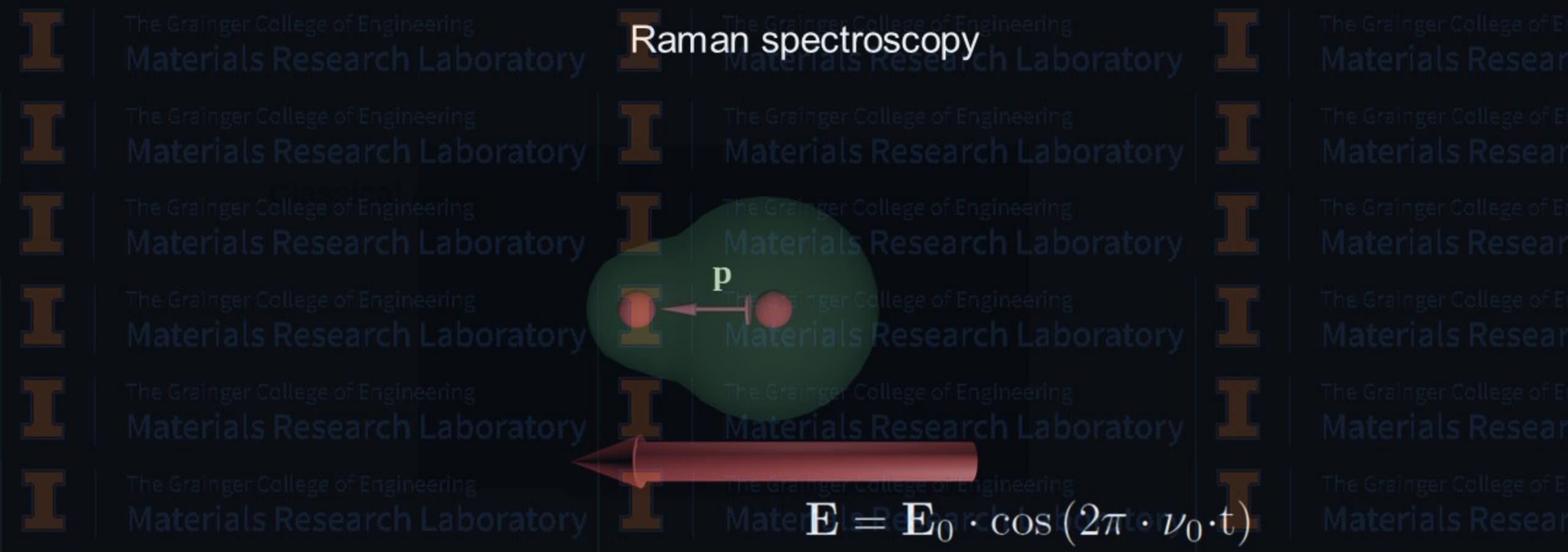
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$$p = \alpha \cdot E_0 \cdot \cos(2\pi \cdot v_0 \cdot t)$$

The α tensor is dependent on the shape, strength, and dimensions of the chemical bond. Since chemical bonds change during vibrations, α is dependent on the vibrations of the molecule:

$$\alpha_k = \alpha_0 + \sum_k \left(\frac{\partial \alpha}{\partial Q_k} \right)_0 \cdot Q_k + \frac{1}{2} \sum_{k,l} \left(\frac{\partial^2 \alpha}{\partial Q_k \partial Q_l} \right)_0 \cdot Q_k \cdot Q_l + \dots$$

$$Q_k = Q_{k0} \cdot \cos(2\pi \cdot \nu_k \cdot t + \varphi_k)$$

$$\alpha_k = \alpha_0 + \alpha'_k \cdot Q_{k0} \cdot \cos(2\pi \cdot \nu_k \cdot t + \varphi_k)$$

The α tensor is dependent on the shape, strength, and dimensions of the chemical bond. Since chemical bonds change during vibrations, α is dependent on the vibrations of the molecule:

$$\alpha_k = \alpha_0 + \alpha'_k \cdot Q_{k0} \cdot \cos(2\pi \cdot \nu_k \cdot t + \varphi_k)$$

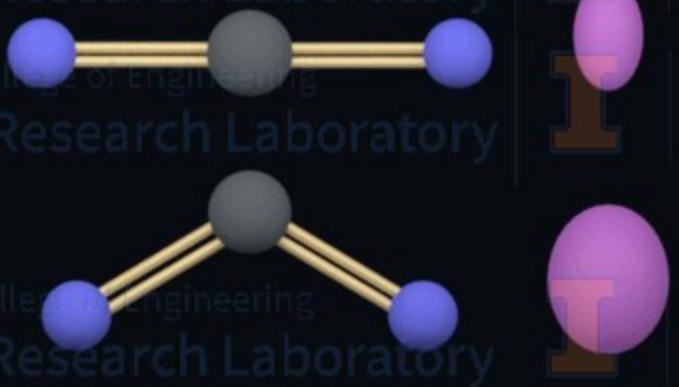
$$p = \alpha_0 \cdot E_0 \cdot \cos(2\pi \cdot \nu_0 \cdot t) +$$

$$+ \frac{1}{2} \cdot \alpha'_k \cdot Q_{k0} \cdot E_0 \cdot [\cos(2\pi \cdot t \cdot (\nu_0 + \nu_k) + \varphi_k) + \\ + \cos(2\pi \cdot t \cdot (\nu_0 - \nu_k) - \varphi_k)]$$

$$\alpha'_k = \left(\frac{\partial \alpha}{\partial Q_k} \right)_0 \neq 0$$

the dipole oscillates with three frequencies simultaneously, corresponding to the three possible scattering modes (Rayleigh, Stokes Raman and anti-Stokes Raman)

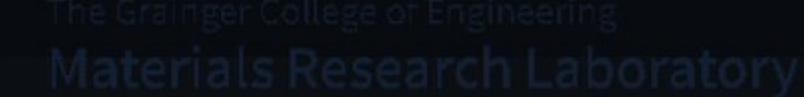
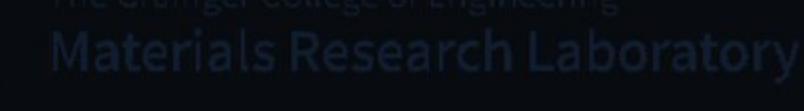
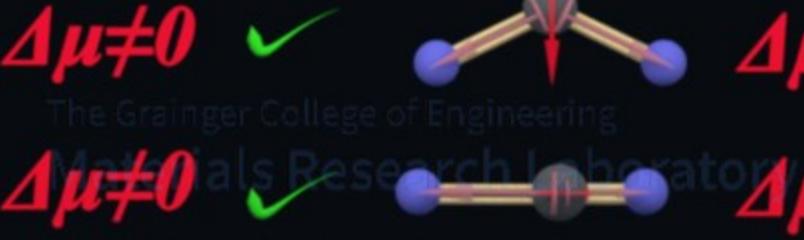
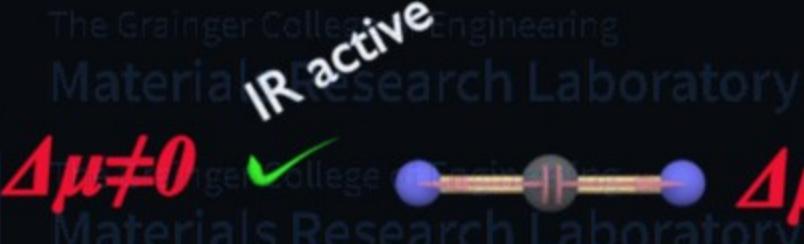
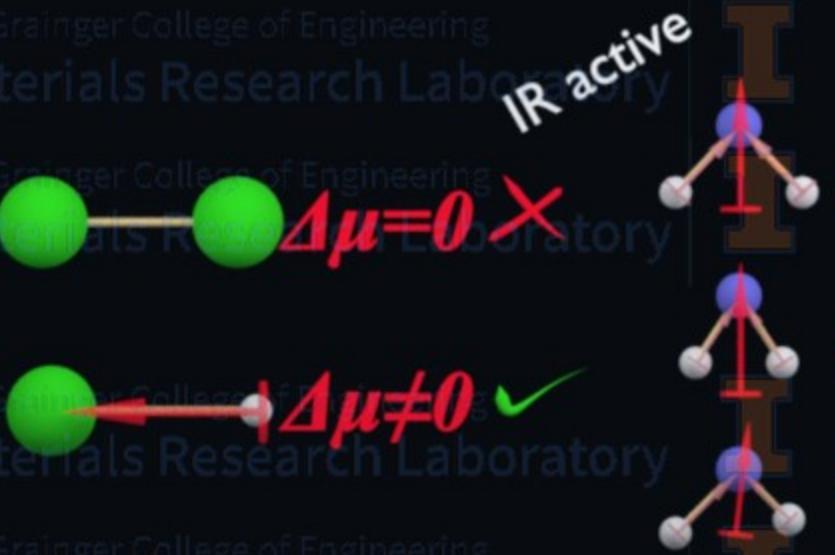
α'_k



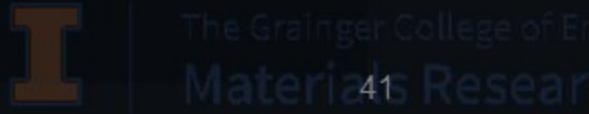
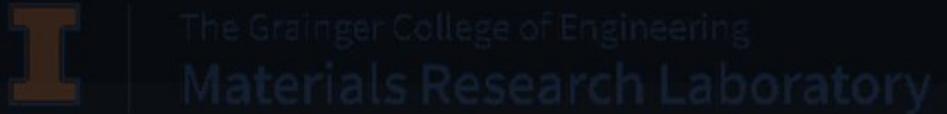
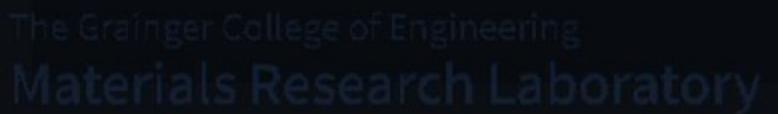
$$\frac{\partial \alpha}{\partial Q_k} \Big|_0 \neq 0$$



IR active vibrations



Raman active vibrations

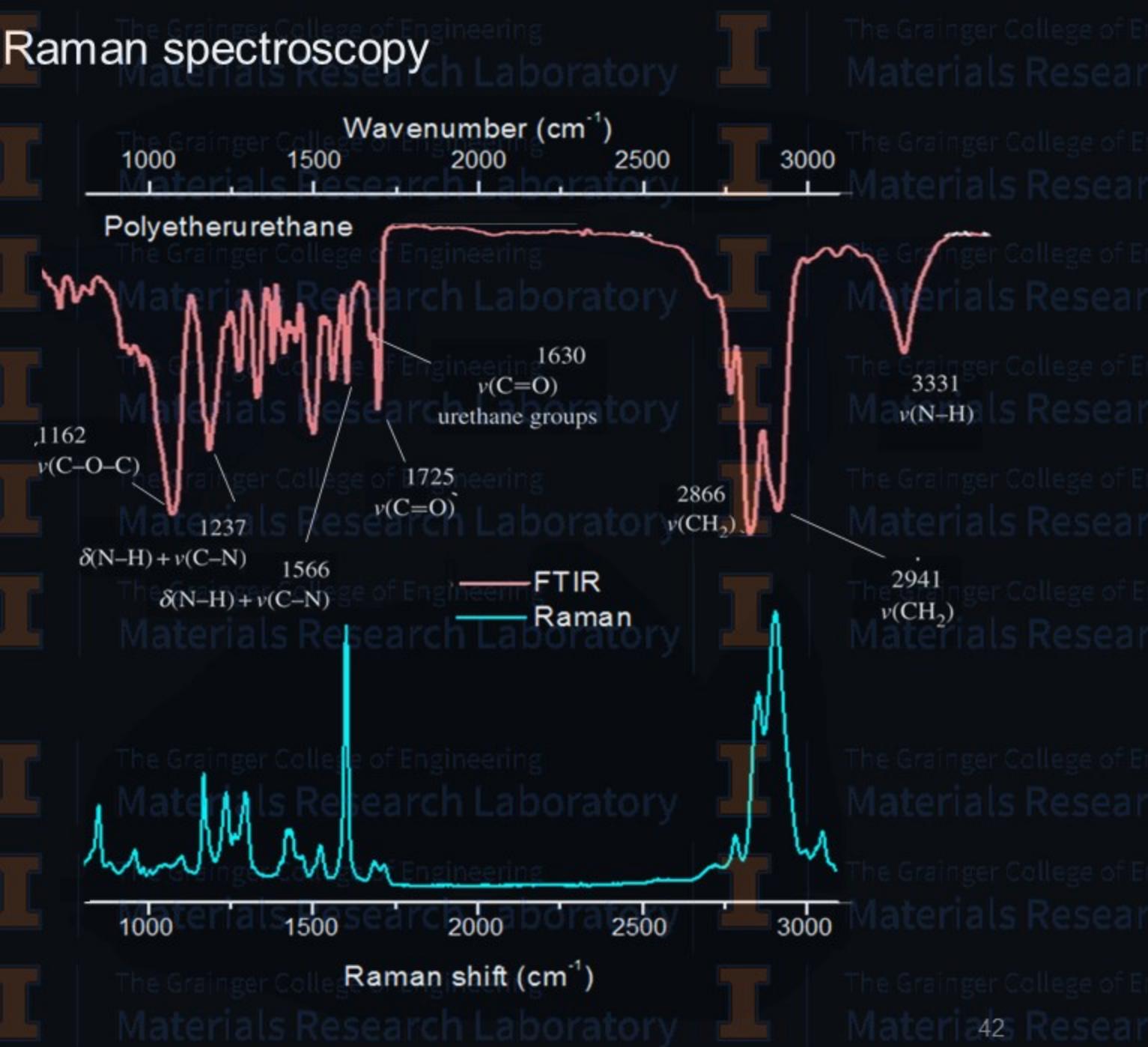


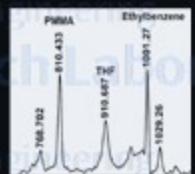
FTIR and Raman:

The two techniques are complementary
(different selection rules).

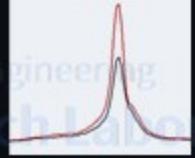
$$\Delta\mu \neq 0$$

$$\alpha'_k = \left(\frac{\partial \alpha}{\partial Q_k} \right)_0 \neq 0$$

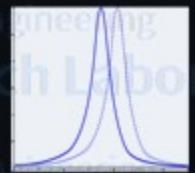




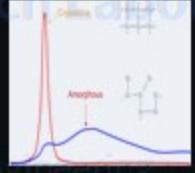
Characteristic Raman
frequencies



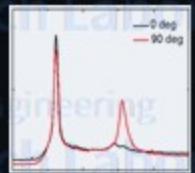
Raman peak intensity



Raman peak frequency
shift



Raman peak width



Raman peak polarization
dependency

... we can estimate ...

Identity and composition
of materials

Volume of material
probed

Strain, stress, crystal
lattice distortion

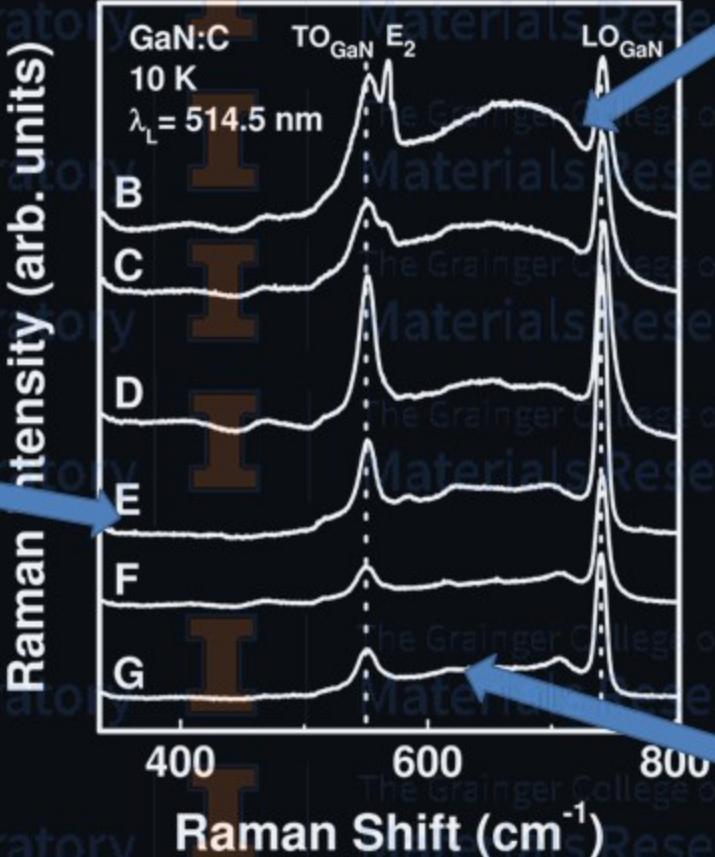
Crystallinity of material

Crystal orientation and
symmetry

Molecular and crystalline structure characterization

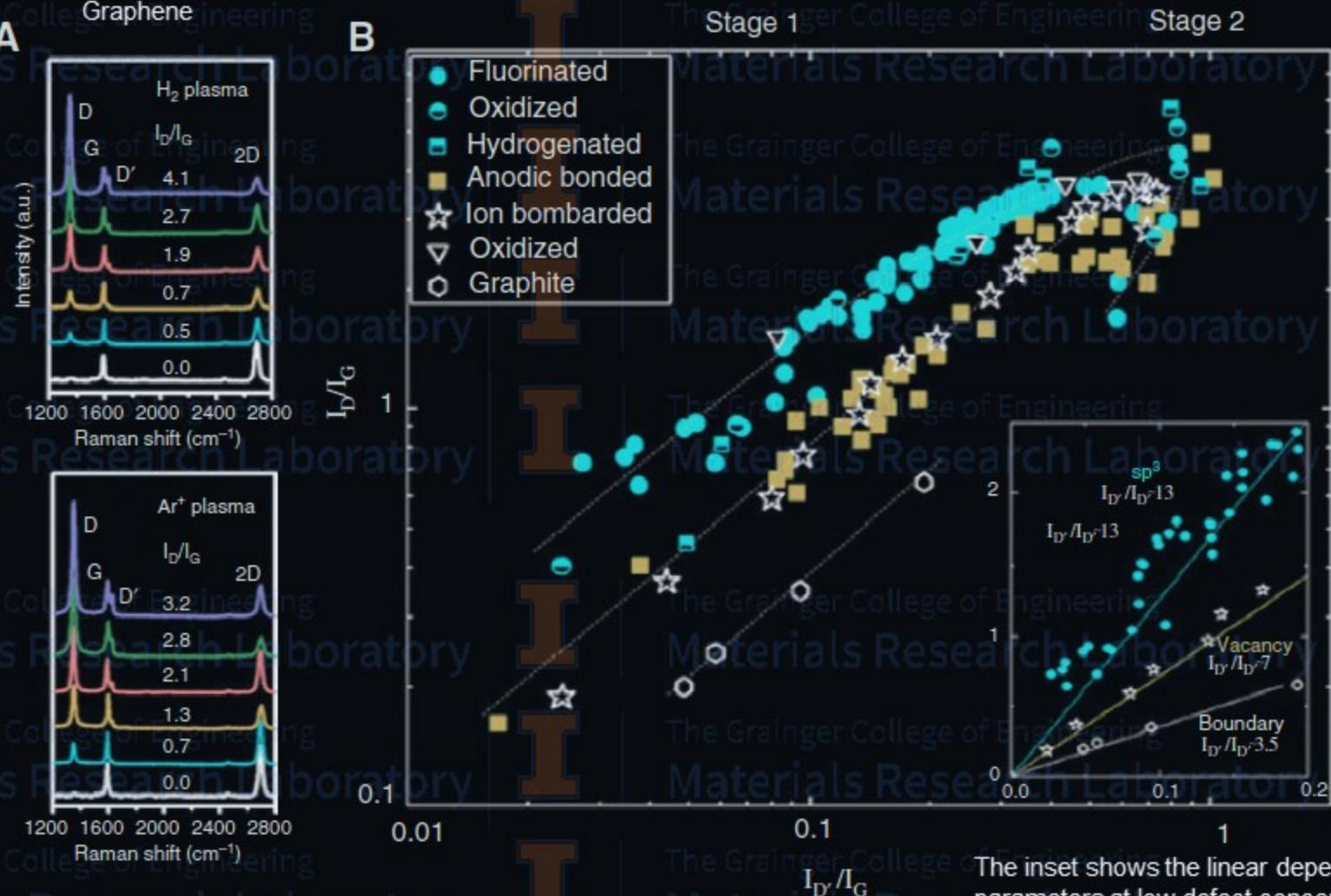
Presence of N vacancies yields poor crystallinity

