

**SMB** Consider a system of  $N$  quantum spins, each with with two possible energy levels. One of the levels has zero energy and the other has energy  $\epsilon > 0$ .

Initially work in the *microcanonical ensemble* in which the system as a whole is thermally isolated and has fixed total energy  $E$ , but the individual spins can exchange energy with each other.

- a) Compute  $\Omega(E, N)$ , the number of microscopic states accessible to this system at a given  $E$  and  $N$ , and from it write down an expression for the system's entropy per spin  $S/N$ . In the limit that  $E$  and  $N$  both become large while  $u = E/N$  remains fixed, use Stirling's approximation  $n! = n \ln n - n$  to find  $s$  as a function of  $u$ , and sketch a graph of  $s(u)$ .
- b) Recall that if the system is large enough, we can assign to it a temperature  $T$ . Use this definition of  $T$  in terms of  $S$  and  $E$  to compute  $u(T)$  and hence the specific heat  $C = dE/dT$  as a function of  $\epsilon$  and  $T$ . Sketch a graph of  $u(T)$  and discuss the physics corresponding to the case where  $u > \epsilon/2$ .

Consider now the situation where the system is no longer isolated, but is instead in contact with a reservoir at temperature  $T$ .

- c) Compute the canonical partition function  $Z$  for the system by computing the  $Z_{\text{spin}}$ 's for each individual spin and by suitably combining them. Show that the  $Z$  of the full system can also be obtained using the same  $\Omega(E, N)$  as in part (a).
- d) From  $Z$  write down the Helmholtz free energy per particle as a function of  $\epsilon$  and  $T$  and use it to compute  $u(T)$ . Compare your result with that from part (b). Are they compatible?