Substitutional C fills N vacancies improving the crystallinity

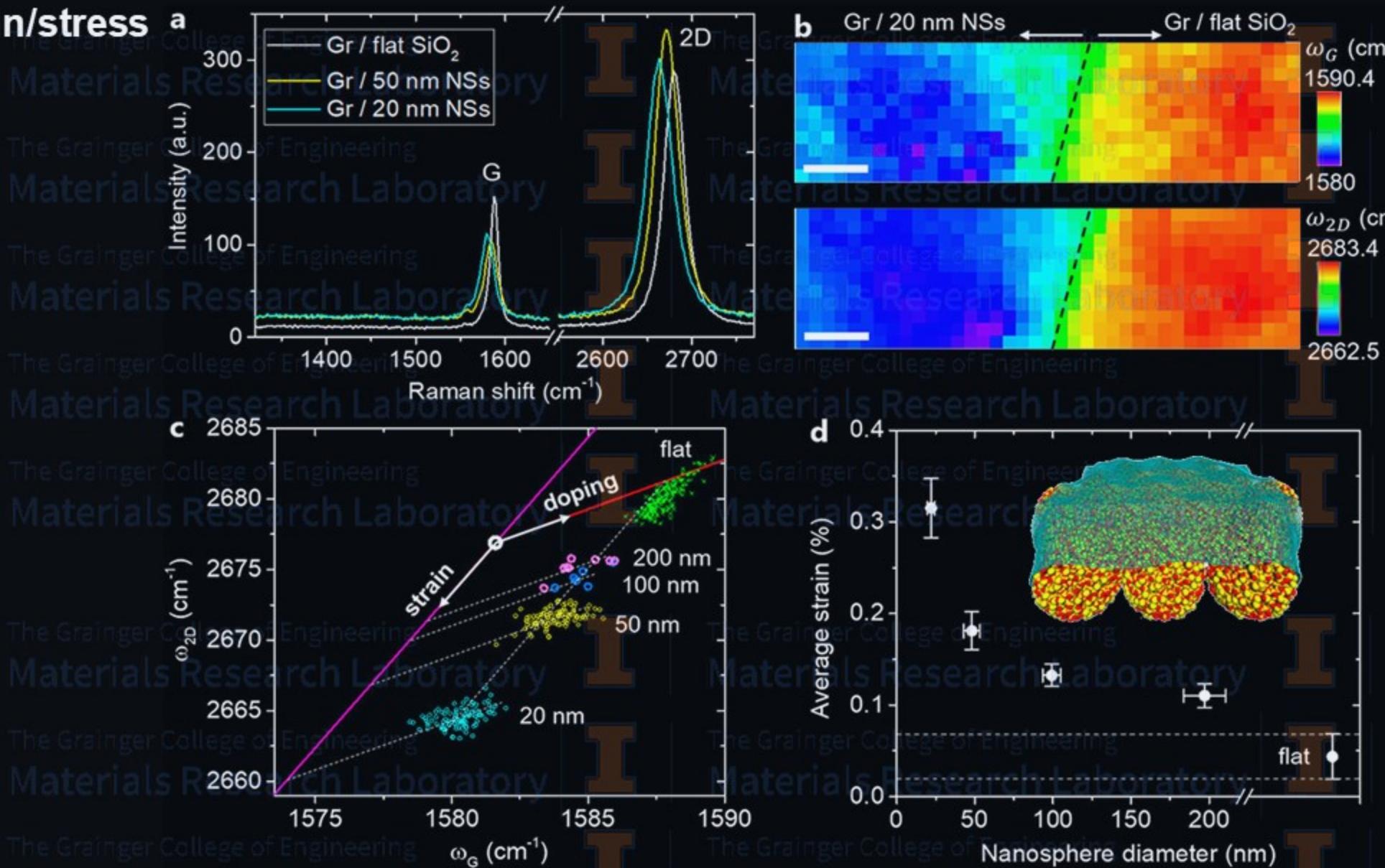


C incorporates interstitially causing a degradation of the crystal lattice

Crystalline structure and defect characterization



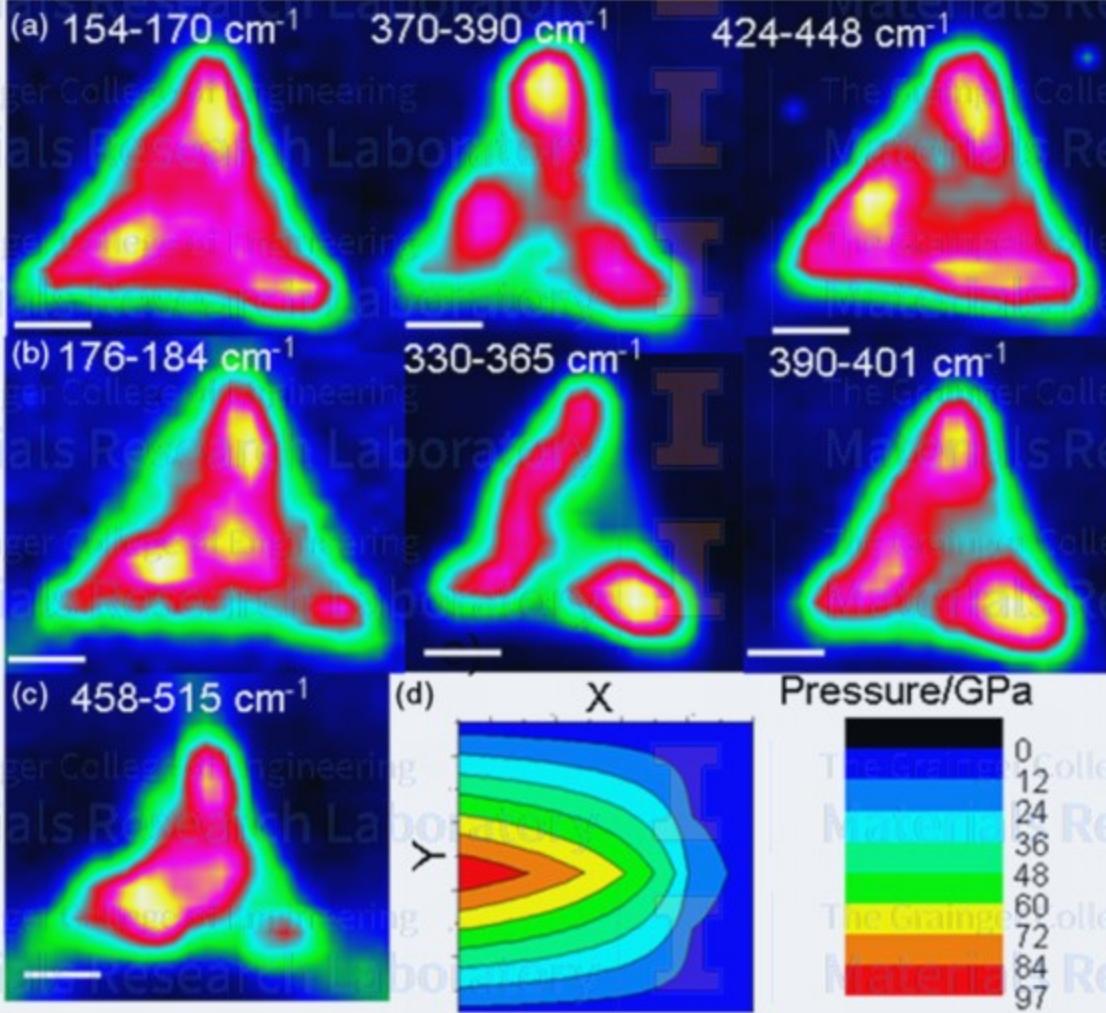
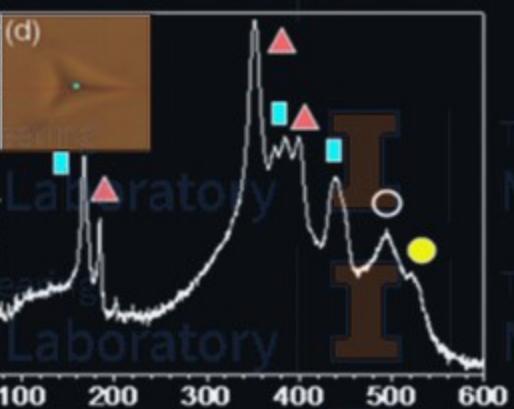
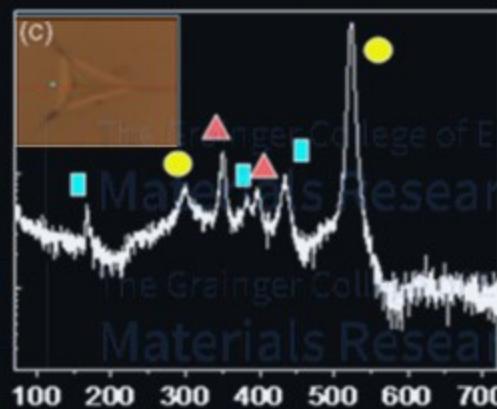
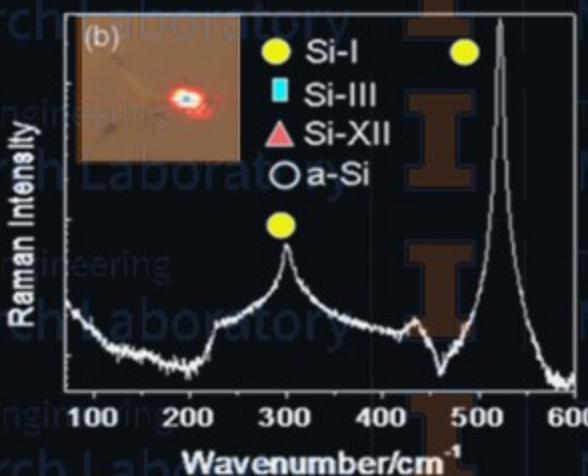
Strain/stress



Phase transitions

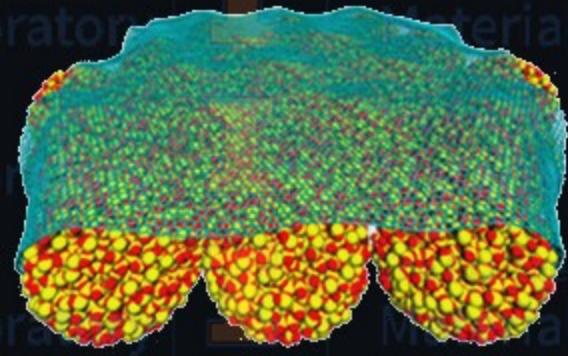
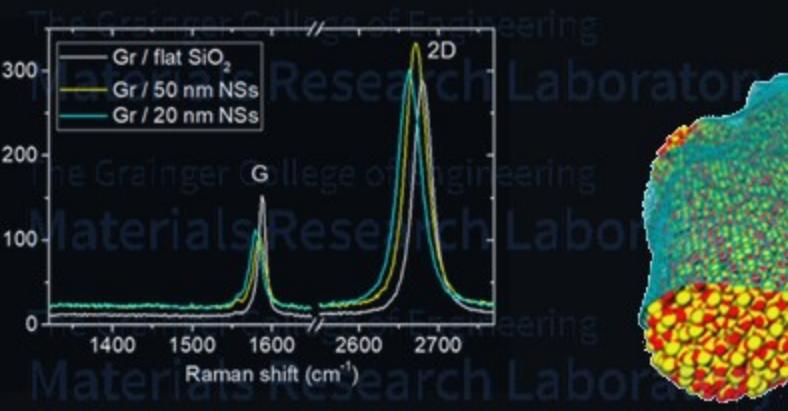
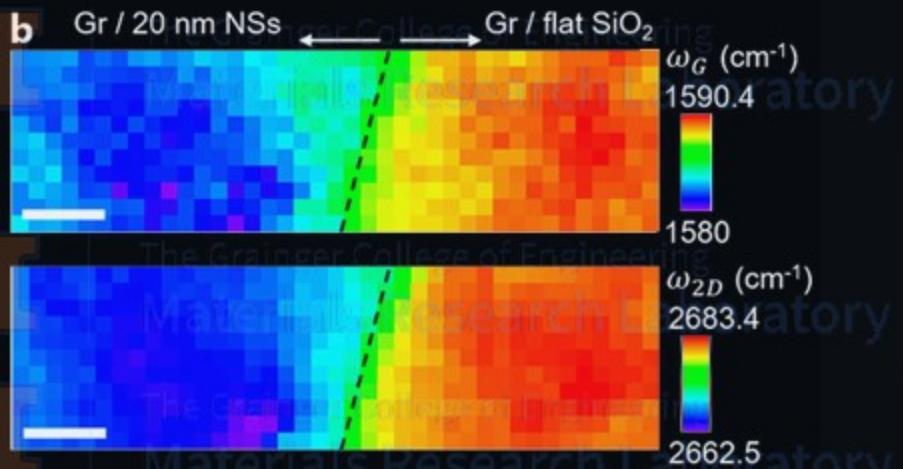
(a) Different Raman Modes

Si-I (cm ⁻¹) [†]	Si-III (cm ⁻¹) [†]	Si-XII (cm ⁻¹) [†]
300, 520	166, 171	182
a-Si (cm ⁻¹) [†]	384	352
475, 510	432, 463	397



Primary Strengths:

- Very little sample preparation.
- Structural characterization.
- Non destructive technique.
- Chemical information.
- Complementary to FTIR.



Primary Limitations:

- Expensive apparatus (for high spectral/spatial resolution and sensitivity).
- Weak signal, compared to fluorescence.
- Limited spatial resolution (diffraction limited).

Complementary techniques:

FTIR, EELS, Mass spectroscopy, EXAFS, XPS, AES, SIMS, XRD, SFG.

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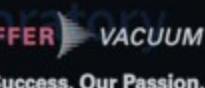
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Optical Characterization Methods Part II

Julio A. N. T. Soares

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University of Illinois at Urbana-Champaign



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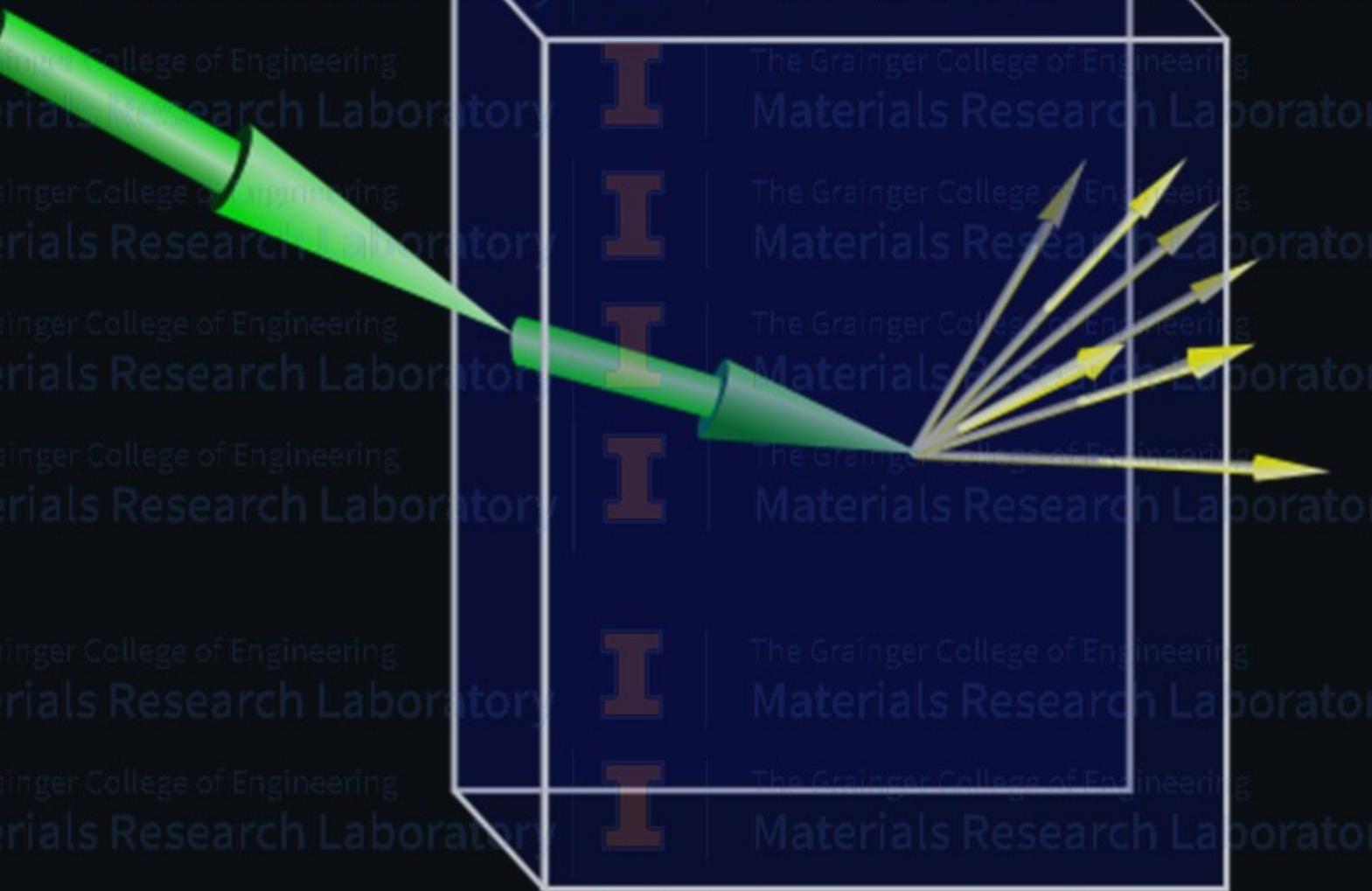


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**Lifetime: Phosphorescence, fluorescence
Mechanism: Photoluminescence, bioluminescence, chemoluminescence,
thermoluminescence, piezoluminescence, etc.**



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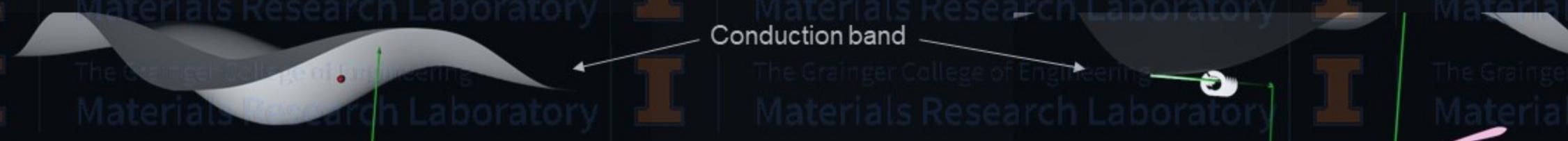
Trevor Morris

52

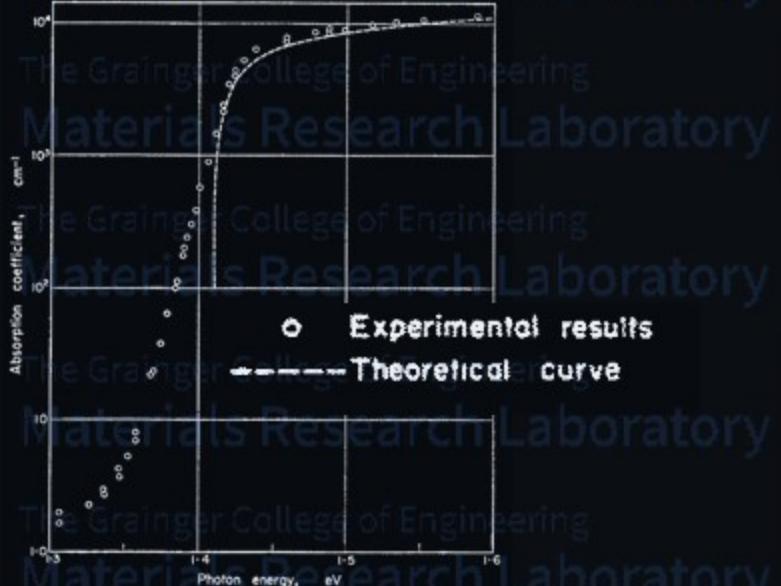
What is measured:

The emission spectra of materials due to radiative recombination following photo-excitation.

Basic principle:

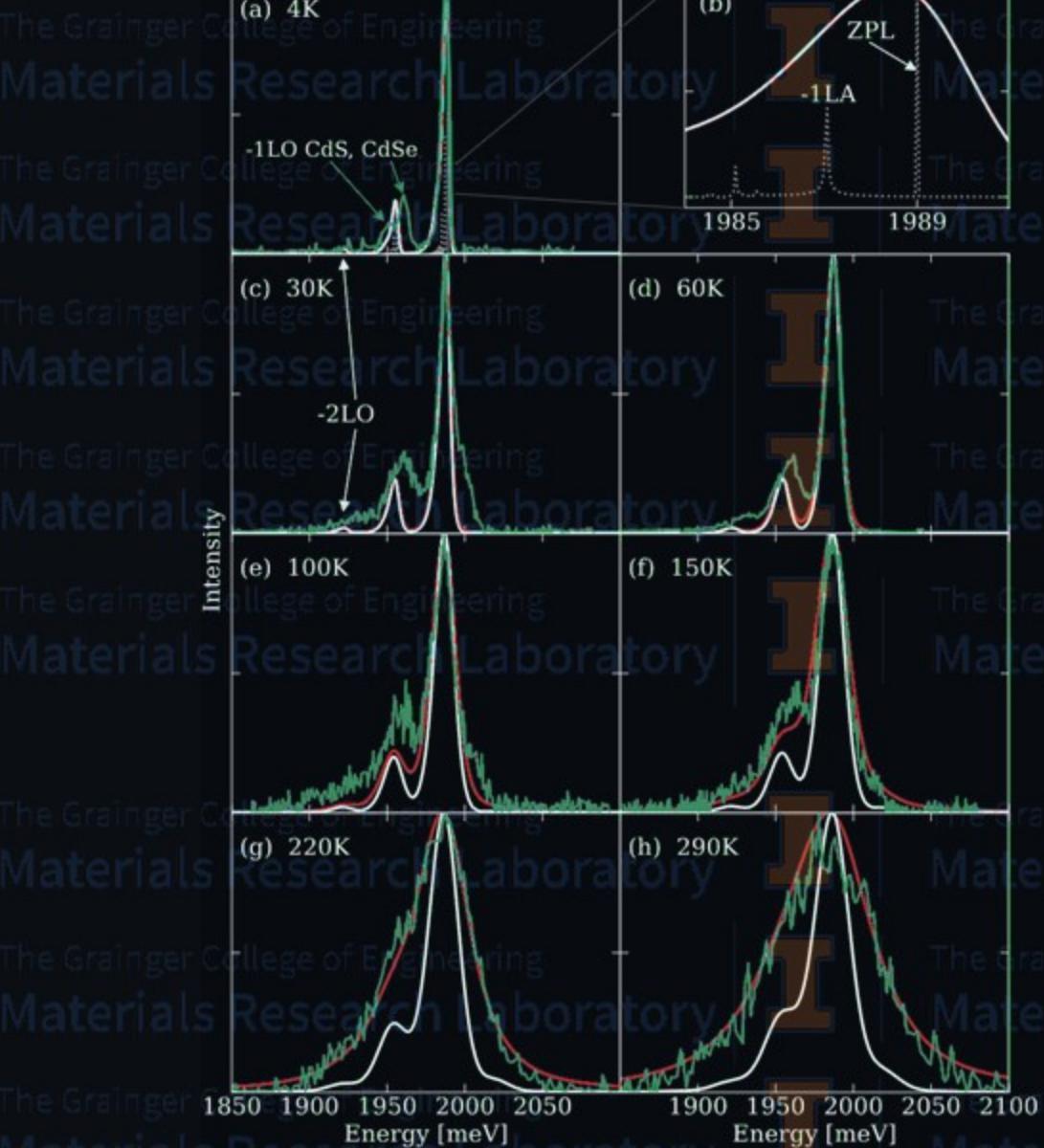






Room temperature absorption in GaAs
Infrared Physics, 1961, Vol. 1, pp. 111

Photoluminescence



Photoluminescence

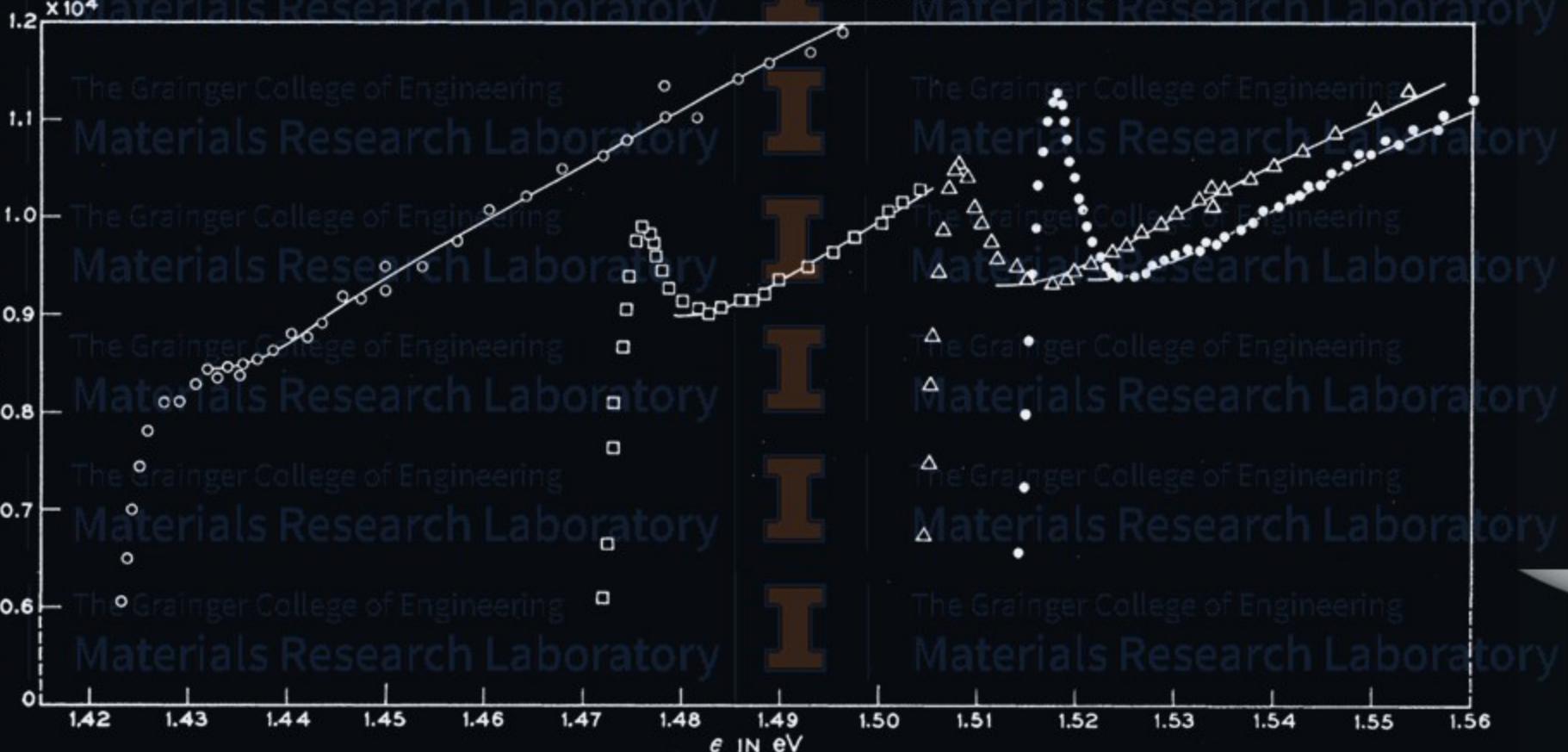


FIG. 3 Exciton absorption in GaAs; \circ 294°K, \square 186°K, Δ 90°K, \bullet 21°K.

Lattice constant

Photoluminescence

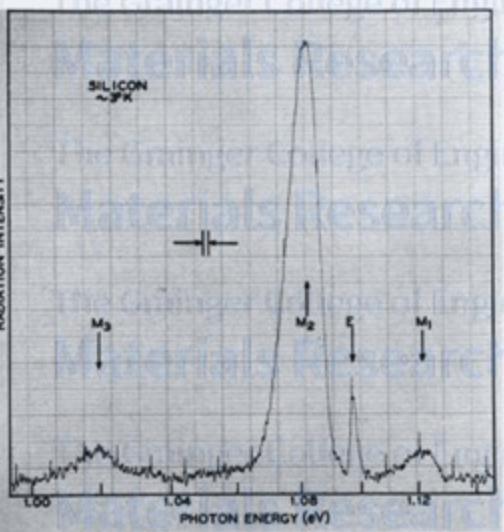


FIG. 1. Spectrogram of a Si specimen at $\sim 3\text{K}$. The horizontal axis is the energy of the emitted photons in eV. The vertical response is nearly proportional to the number of photons per unit energy interval. The specimen resistivity at room temperature was $9 \times 10^3 \text{ }\Omega \text{ cm}$.

Excitonic molecule emission in Si

Phys. Rev. Lett. 17, 860 (1966)

PHOTON ENERGY (eV)
1.52 1.51 1.50 1.49 1.48

(D°,X) (D°,X)

(A°,X)

P = 17 Torr
n = $2.3 \times 10^{14} \text{ cm}^{-3}$
 $\mu_{77} = 105,000 \text{ cm}^2/\text{V}\cdot\text{s}$

(e,C°) (D°,C°)

PL INTENSITY (arb. unit)

820 830 840
WAVELENGTH (nm)

Jap. J. App. Phys. 23, L100 (1984)

Defect emission in GaAs



Conduction band

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Valence band

A

D

(D^+, X)

(A°,X)

FE

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(D°,C°)

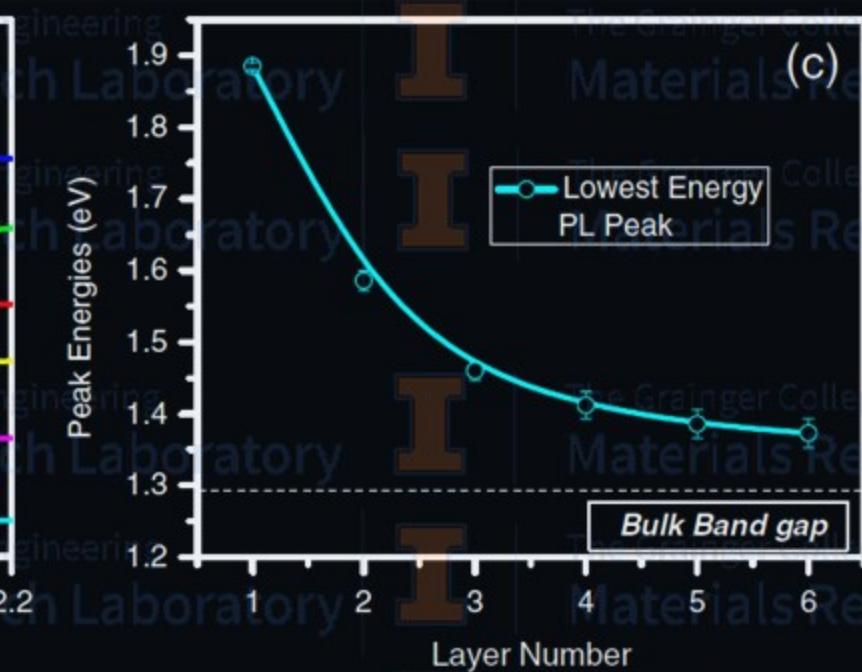
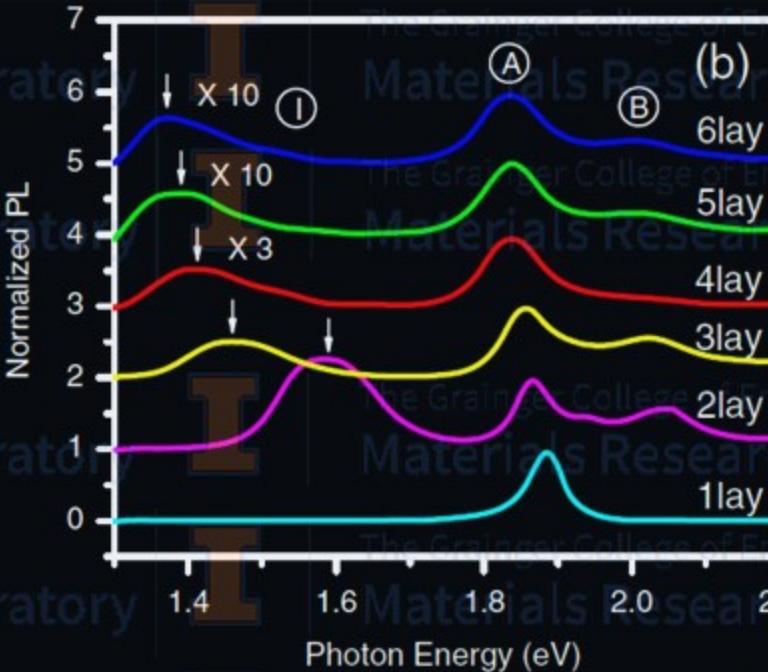
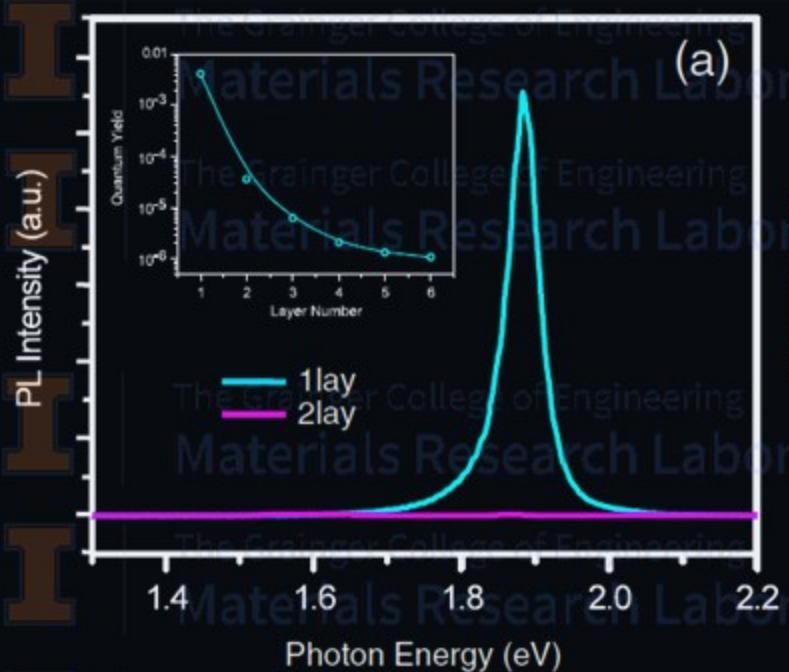
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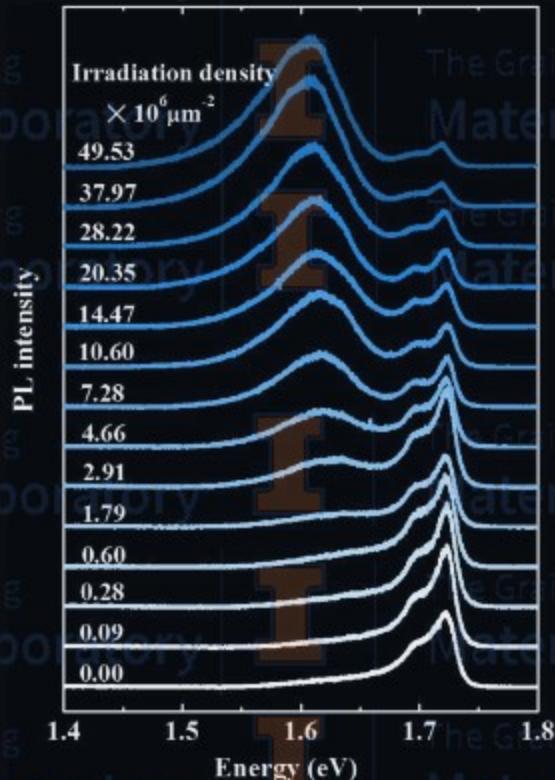
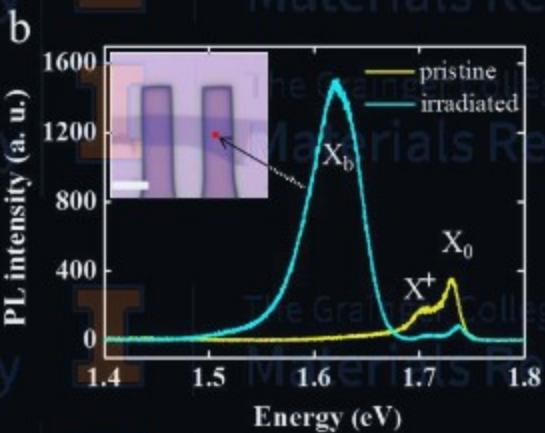
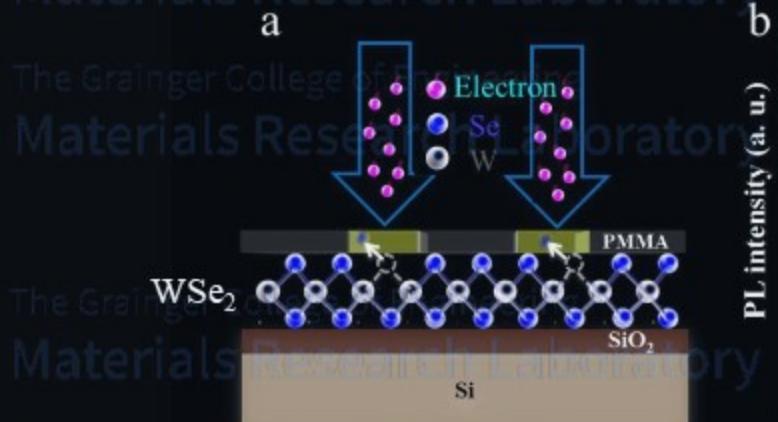
Number of layers in 2D materials

- a) PL spectra for mono- and bilayer MoS₂. Inset: PL QY of thin layers for N = 1–6.
- b) Normalized PL spectra by the intensity of peak A of thin layers of MoS₂ for N = 1–6. Feature I for N = 4–6 is magnified for clarity.
- c) Band-gap energy of thin layers of MoS₂, inferred from the energy of the PL feature I for N = 2–6 and from the energy of the PL peak A for N = 1. The dashed line represents the (indirect) band-gap energy of bulk MoS₂.



Defects in 2D materials

- Schematic diagram of electron beam irradiation on monolayer WSe₂ sample during the EBL process.
- PL spectrum of pristine monolayer WSe₂ and monolayer WSe₂ after EBL.
The inset shows optical image of WSe₂ with PMMA patterned by EBL, scale bar is 5 μm
- PL spectra of a pristine WSe₂ under different e⁻ beam irradiation density.



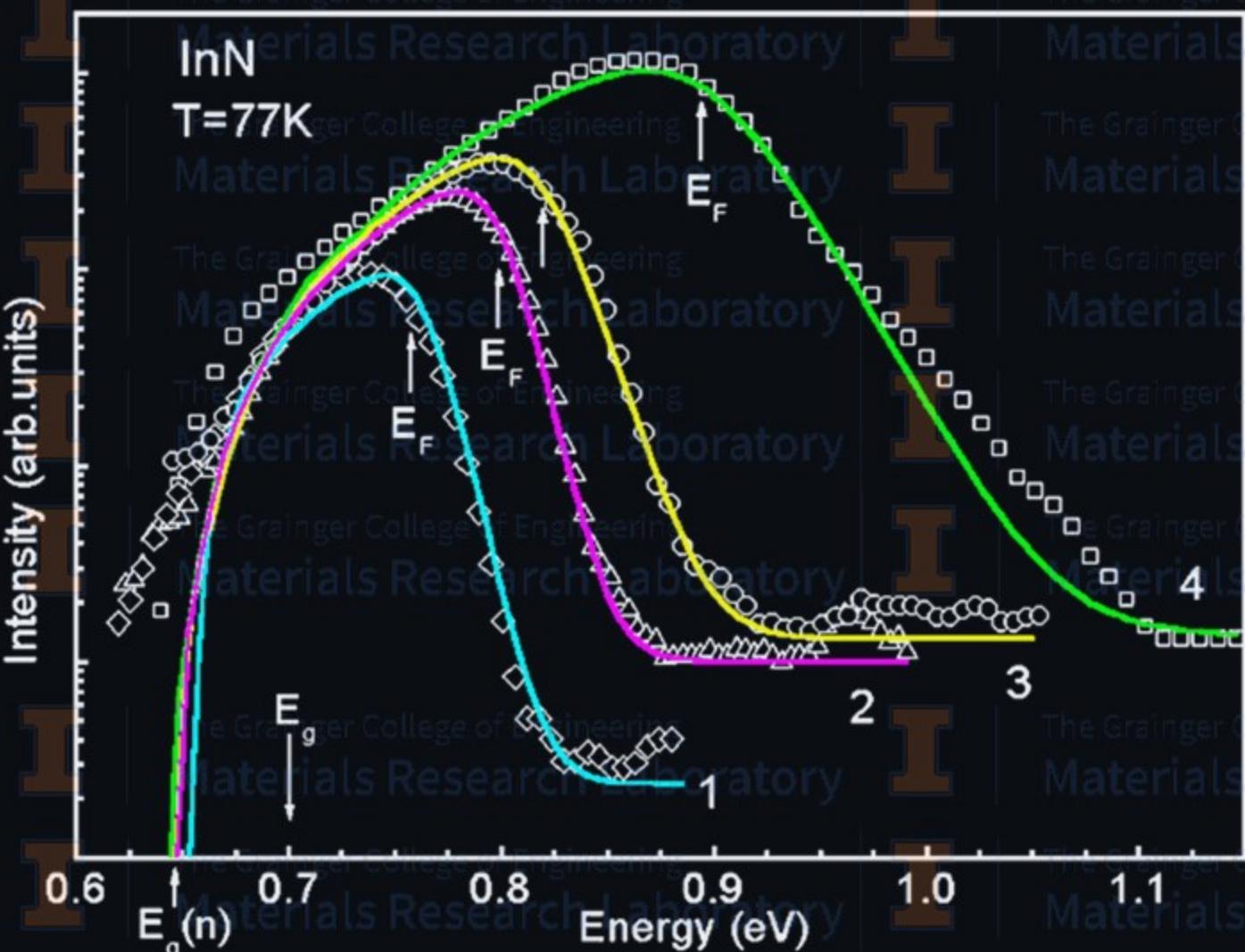
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Carrier concentration

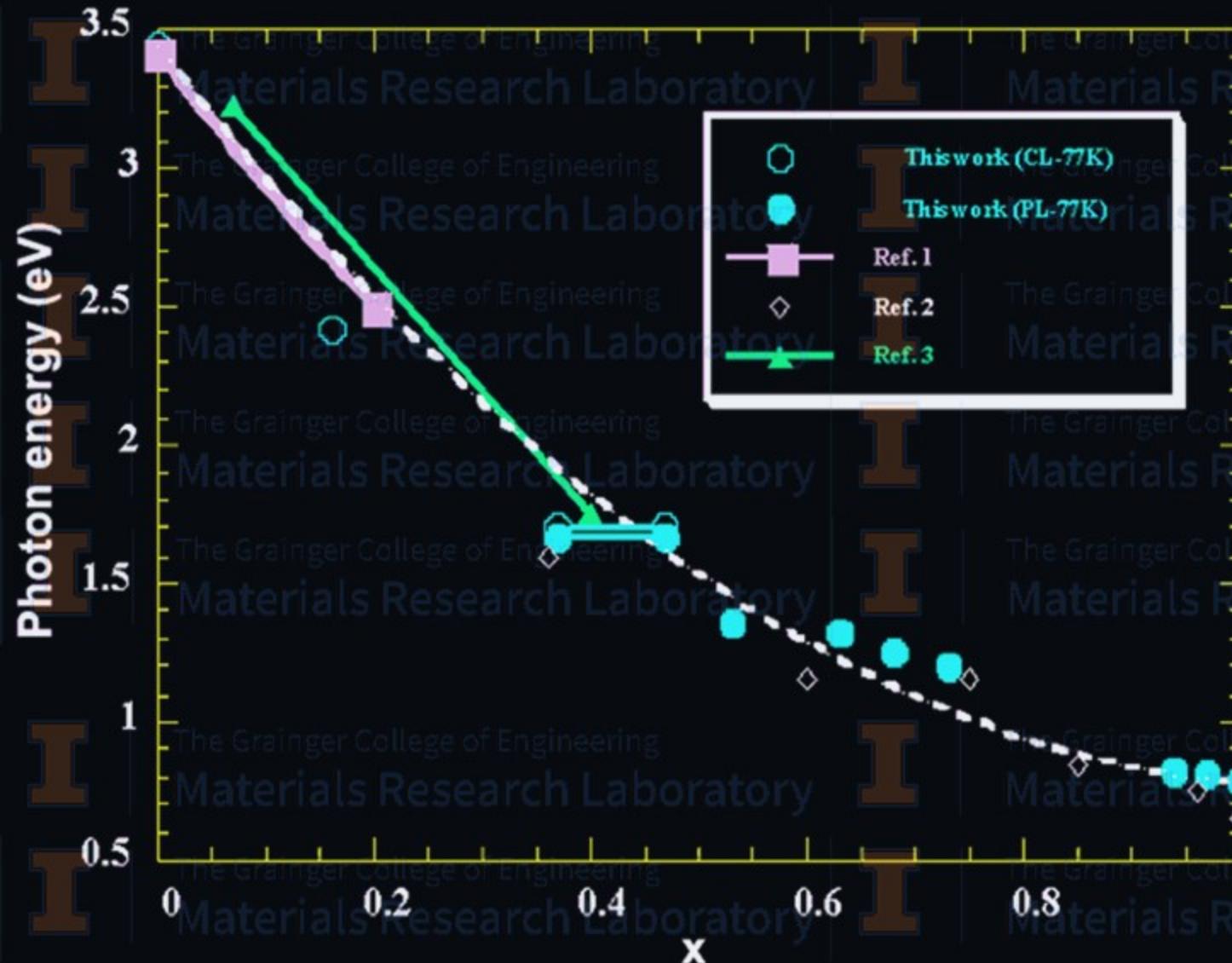
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- 1 - $n = 6 \times 10^{18} \text{ cm}^{-3}$ (MOCVD); Engineering
- 2 - $n = 9 \times 10^{18} \text{ cm}^{-3}$ (MOMBE);
- 3 - $n = 1.1 \times 10^{19} \text{ cm}^{-3}$ (MOMBE);
- 4 - $n = 4.2 \times 10^{19} \text{ cm}^{-3}$ (PAMBE)



Alloy composition



In_xGa_{1-x}N alloys. Luminescence peak positions of catodoluminescence and photoluminescence spectra vs. concentration x.

The plots of luminescence peak positions can be fitted to the curve

$$E_g(x) = 3.48 - 2.70x - bx(1-x)$$

with a bowing parameter of **b=2.3 eV**

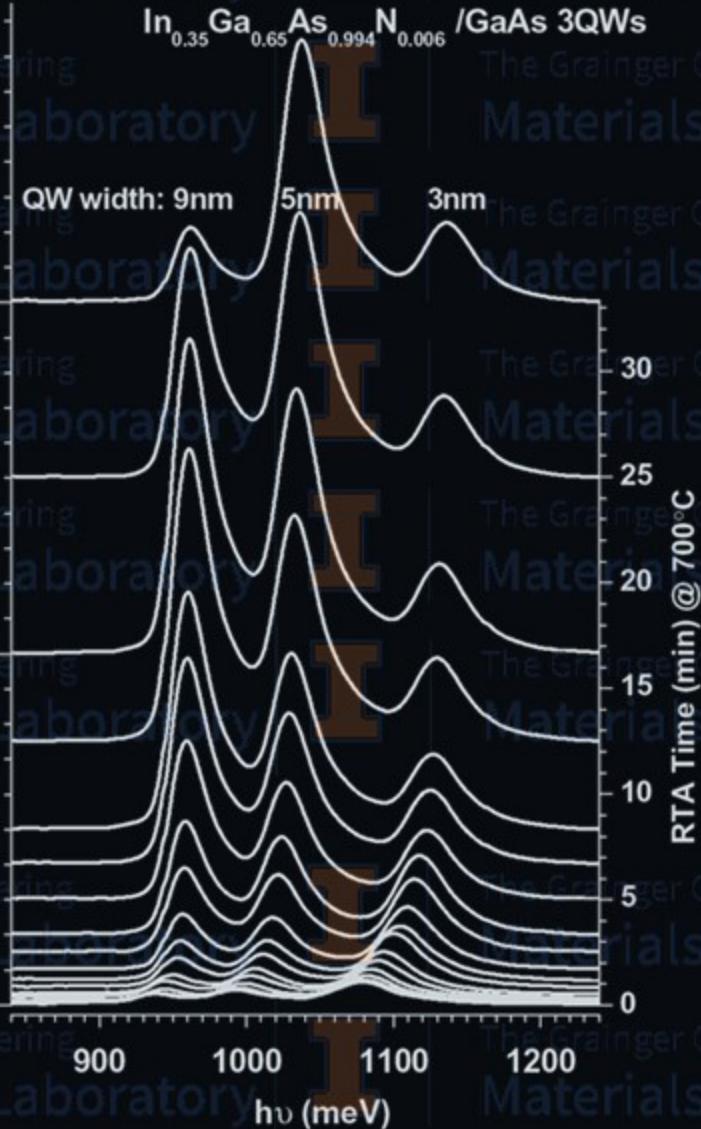
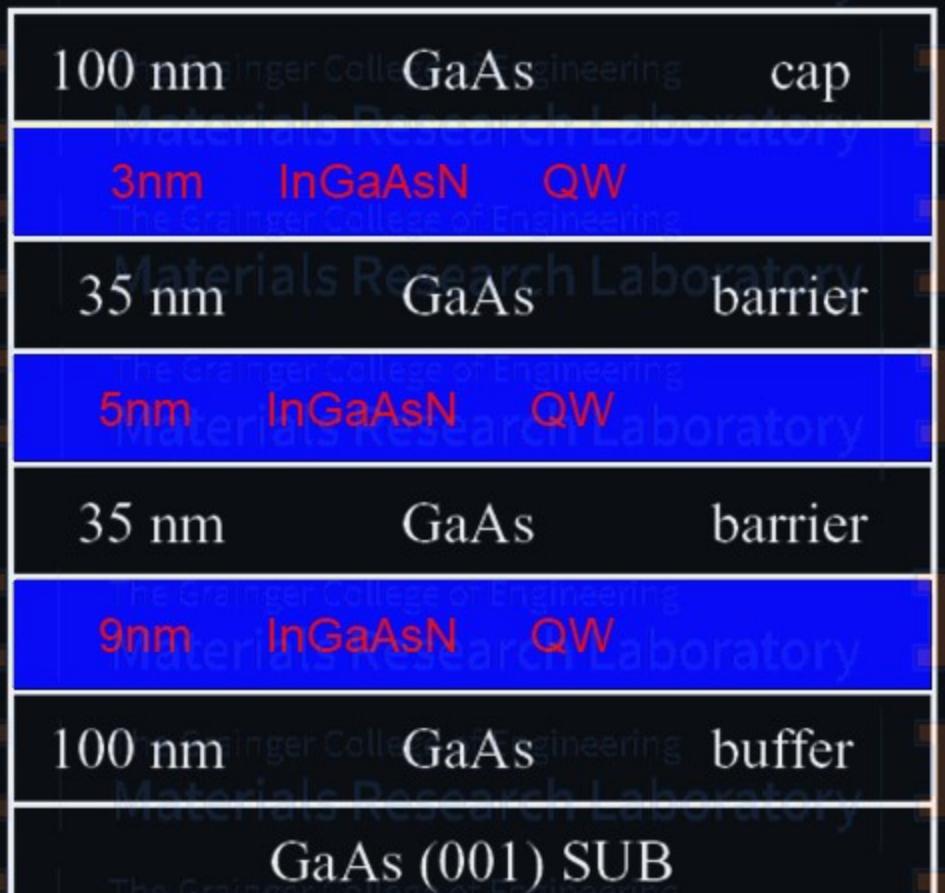
Ref.1 - Wetzel., *Appl. Phys. Lett.* **73**, 73 (1998).

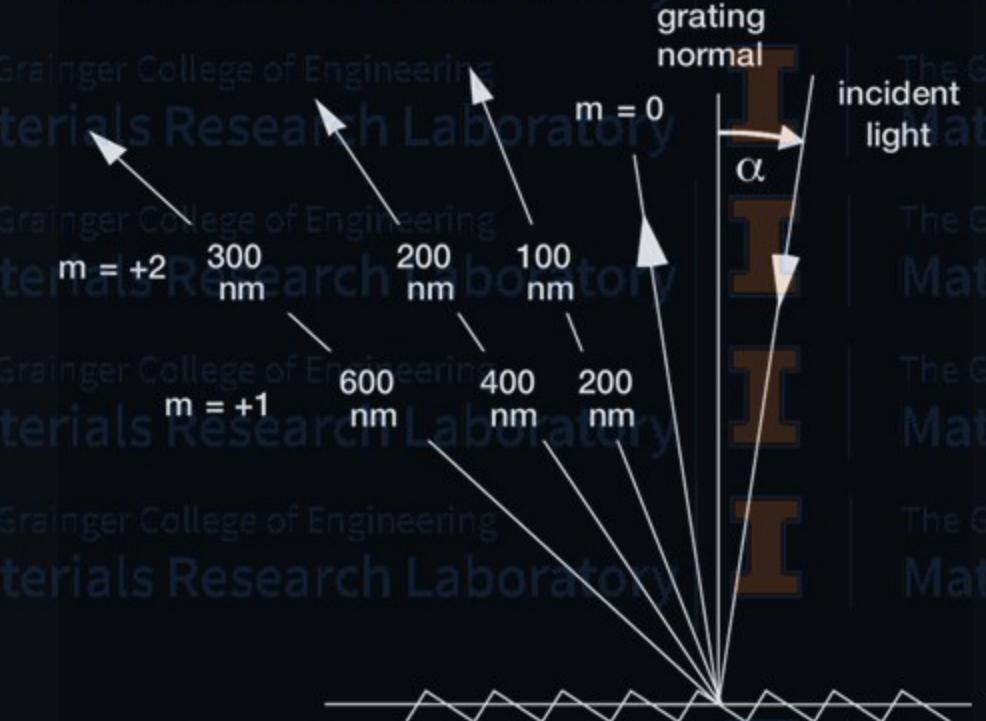
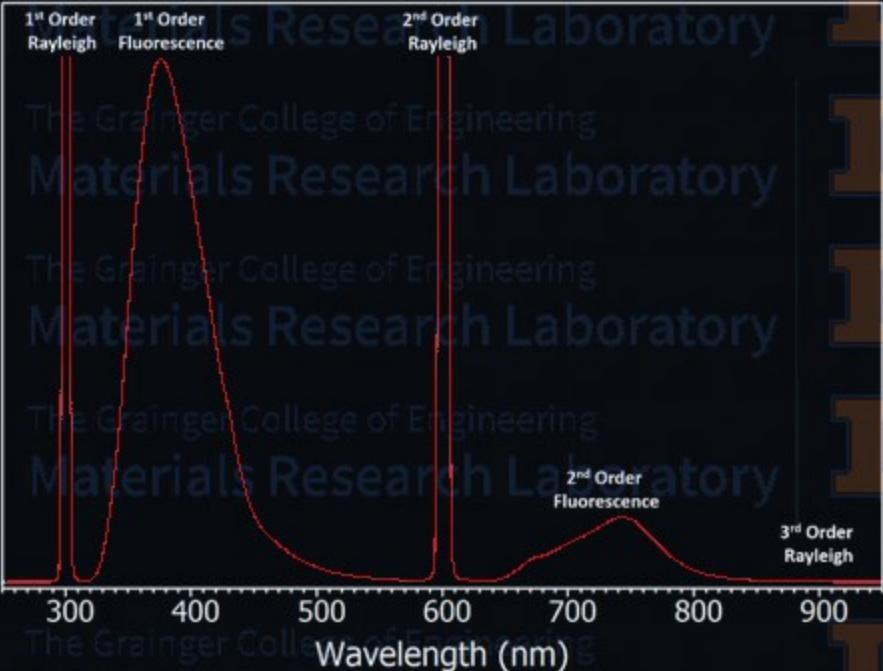
Ref.2 - V. Yu. Davydov., *Phys. Stat. Sol. (b)* **230**, R4 (2002).

Ref.3 - O'Donnell., *J. Phys. Condens. Matt.* **13**, 1994 (1998).

Extracted from *Phys. Stat. Sol. (b)* **234** (2002) 750

Width and quality of semiconductor quantum wells.

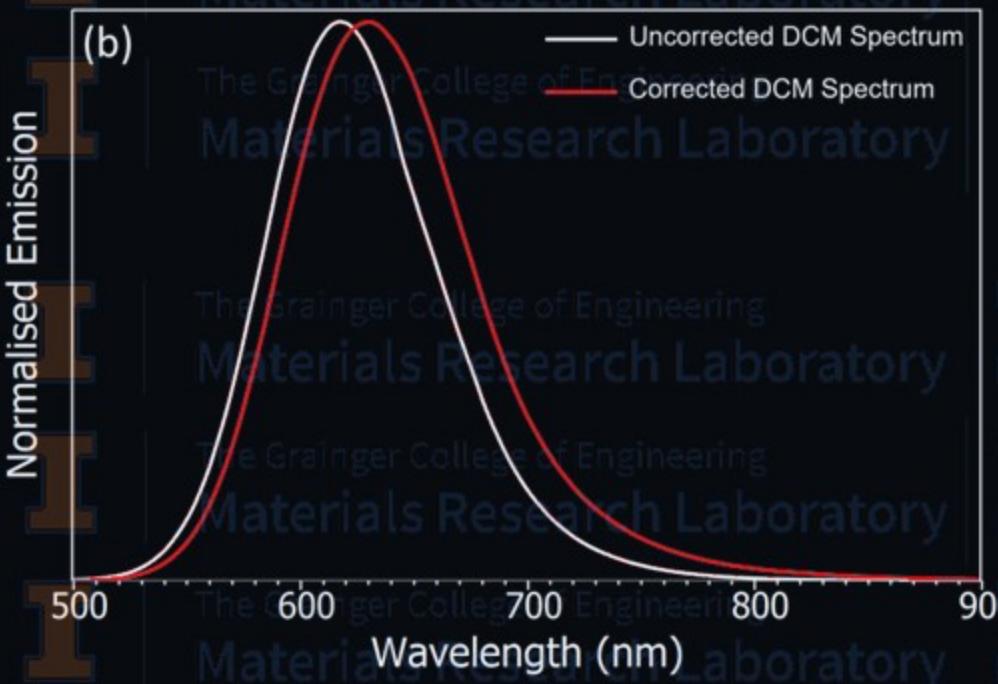
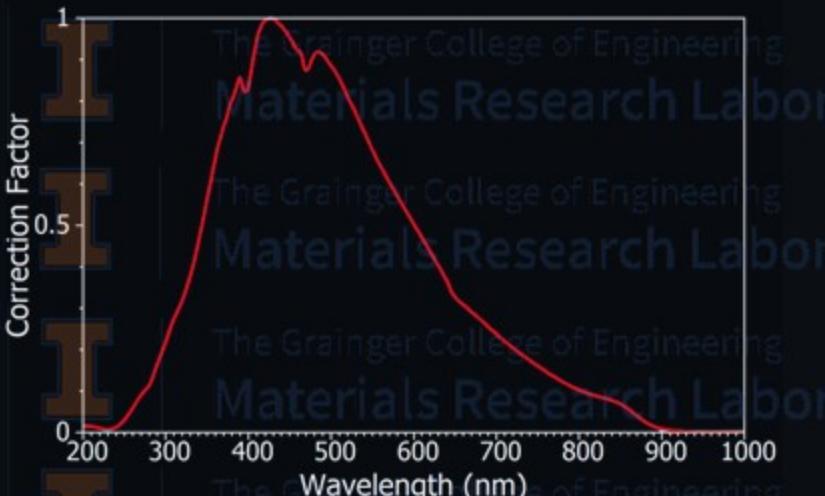
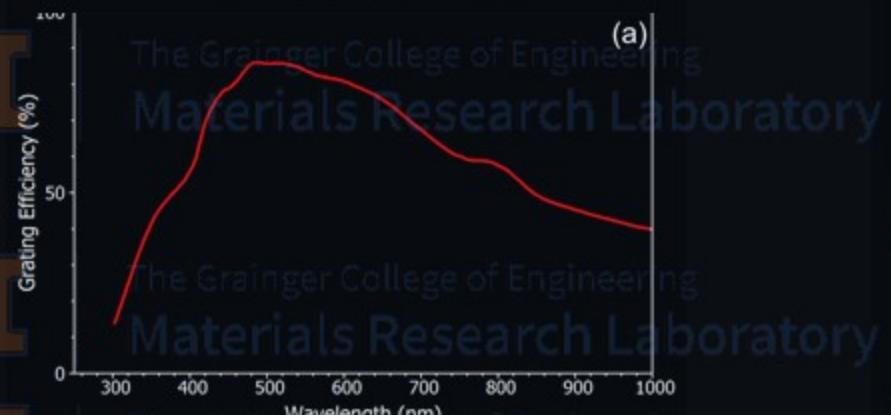
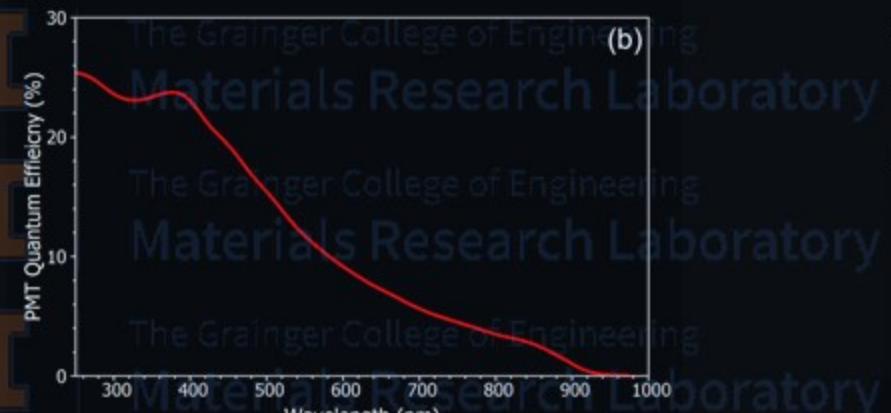




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www.newport.com

Pitfalls, artifacts, corrections ...



Non-ideal components
introduce spectral distortions

Strengths:

- Very little to none sample preparation.
- Non destructive technique.
- The very informative spectrum.

Limitations:

- Often requires low temperature.
- Data analysis may be complex.
- Many materials luminescence weakly.

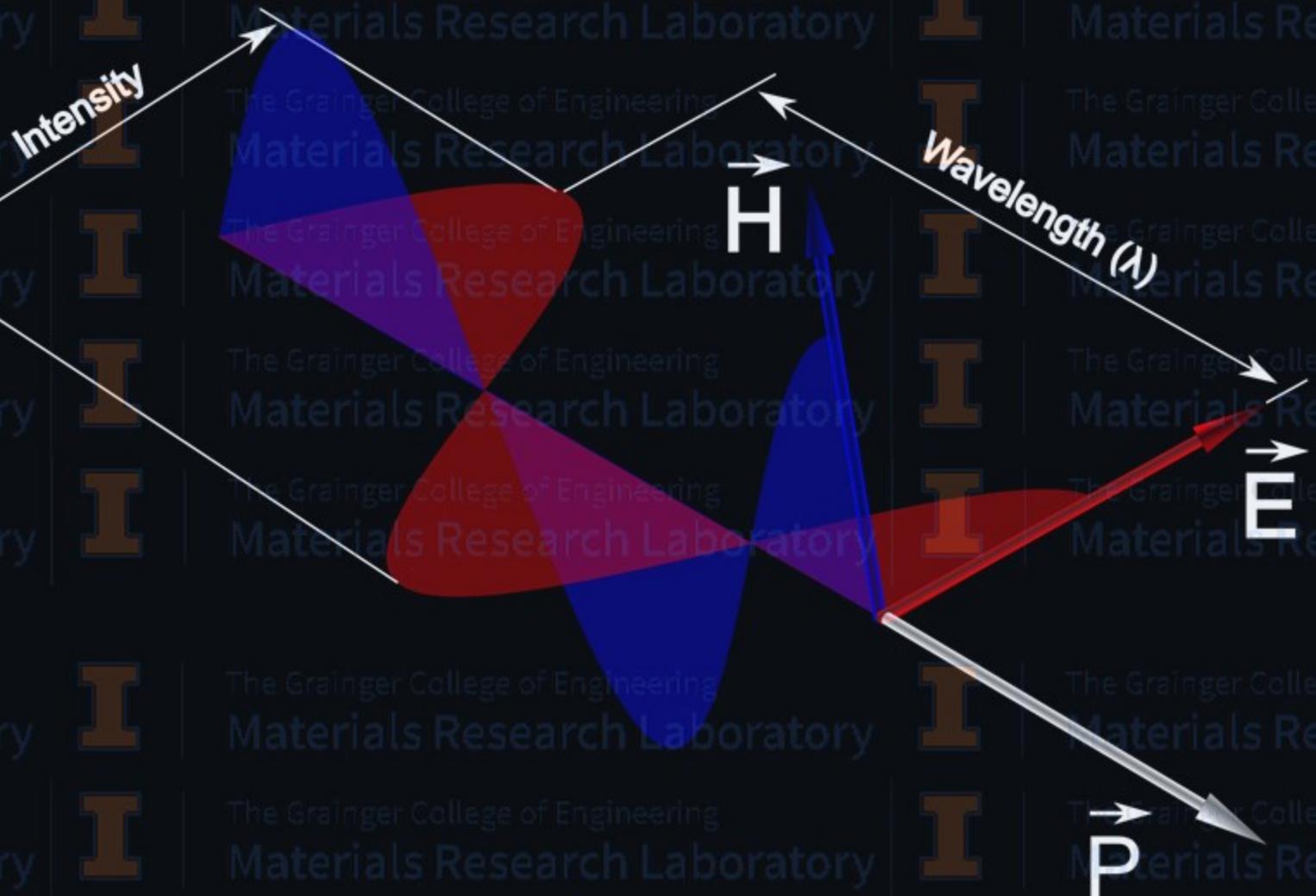
www.glofish.com



Complementary techniques:

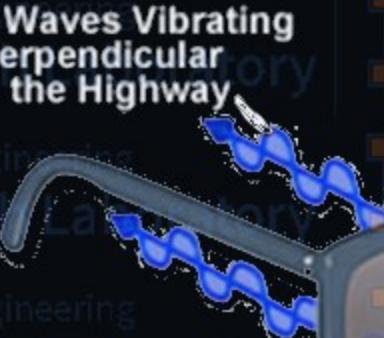
Ellipsometry, Modulation spectroscopies,
Spectrophotometry, Raman.

- Direction of propagation
- Electric field direction or polarization
- Photon energy or wavelength
- Intensity



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**Light Waves Vibrating
Perpendicular
to the Highway**



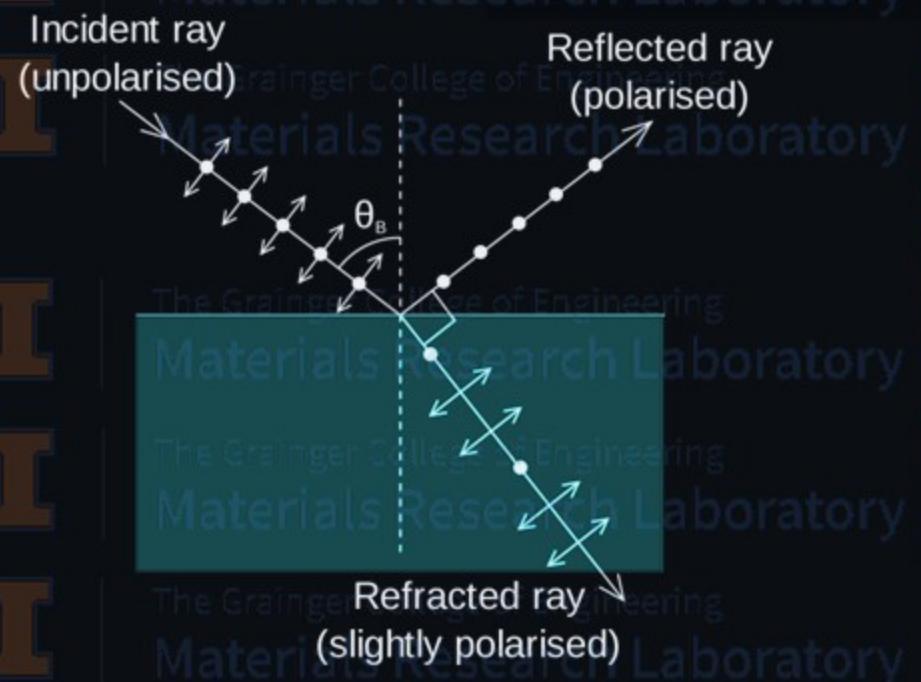
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**Light Waves
Vibrating
Parallel
to the Highway**

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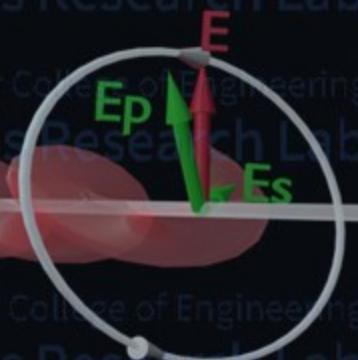
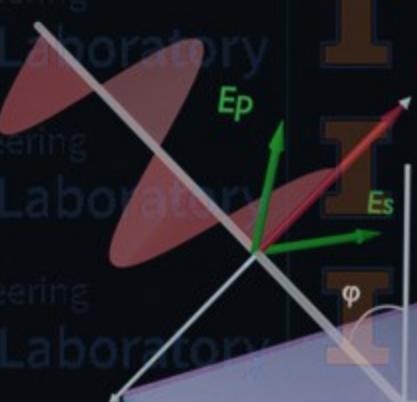
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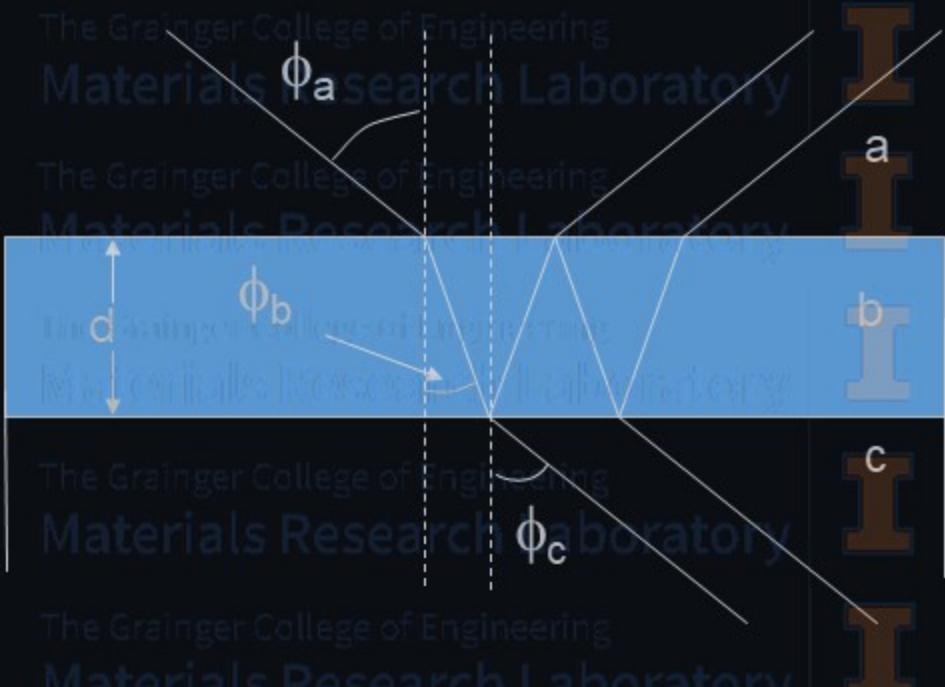
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Basic principle:

The reflected light emerges from the surface elliptically polarized, i.e. its p and s polarization components are generally different in phase and amplitude.



$$\tan(\Psi)e^{i\Delta} = \frac{\tilde{R}_p}{\tilde{R}_s}$$



$$\tilde{R}_{p,s} = \frac{E_{p,s}^r}{E_{p,s}} e^{i(\delta_{p,s}^r - \delta_{p,s}^i)}$$

$$\tilde{n} = n - ik \quad \tilde{n}_1 \sin \phi_1 = \tilde{n}_2 \sin \phi_2$$

$$\tilde{r}_{12}^{p,s} = \frac{\tilde{n}_{2,1} \cos \phi_1 - \tilde{n}_{1,2} \cos \phi_2}{\tilde{n}_{2,1} \cos \phi_1 + \tilde{n}_{1,2} \cos \phi_2}$$

$$\tilde{R}_{p,s} = \frac{\tilde{r}_{ab}^{p,s} + \tilde{r}_{bc}^{p,s} e^{-2i\beta}}{1 + \tilde{r}_{ab}^{p,s} \tilde{r}_{bc}^{p,s} e^{-2i\beta}}$$

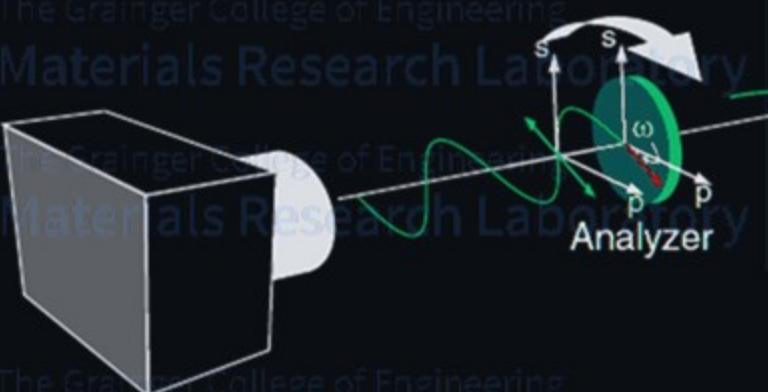
$$\beta = \frac{2\pi d}{\lambda} \tilde{n}_b \cos \phi_b$$

$$\tan(\Psi) e^{i\Delta} = \frac{\tilde{R}_p}{\tilde{R}_s} \Rightarrow \begin{cases} \tan(\Psi) = \frac{|\tilde{R}_p|}{|\tilde{R}_s|} \\ \Delta = \delta^r - \delta^i \end{cases}$$

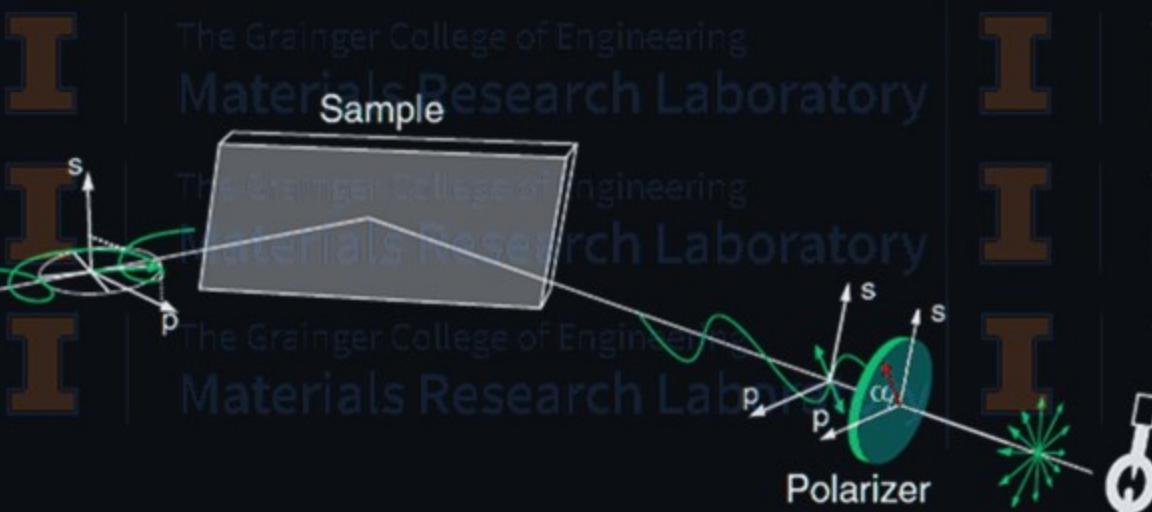
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What is measured:

The changes in the polarization state of light upon reflection from a mirror like surface.



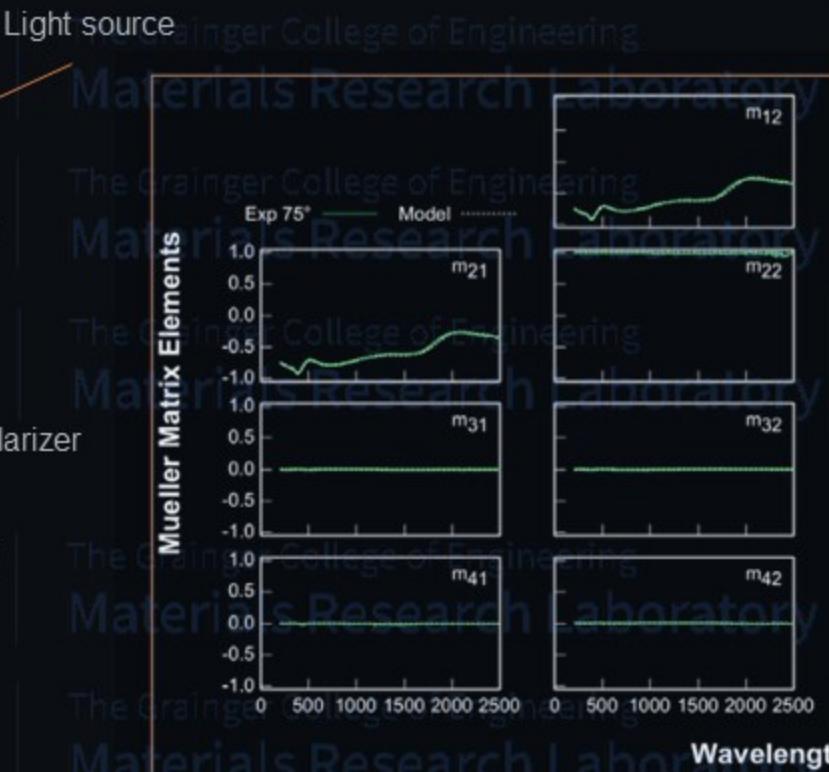
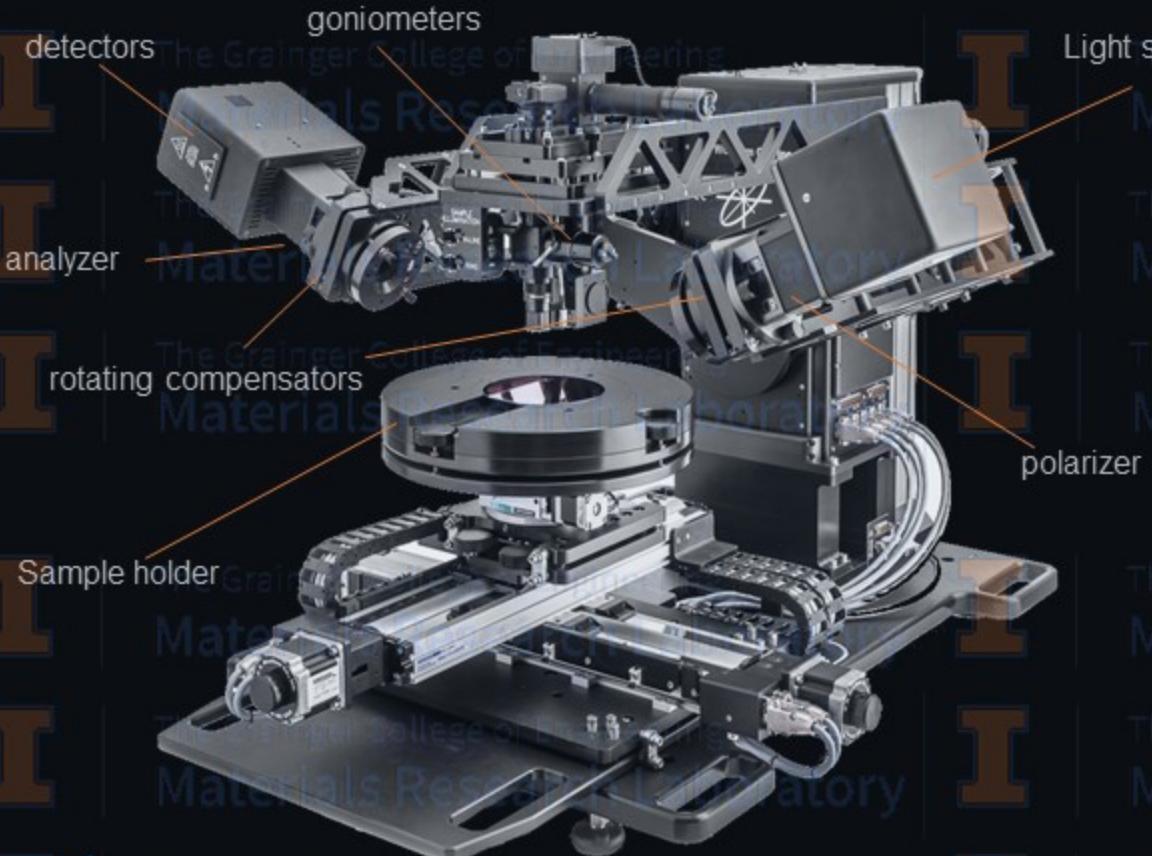
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Detector

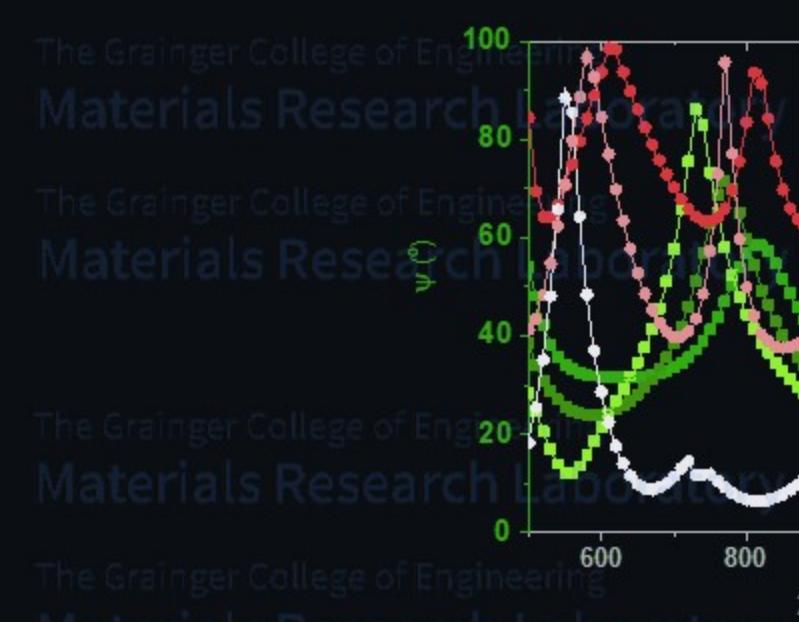
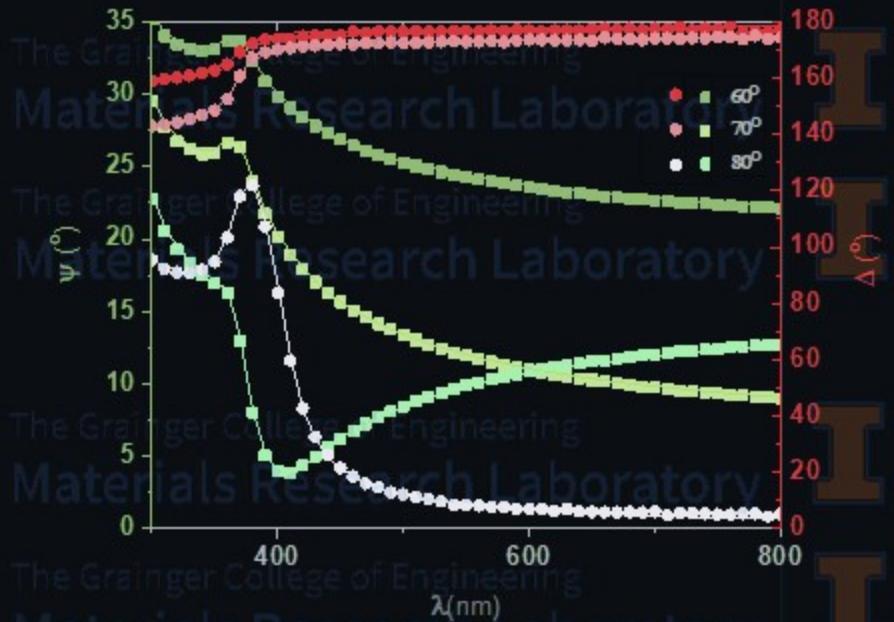


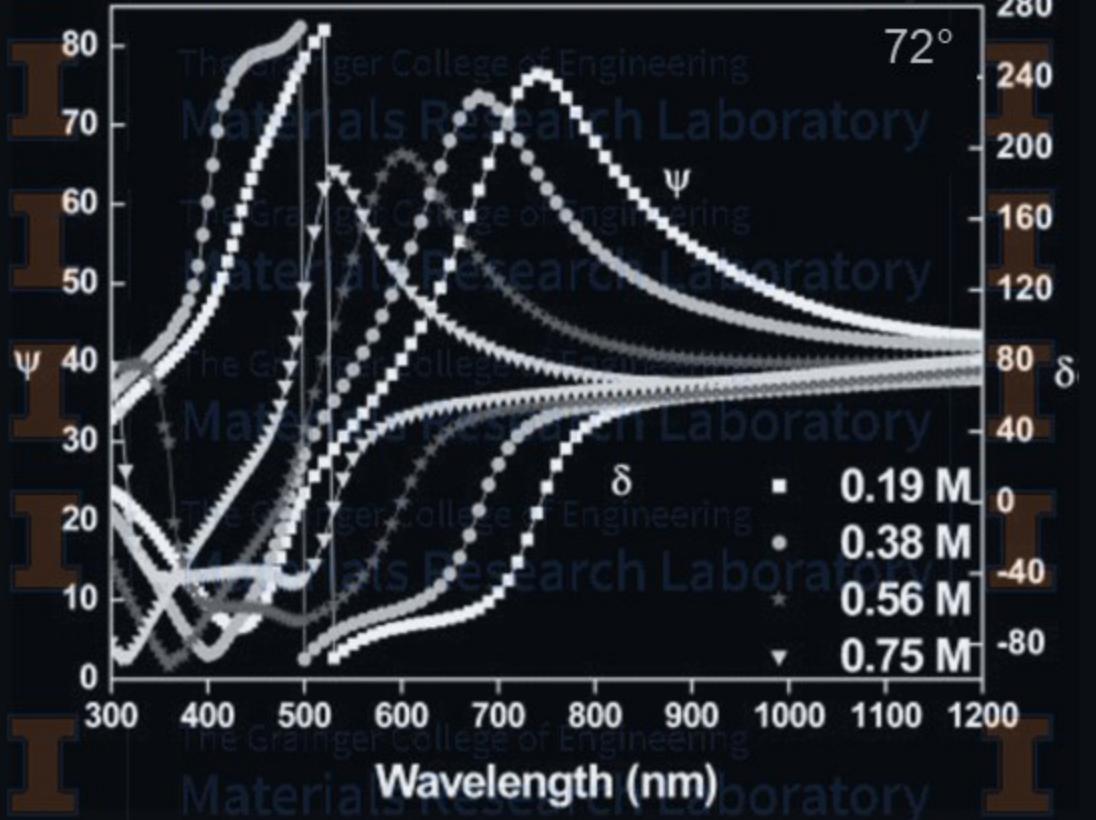
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What is measured:

The changes in the polarization state of light upon reflection from a mirror like surface.



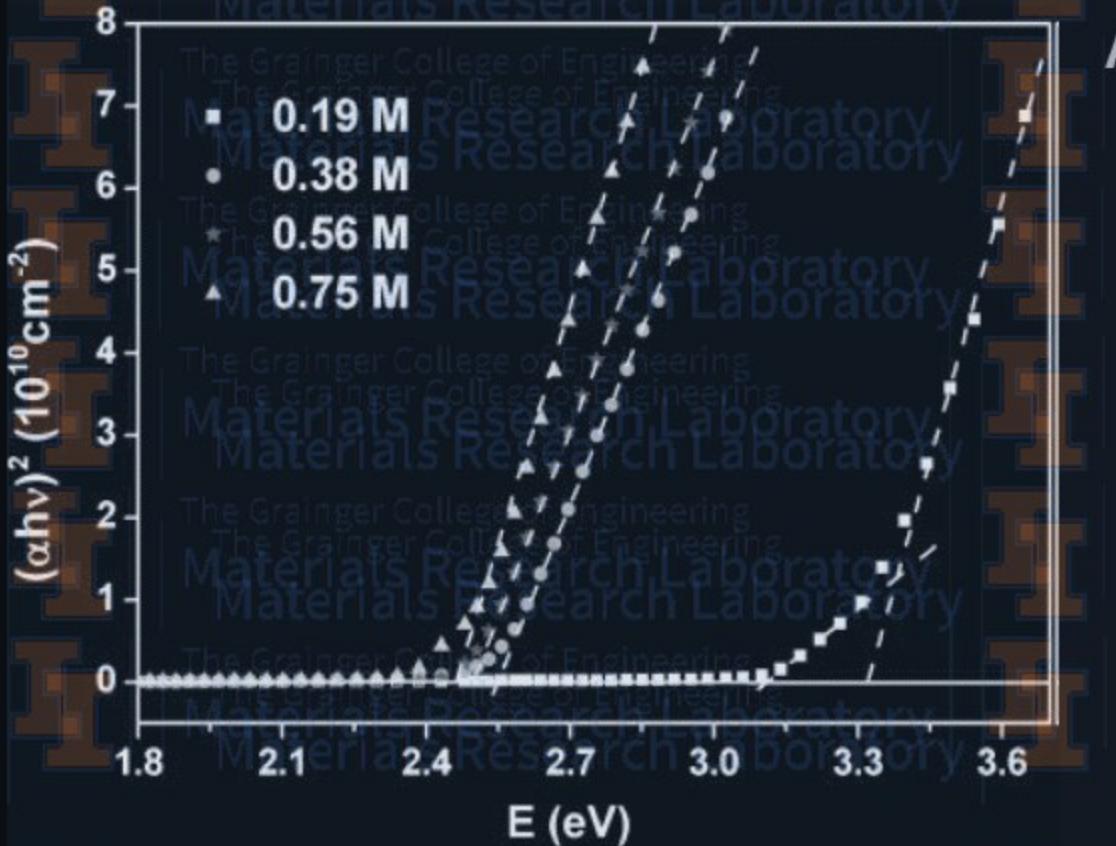




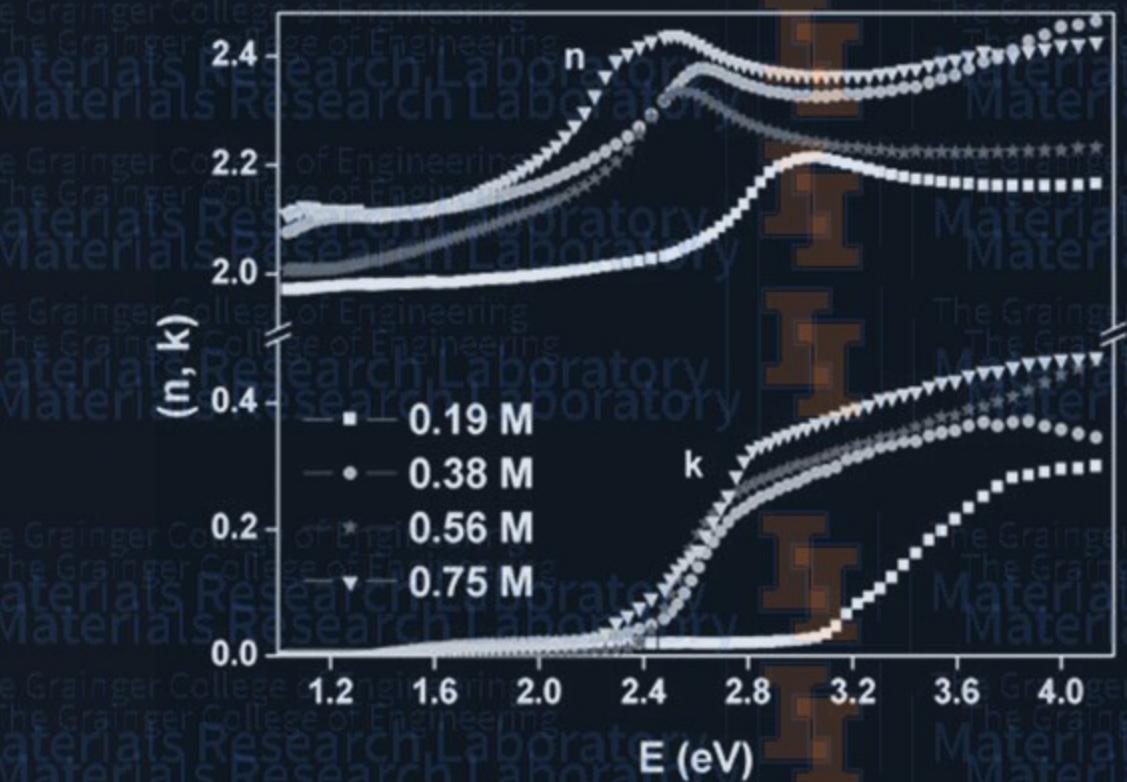
Ellipsometric $\Psi(\lambda)$ and $\Delta(\lambda)$ spectra of $\text{Cd}_{1-x}\text{Zn}_x\text{S}$ thin films deposited under the different concentration of ammonia: 0.19, 0.38, 0.56, and 0.75 M

Ellipsometry

Applications



- Composition
- Surface roughness
- Film thickness
- Band gap energy
- Optical constants (dielectric function)

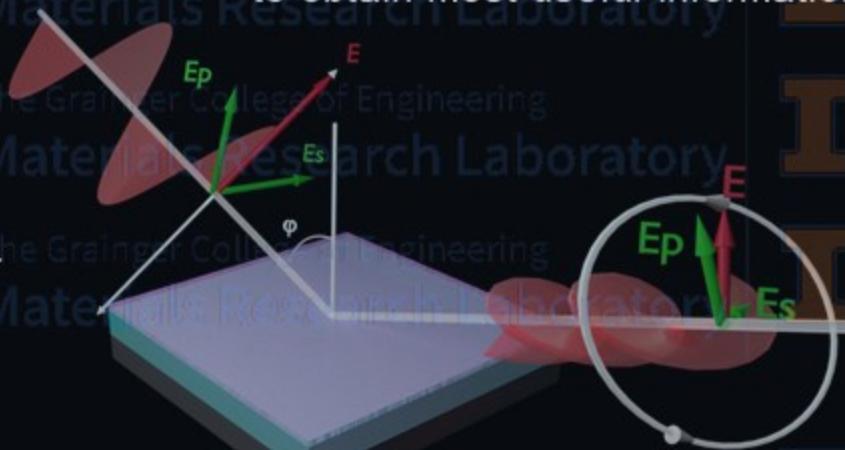


Strengths:

- Fast.
- Measures a ratio of two intensity values and a phase.
 - Highly accurate (even in low light levels).
 - No reference sample necessary.
 - Not susceptible to scatter, lamp or purge fluctuations.
 - Increased sensitivity, especially to ultrathin films (<10nm).
- Can be used in-situ.

Limitations:

- Flat and parallel surface and interfaces with measurable reflectivity.
- A realistic physical model of the sample is required to obtain most useful information.

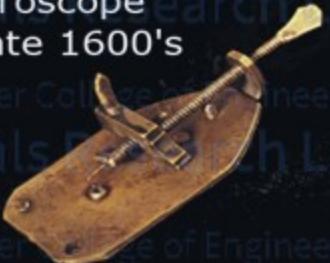


Complementary techniques:

PL, Modulation spectroscopies, X-Ray Photoelectron Spectroscopy, Secondary Ion Mass Spectroscopy, XRD, Hall effect.

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Von Leeuwenhoek
microscope
ca. late 1600's



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Optical microscopy
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Hand-held microscope ca. early 1700's



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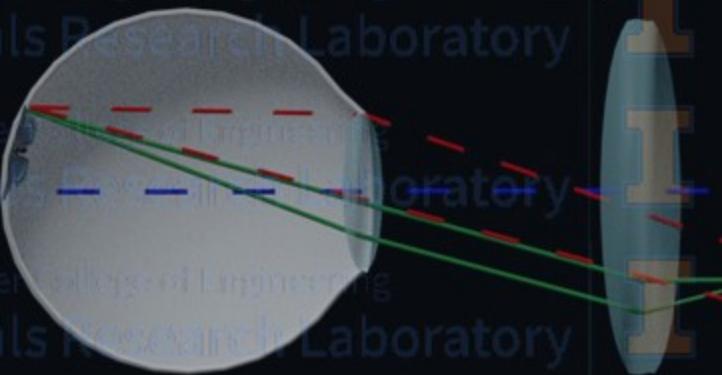
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Zeiss microscope
ca. 1930



Modern scientific microscope

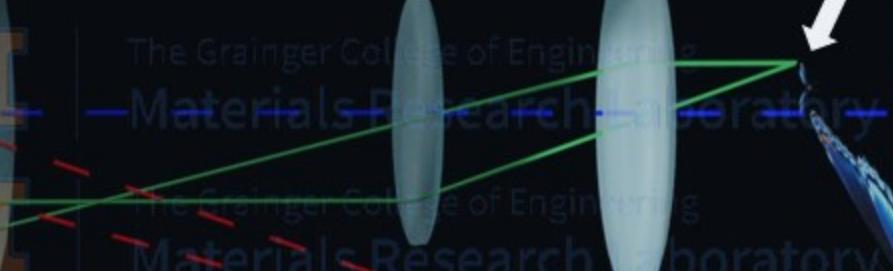
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Eyepiece

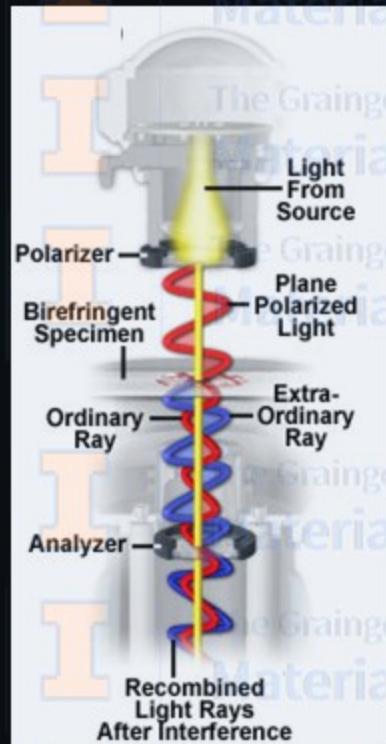
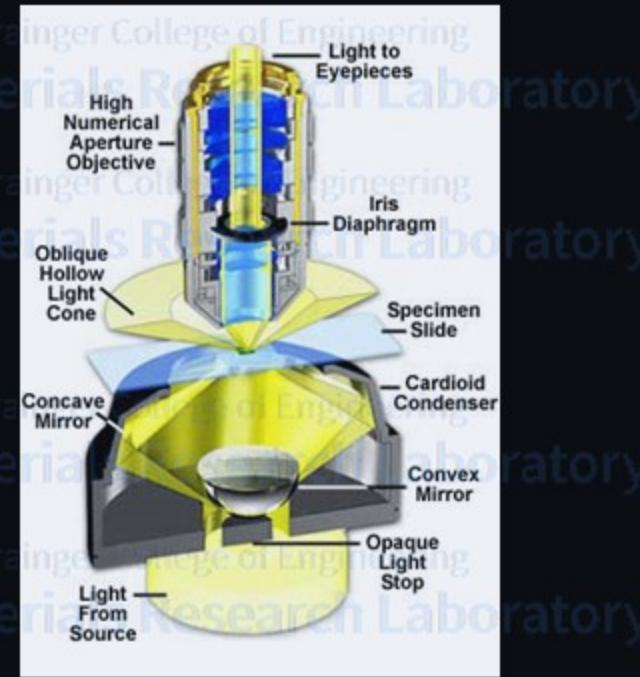
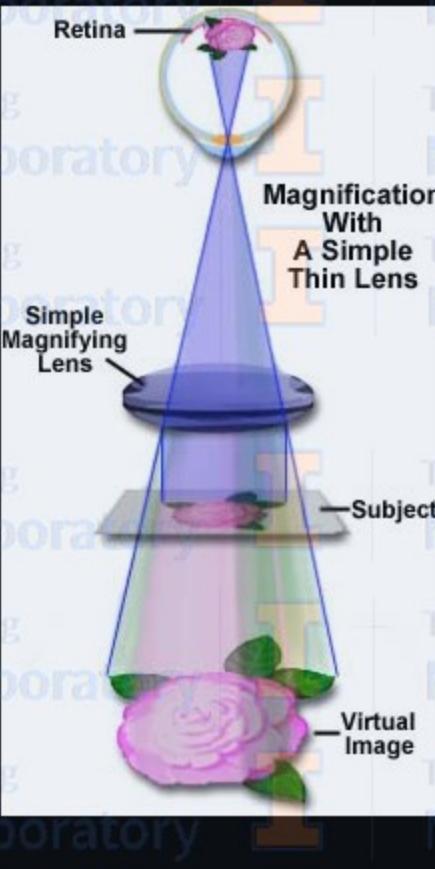
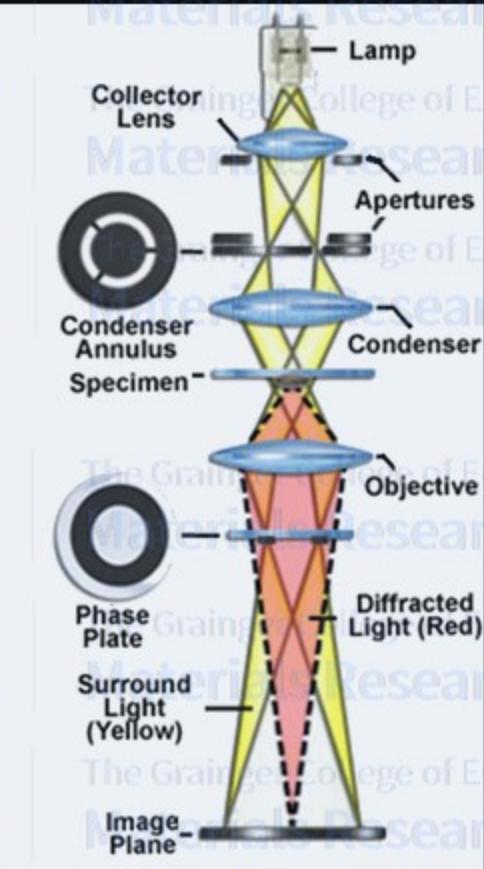
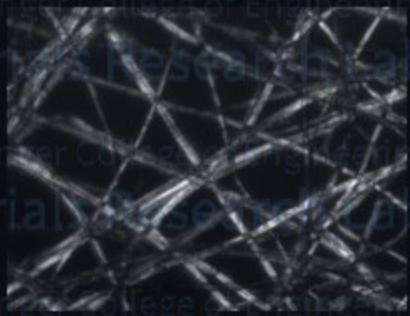
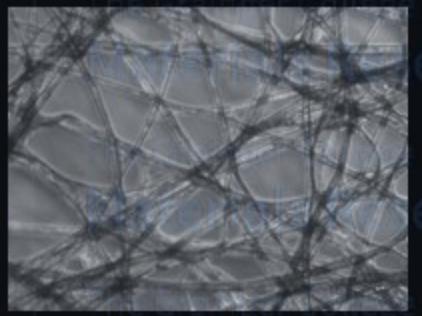


Eye

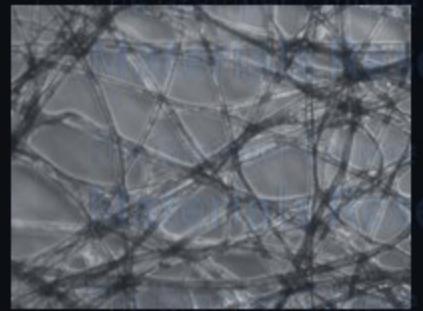
Tube lens



Perceived image



Phase contrast

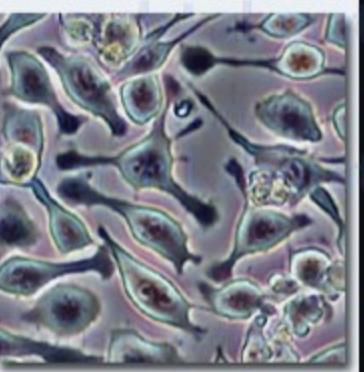


Bright field



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Living Cells in Brightfield and Phase Contrast



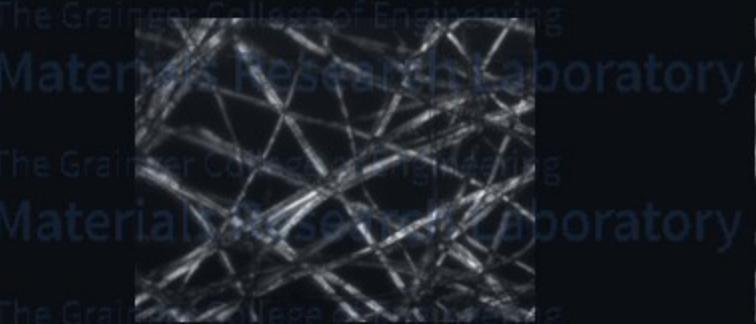
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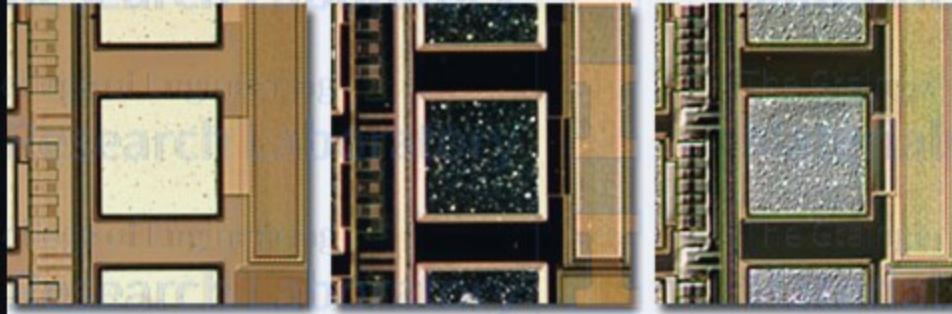
Optical microscopy

Dark field



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Integrated Circuit in Brightfield, Darkfield, and DIC with Reflected Light



Phyllite Thin Section in Polarized Light



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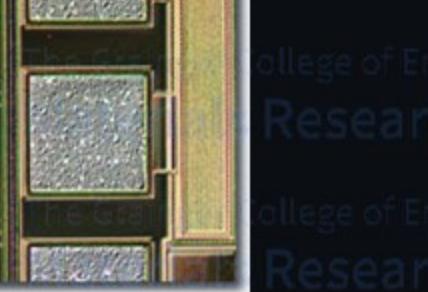
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Polarizing



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• But how small a thing can we see?

• But how small a thing can we see?



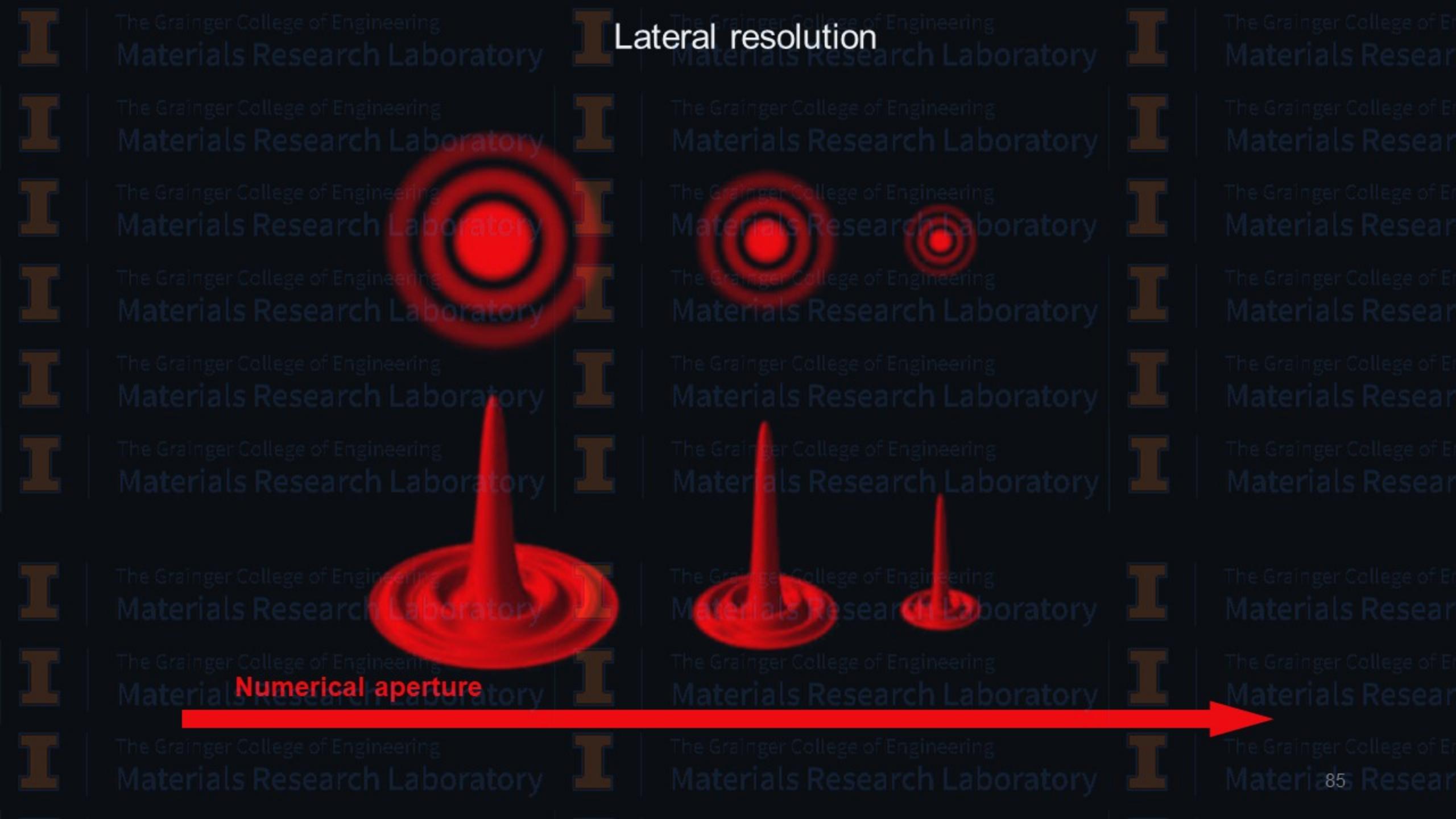
bing

Diameter of
Airy Disc

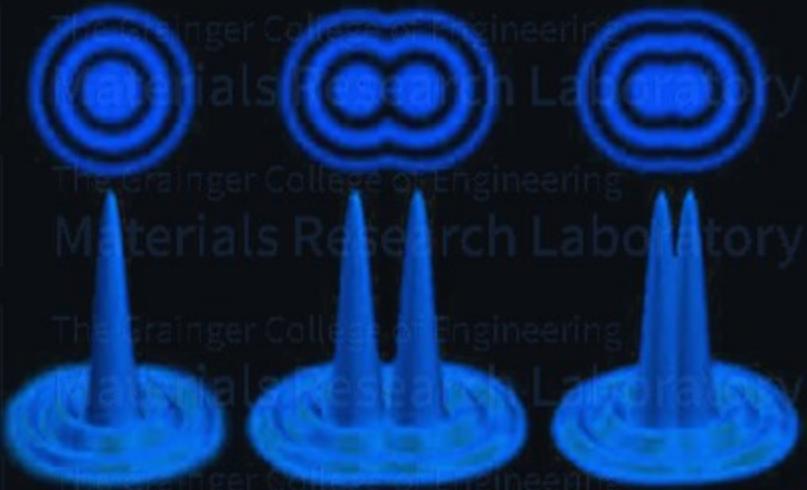


The Airy pattern



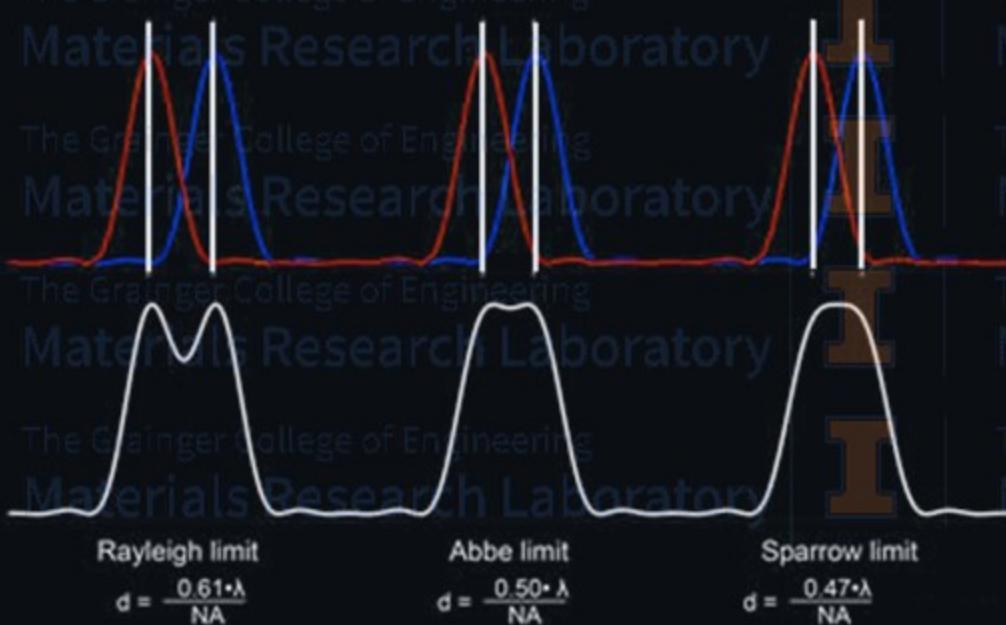


Airy Discs



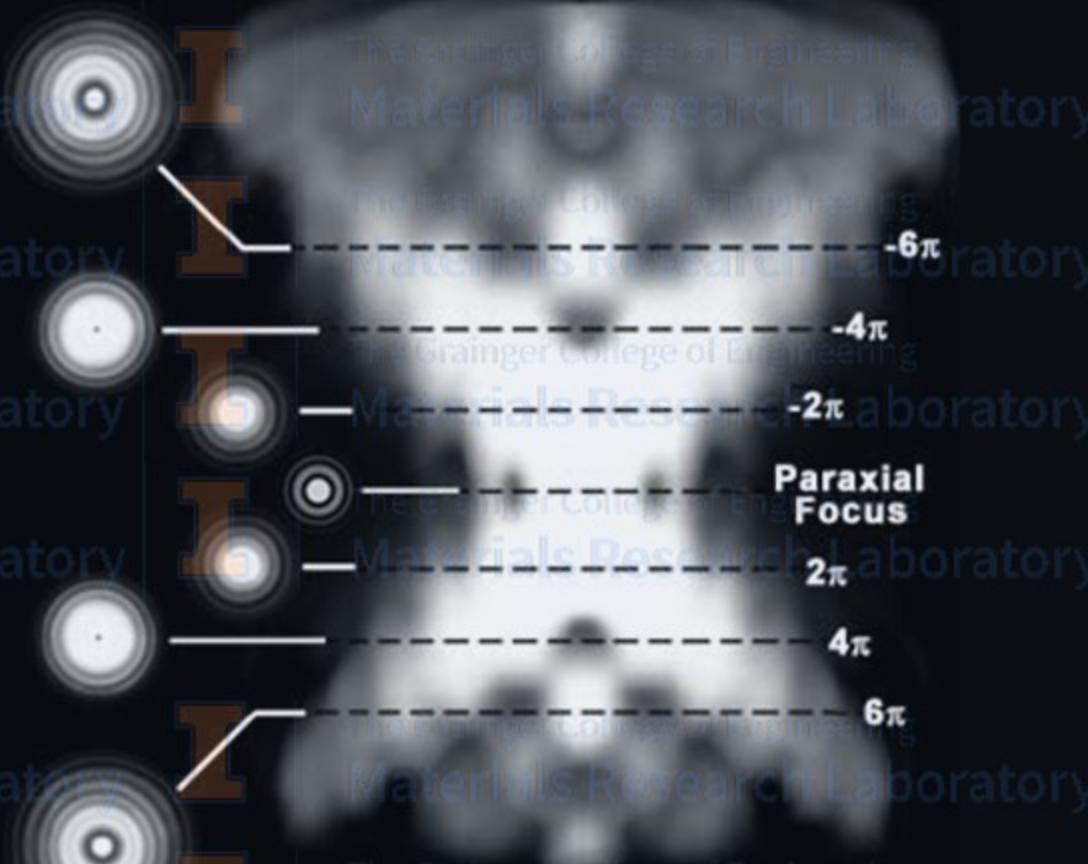
$$d \approx \frac{0.61\lambda}{NA}$$

Rayleigh criterion



Abbé criterion

Figure

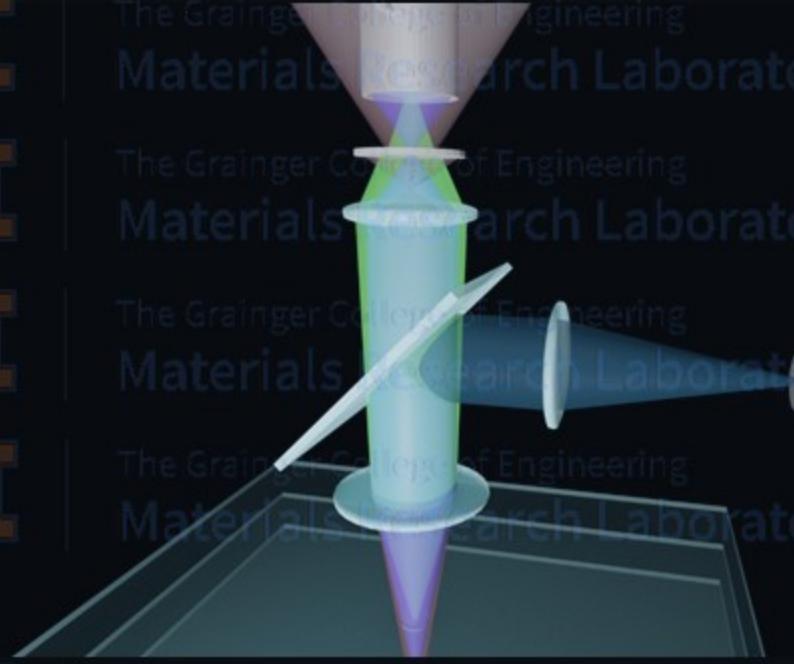




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- Increased contrast => 200:1.
- Slightly increased in plane resolution (1.5 x)
- Significantly increased resolution along the optical axis.
- Scanning image formation.

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Confocal microscopy

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Increased contrast => 200:1
Slightly increased in plane resolution (1.5 x)
Significantly increased resolution along the o
Scanning image formation.

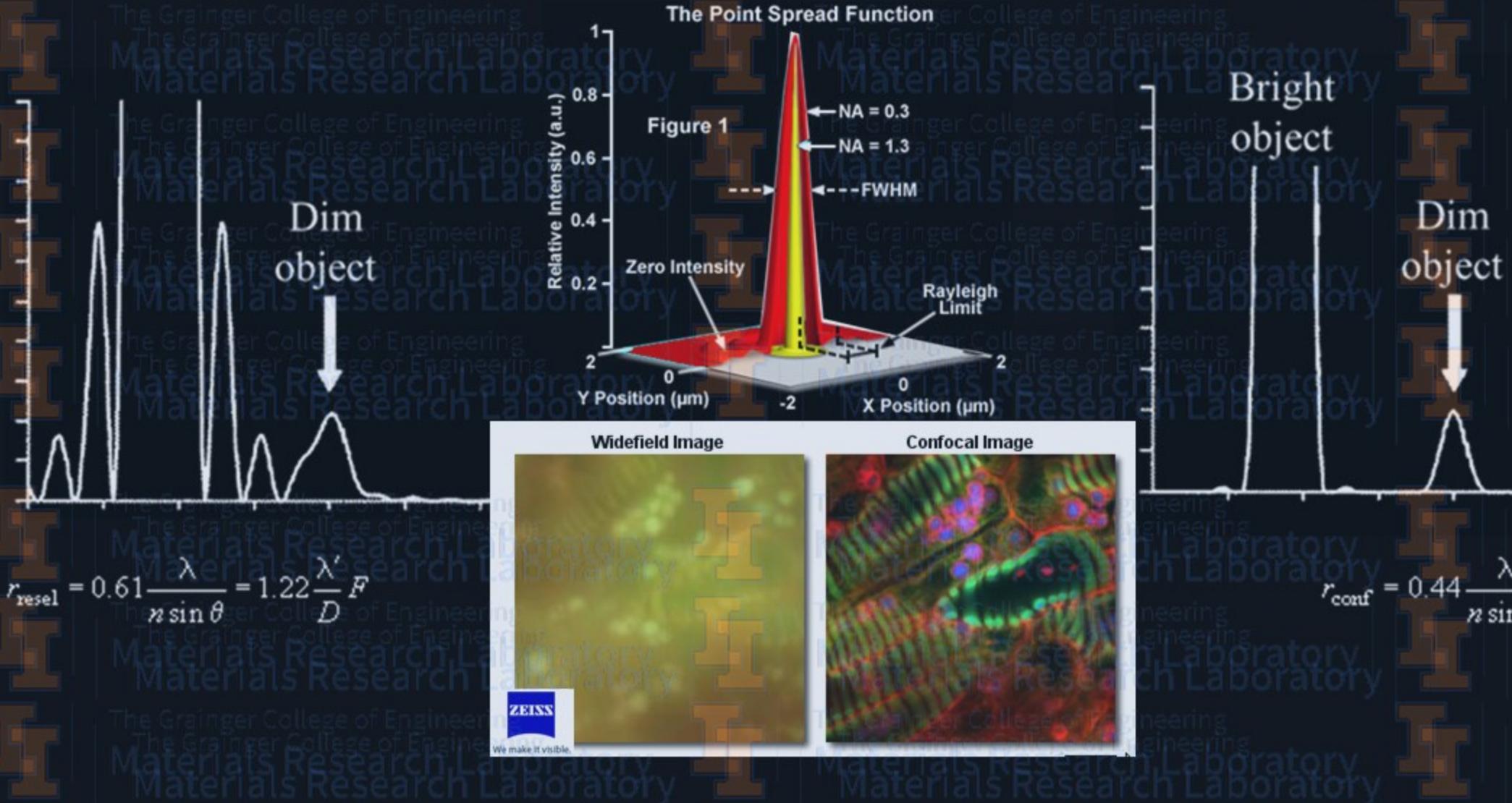


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Confocal microscopy

The relation of the first ring maximum amplitude to the amplitude in the center is 2% in case of conventional point spreading function (PSF) in a focal plane, while in case of a confocal microscope this relation is 0.04%.

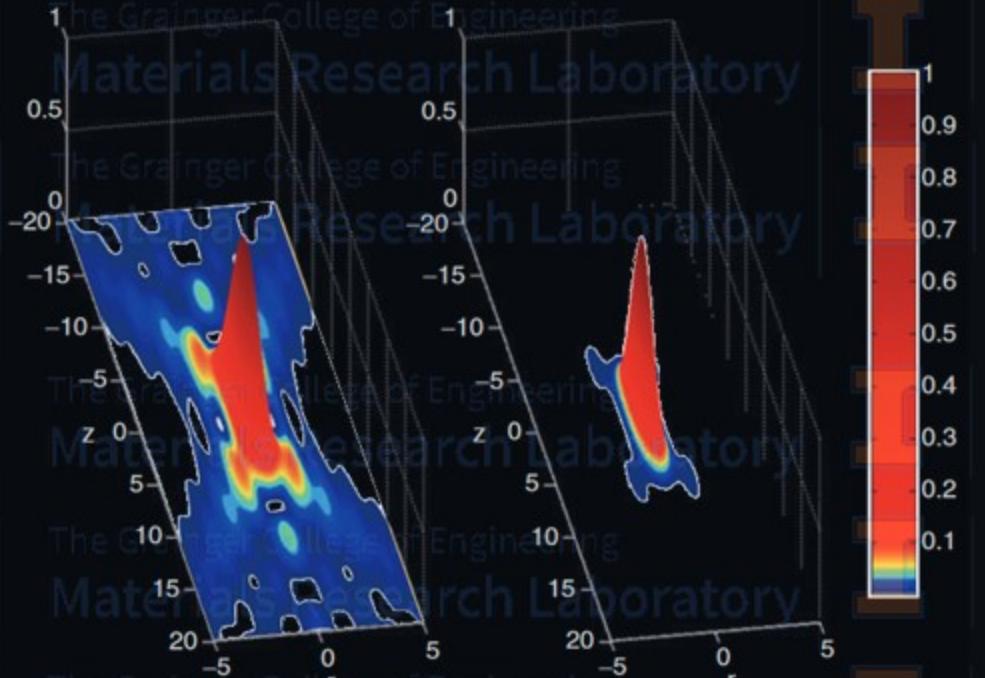


Confocal microscopy combined with spectroscopy

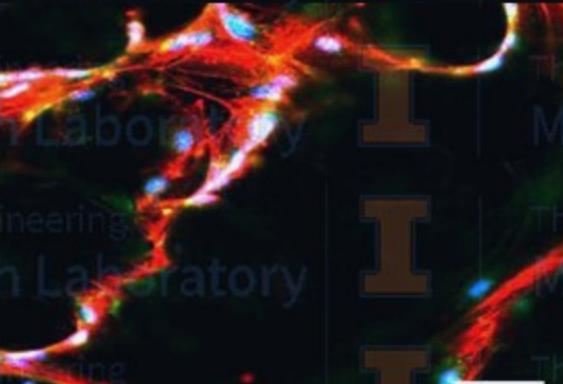
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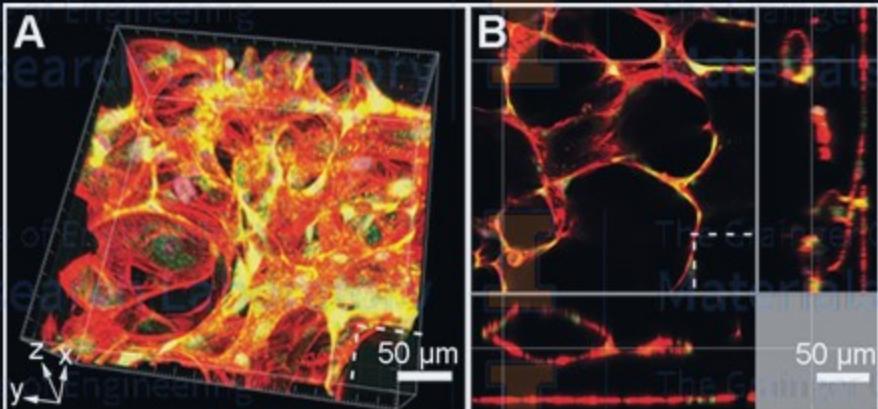
Widefield microscope Confocal microscope
PSF PSF



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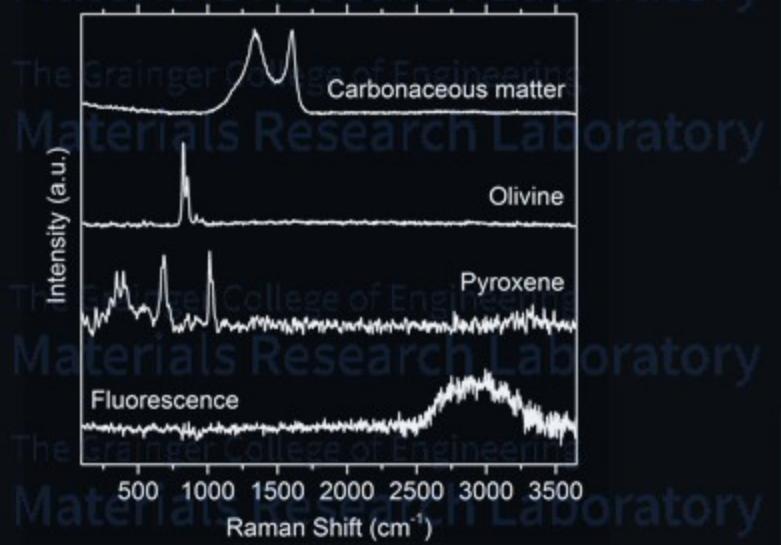
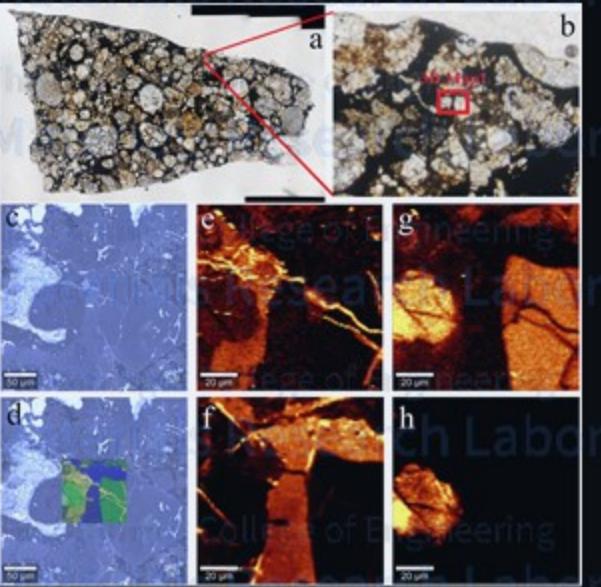
Confocal microscopy reconstruction of a 3D capillary bed



PLOS ONE 7(12): e50582 (2012)

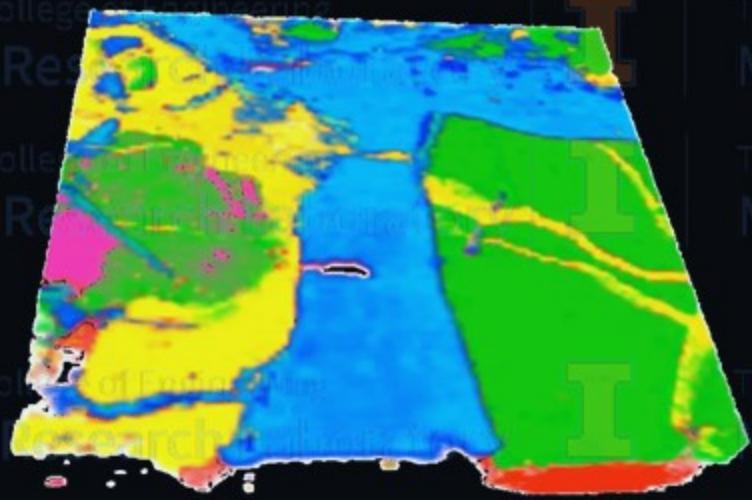
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Confocal microscopy combined with spectroscopy



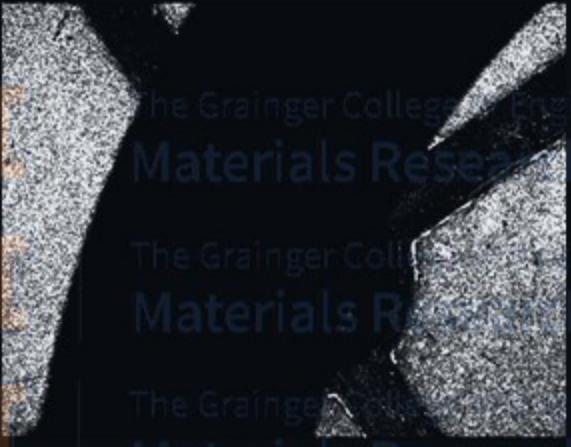
Chemical composition
Component identification
Components distribution

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Confocal microscopy z-stack

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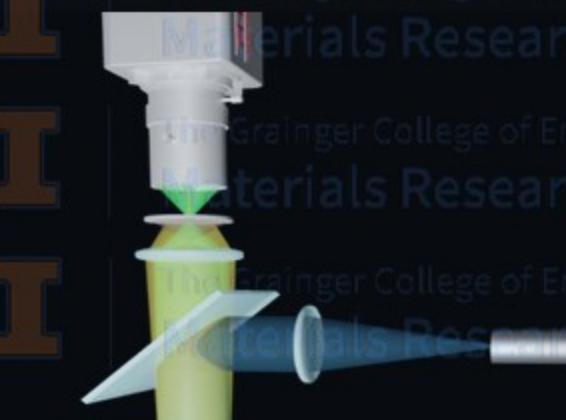
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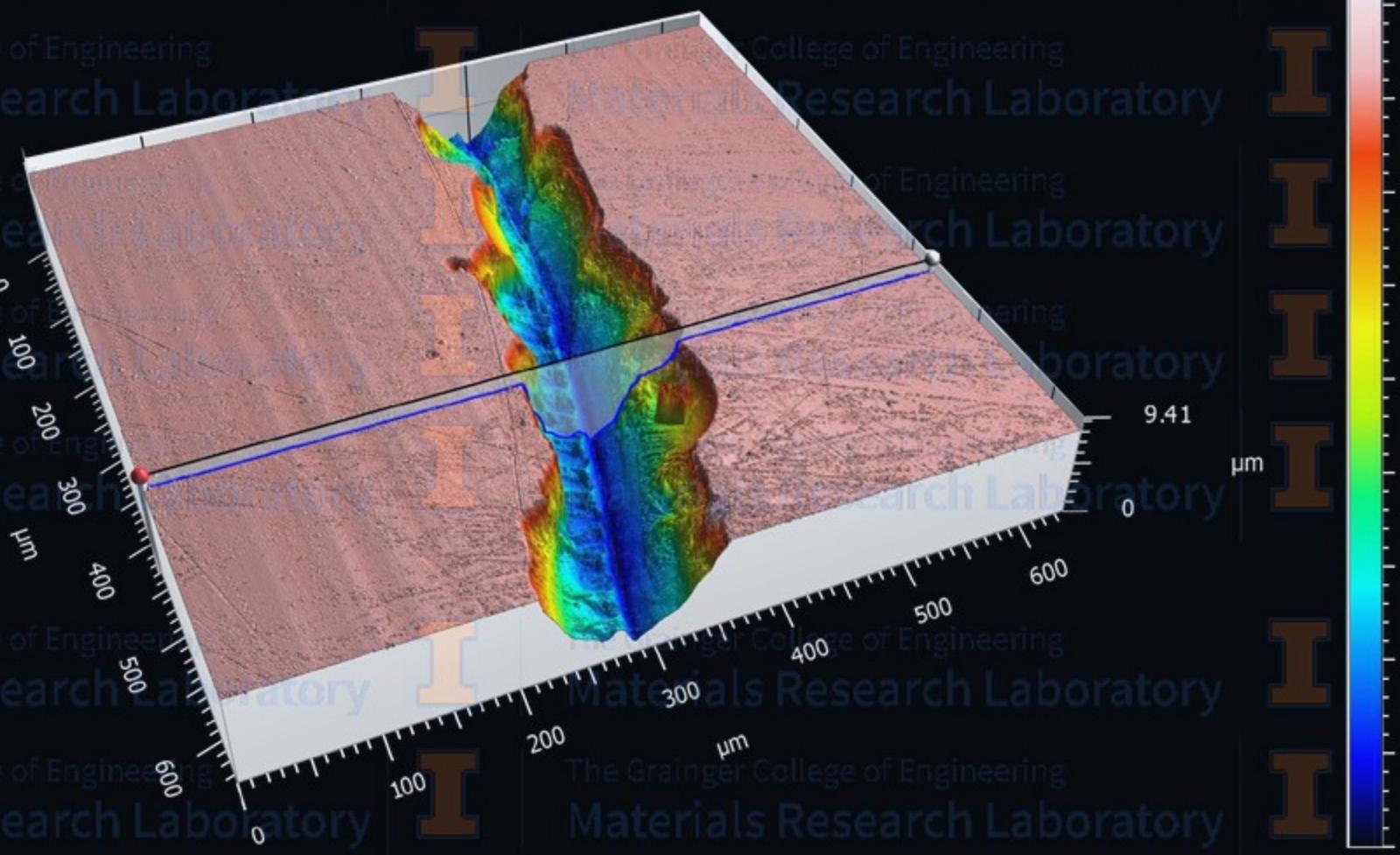


Confocal microscopy for measuring topography

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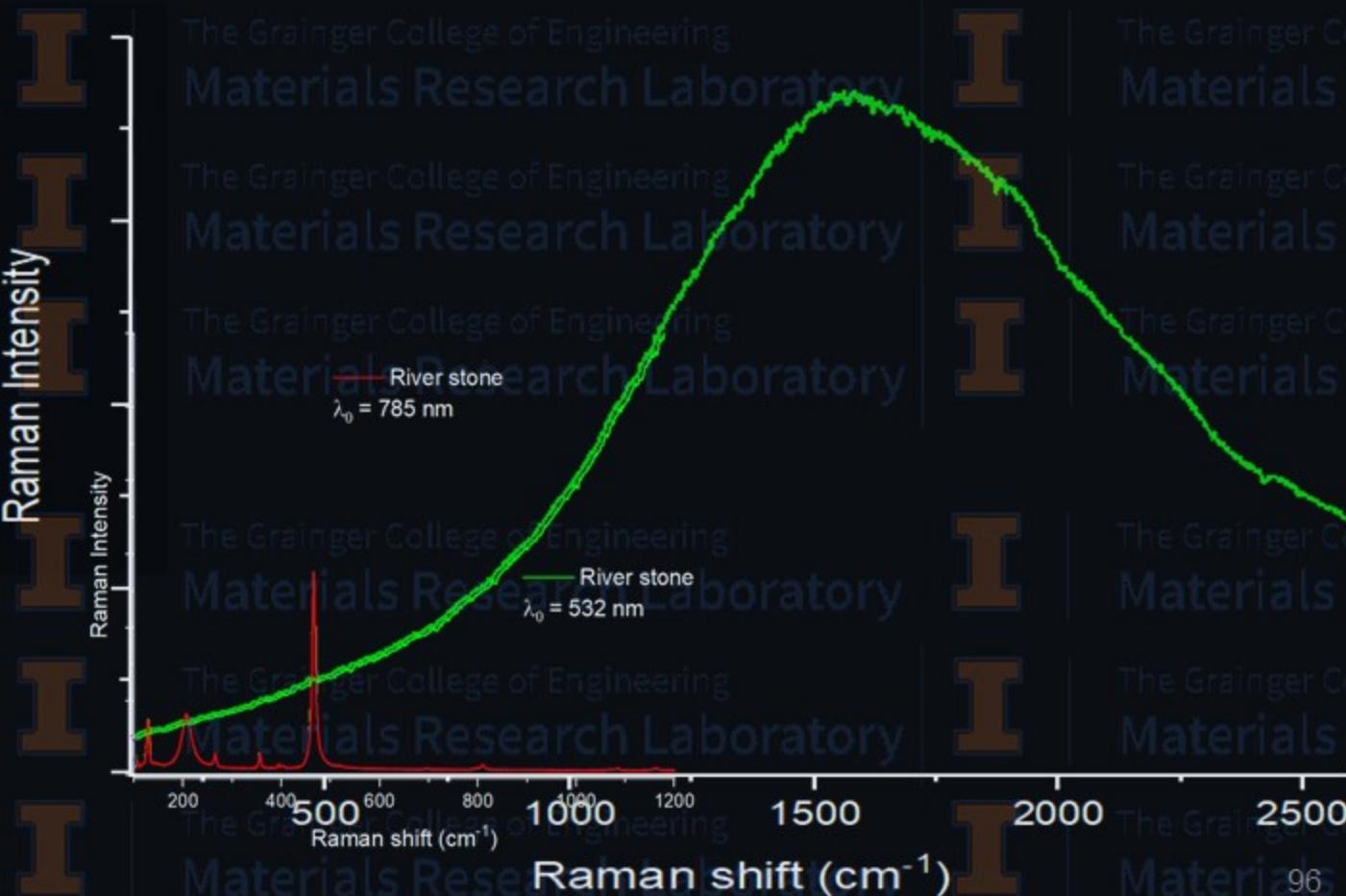
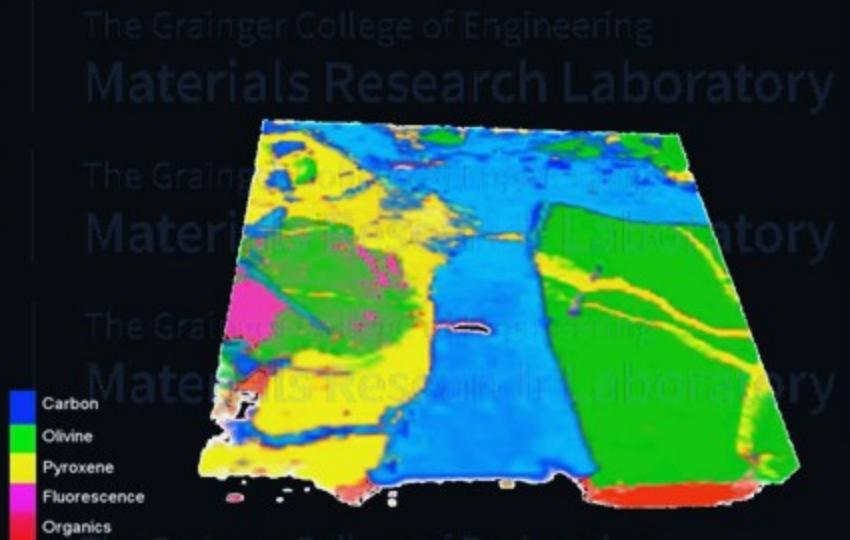
μm

9
8
7
6
5
4
3
2
1
0

95

Primary Limitations:

- Expensive apparatus (for high spectral/spatial resolution).
 - Weak signal, compared to fluorescence.
 - Limited spatial resolution (diffraction limited).



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The More Time Approach

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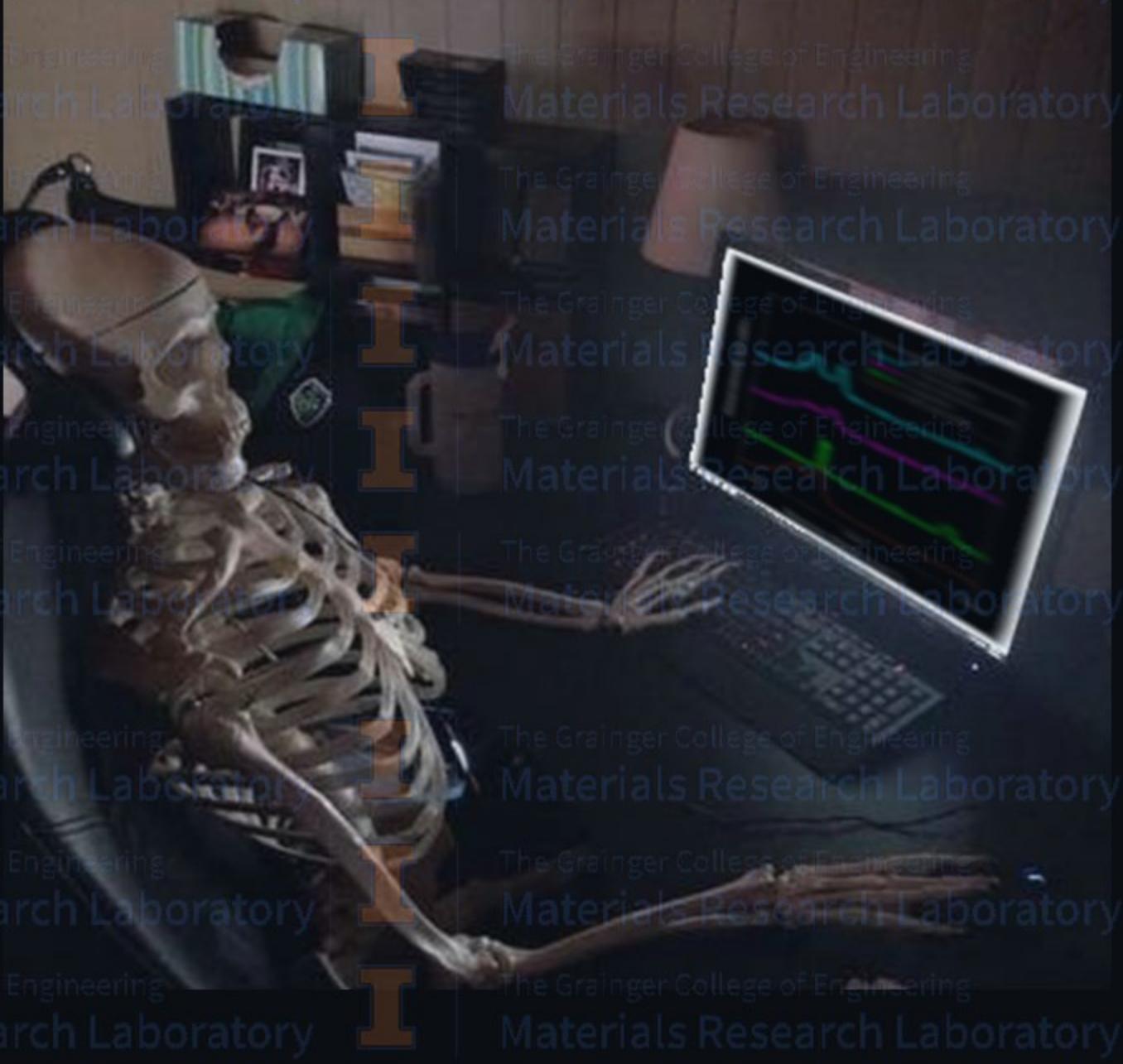
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Dielectric

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Metal

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Plasmons can be driven by photons at resonance to build large standing wave electric fields.

That leads to a strong enhancement of Raman scattering, proportional to fourth power of the E field strength.

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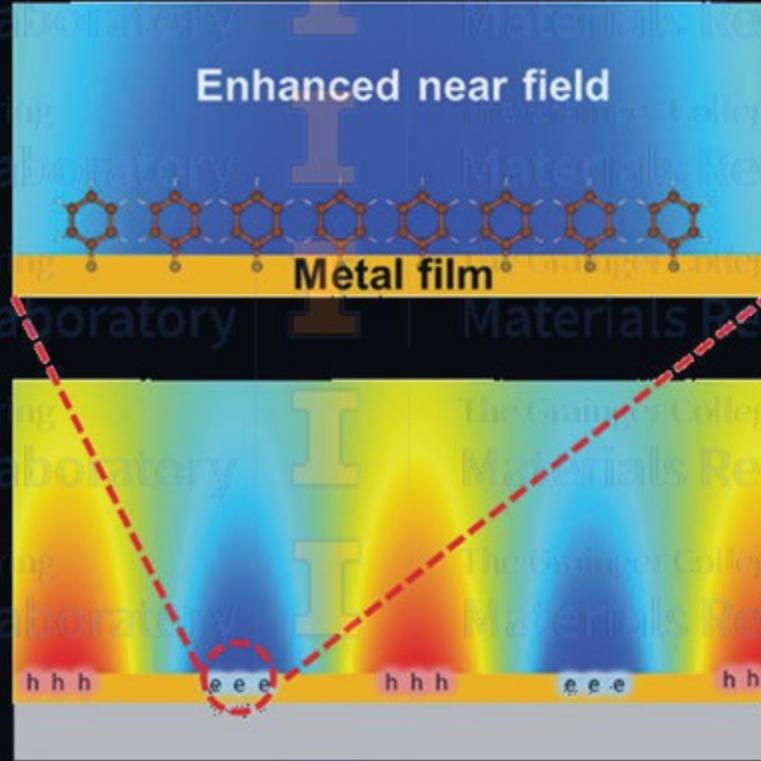
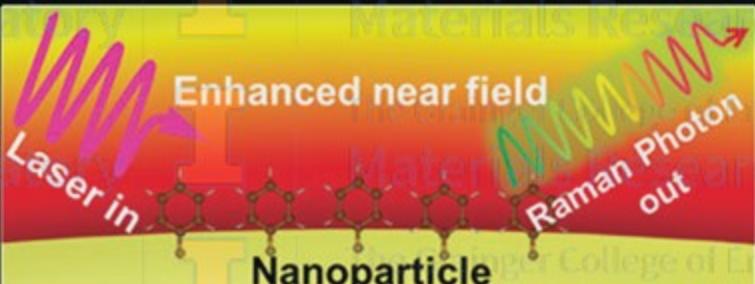
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Surface Enhanced Raman Spectroscopy (SERS)

Typically achieved with corrugated gold/silver surface or gold/silver nanoparticles with molecules of interest attached.

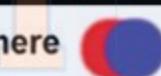
Capable of boosting Raman signal up to **14 Orders of Magnitude** or more! *Science* **275**, 1102 (1997)



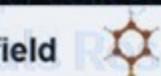
Chem. Rev. **117**, 5002, (2017)



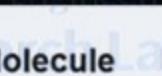
Metal sphere



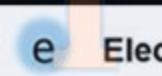
Electric field



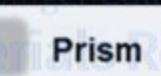
Molecule



Holes

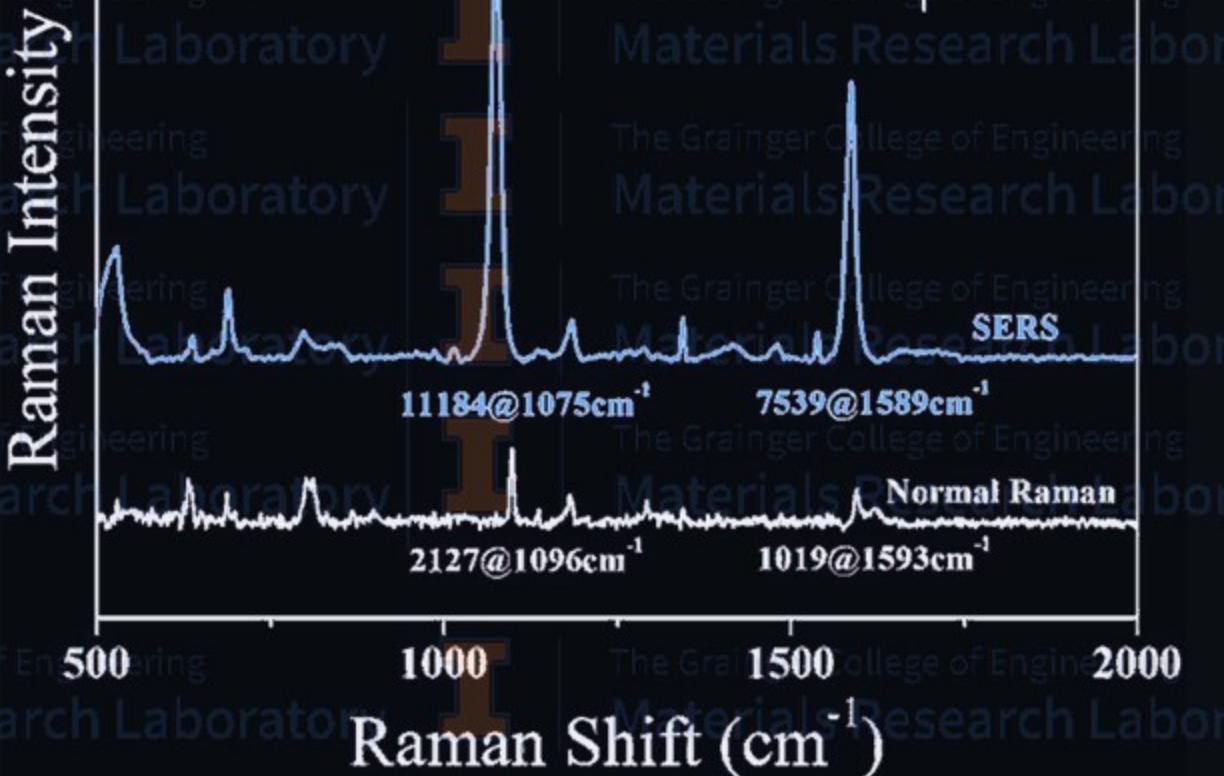


Electrons



Prism

3500 cps



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Anal. Methods, 6, 9547 (2014)

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That's cool, but what about ...

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That's cool, but what about ...

- Limited spatial resolution (diffraction limited).

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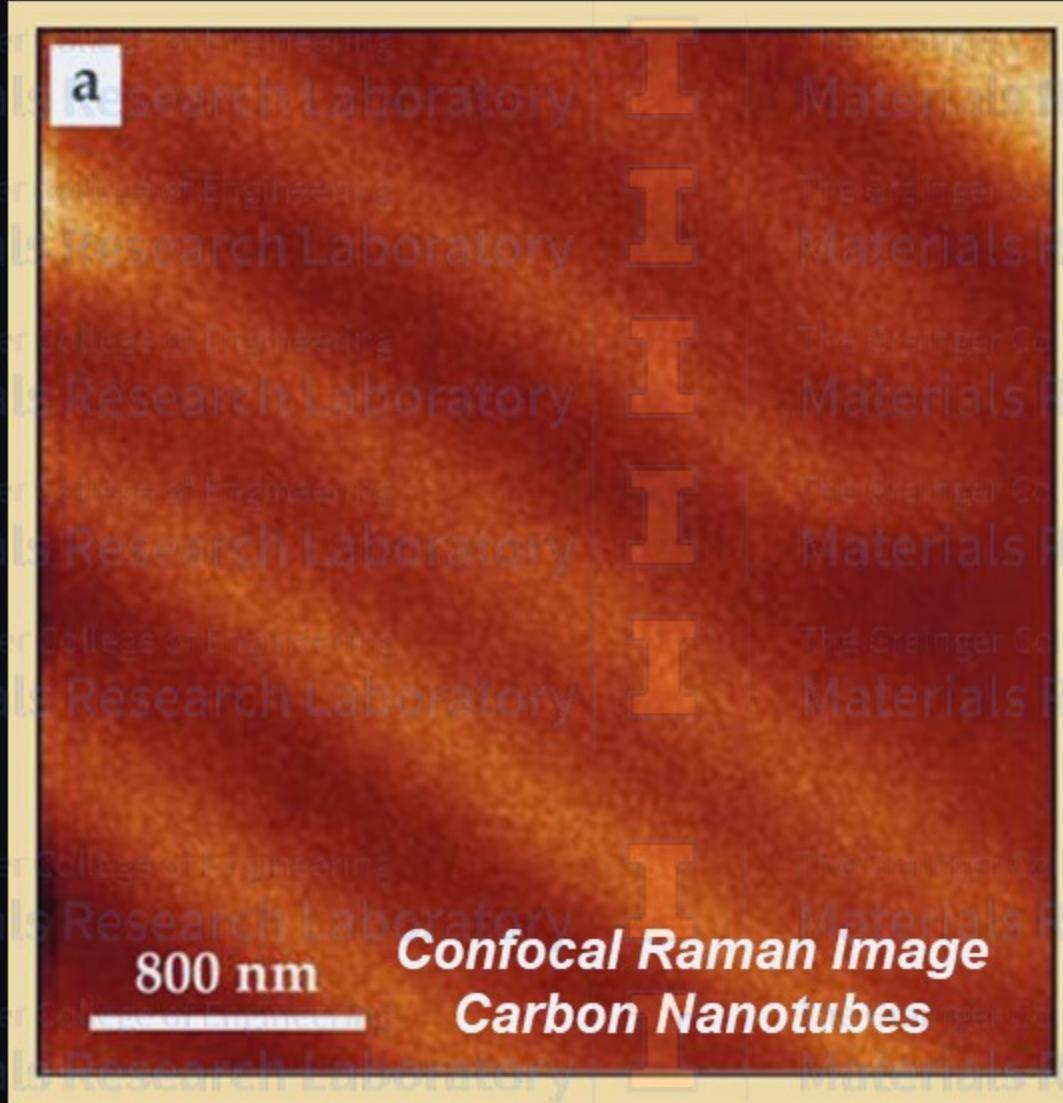
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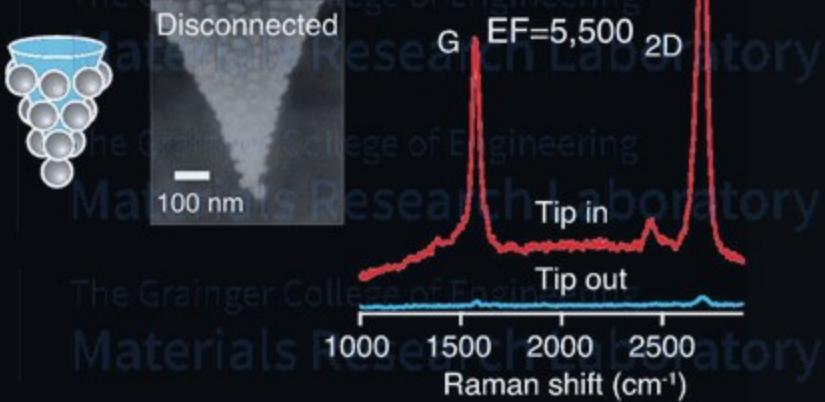


Phys. Rev. Lett. **103**, 186101 (2009)

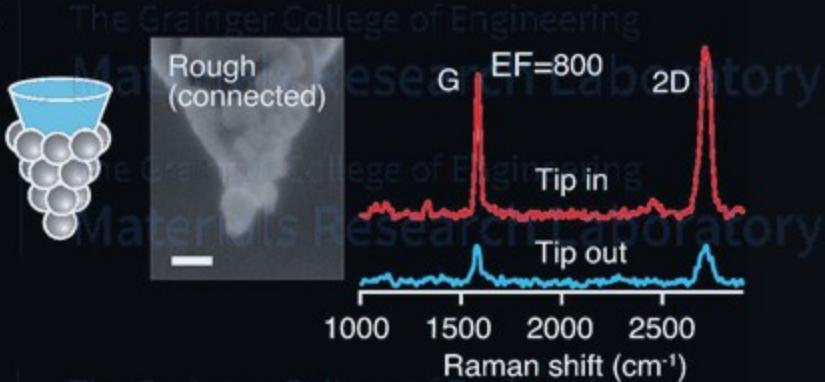
Tip Enhanced Raman Spectroscopy (TERS)

(a)

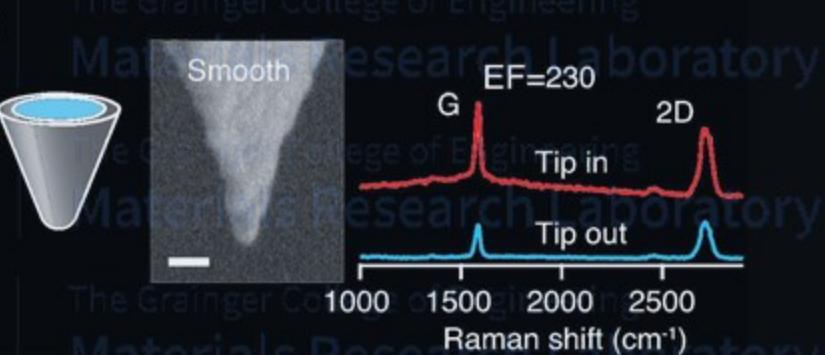
Raman signal TERS enhancement



(b)



(c)



What is really cool is that this also works with a single metalized sharp tip, such as an STM or AFM tip!

Not only do you get the electric field enhancement, but now the source of the Raman signal is extremely localized.

Tip Enhanced Raman Spectroscopy (TERS)

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800 nm

Confocal Raman Image

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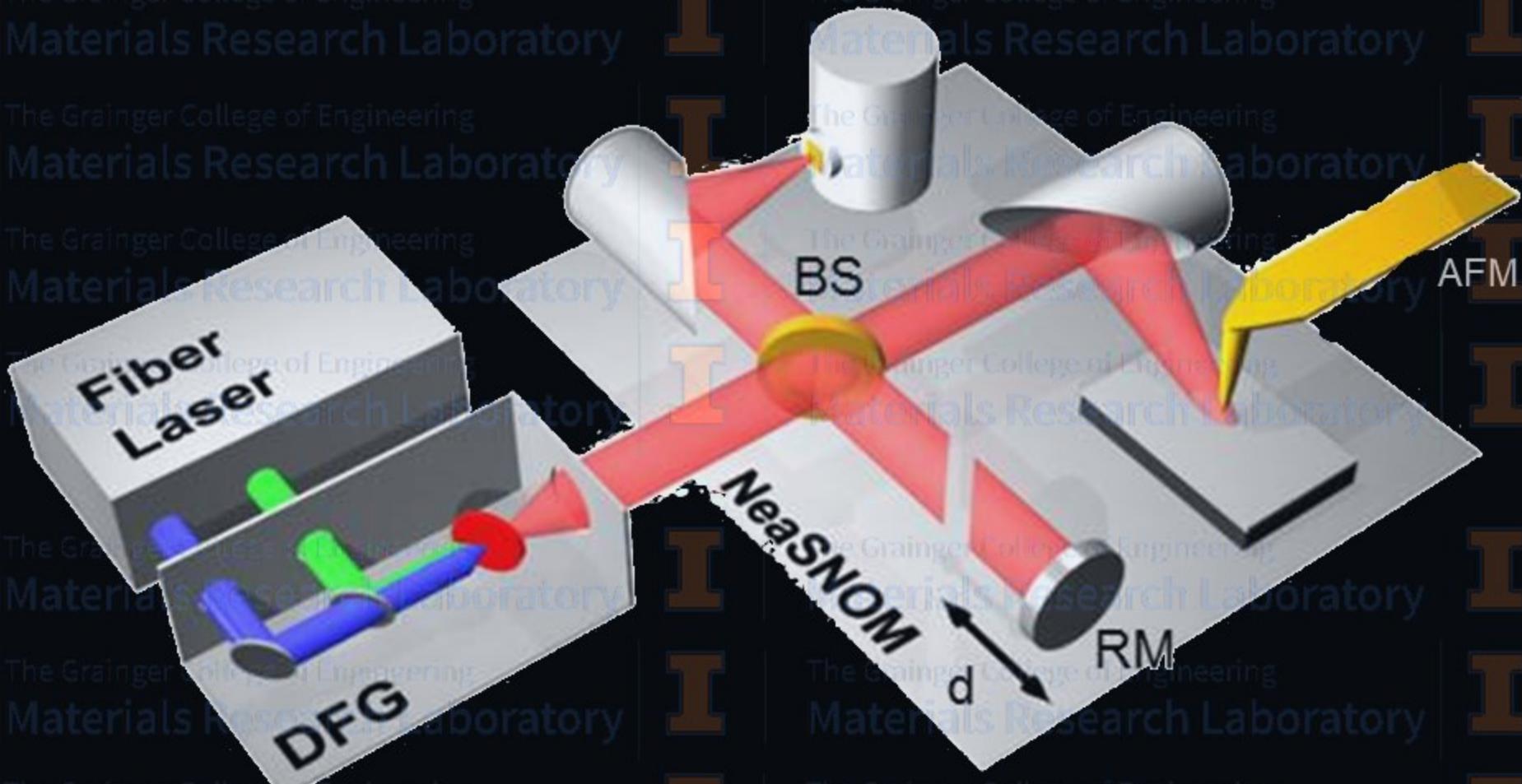
15 nm

Tip Enhanced Raman Image

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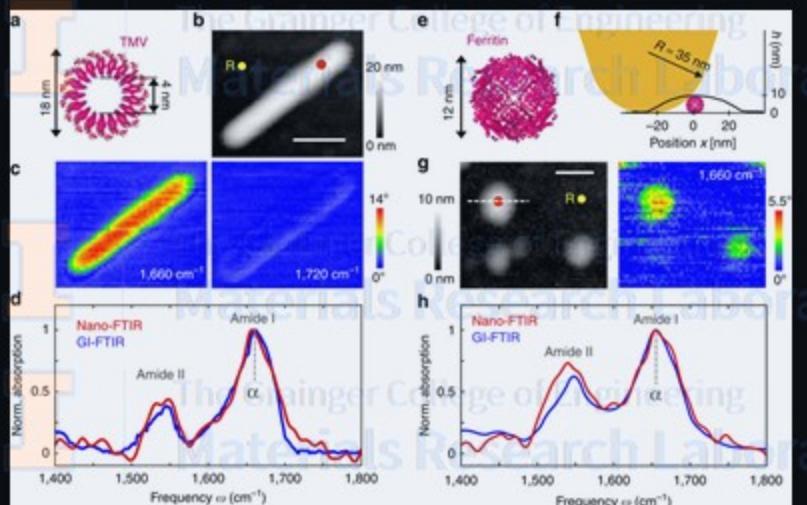
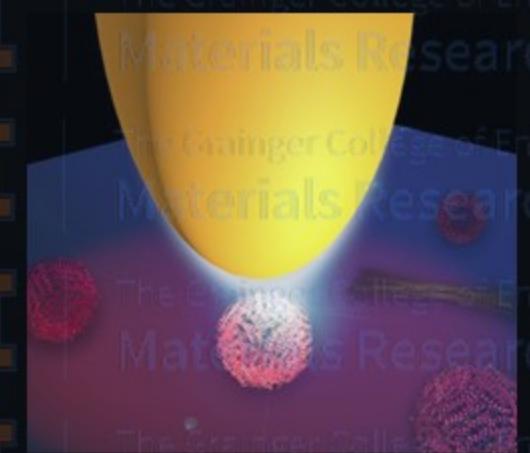
Phys. Rev. Lett. **103**, 186101 (2009)¹⁰⁴

Near-field scanning optical nanospectroscopy

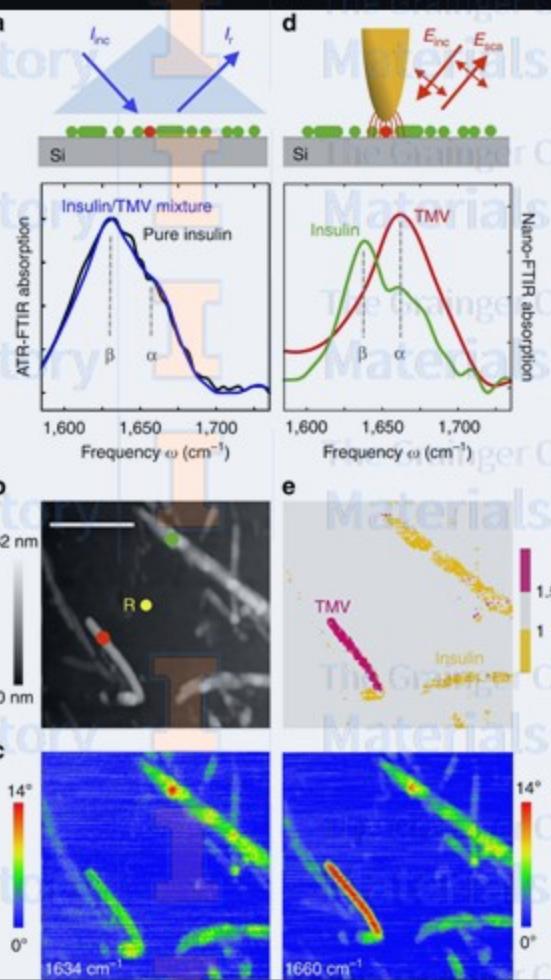


Near-field scanning optical nanospectroscopy

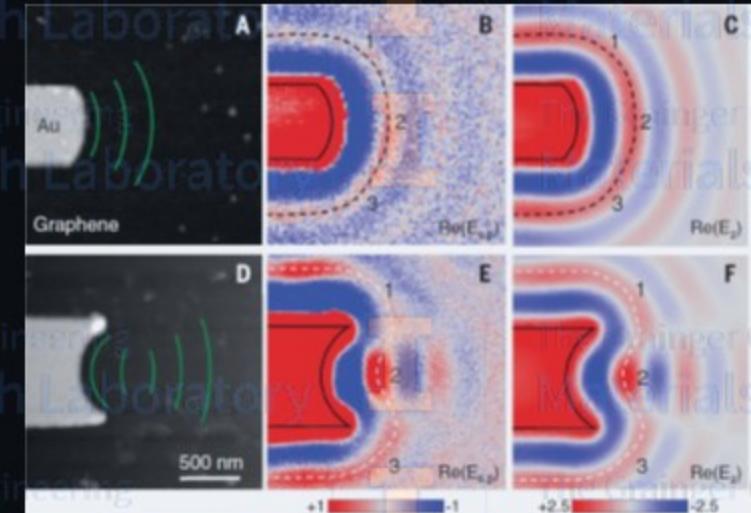
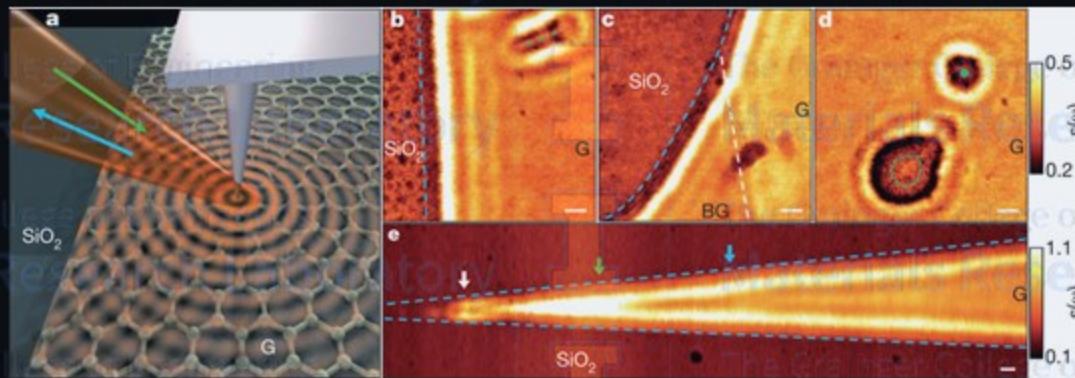
Nano-FTIR



Nature Communications 4, 2890



Nature 000, 1-4 (2012) doi:10.1038/nature11253



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