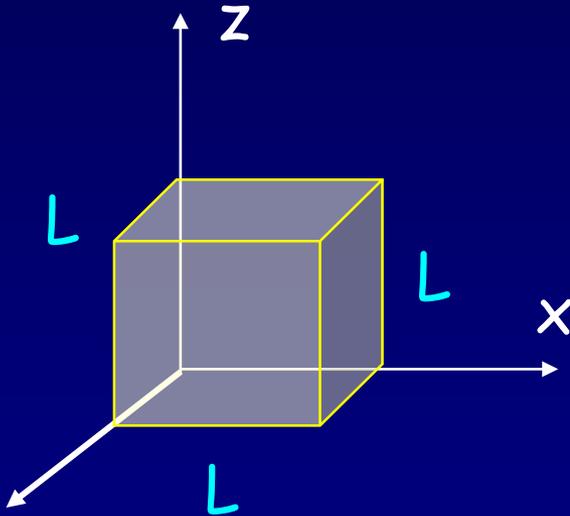


# Lecture 16:

## 3D Potentials and the Hydrogen Atom

3D infinite potential well (cubic)

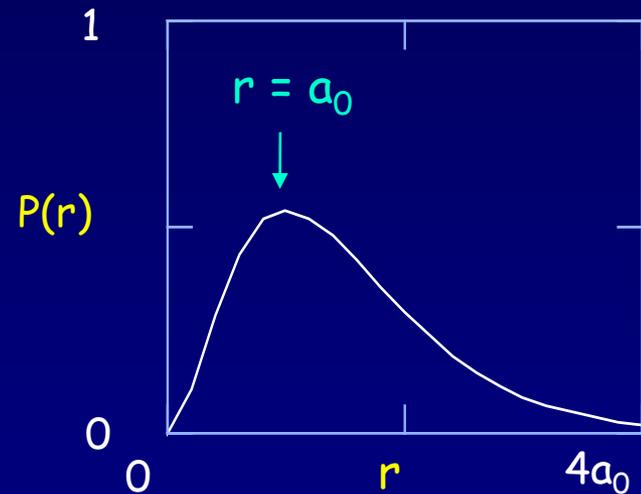
$$\psi(x, y, z) = \varphi(x)\varphi(y)\varphi(z)$$



$$E_{n_x n_y n_z} = \frac{h^2}{8mL^2} \cdot (n_x^2 + n_y^2 + n_z^2)$$

3D Coulomb potential well (hydrogen)

$$\psi(r) = \sqrt{\frac{1}{\pi a_0^3}} e^{-r/a_0}$$



$$E_n = \frac{-13.6 \text{ eV}}{n^2}$$

# Today

## 3-Dimensional Potential Well:

- Separable  $\rightarrow$  Product wave functions
- Degeneracy
- Probability density and normalization

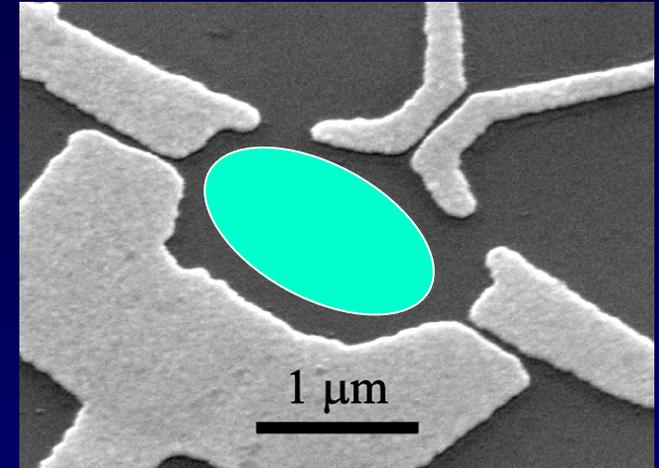
## The Hydrogen Atom:

- Early experiments
- Semi-quantitative picture from uncertainty principle
- Hydrogenic ions
- Next time: Separable  $\rightarrow$  Product wave functions

# Quantum Particles in 3D Potentials

So far, we have considered quantum particles bound in one-dimensional potentials. This situation can be applicable to certain physical systems but it lacks some of the features of most real 3D quantum systems, such as atoms and artificial structures.

A real (2D) “quantum dot”



<http://pages.unibas.ch/phys-meso/Pictures/pictures.html>

One consequence of confining a quantum particle in two or three dimensions is “degeneracy” -- the existence of several quantum states at the same energy.

To illustrate this important point in a simple system, let's extend our favorite potential - the infinite square well - to three dimensions.



# Particle in a 3D Box (1)

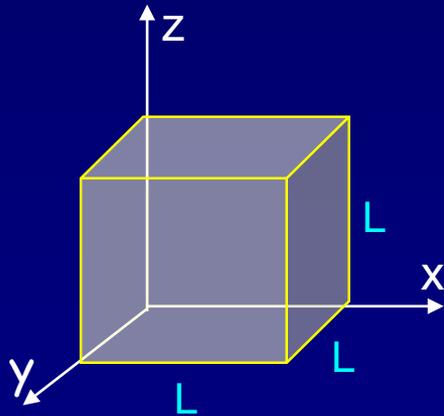
The extension of the Schrödinger Equation (SEQ) to 3D is straightforward in Cartesian (x,y,z) coordinates:

$$-\frac{\hbar^2}{2m} \left( \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} \right) + U(x,y,z)\psi = E\psi$$

where  $\psi \equiv \psi(x,y,z)$

Kinetic energy term:  $\frac{1}{2m}(p_x^2 + p_y^2 + p_z^2)$

Let's solve this SEQ for the particle in a 3D cubical box:



$$U(x,y,z) = \begin{cases} \infty & \text{outside box, } x \text{ or } y \text{ or } z < 0 \\ 0 & \text{inside box} \\ \infty & \text{outside box, } x \text{ or } y \text{ or } z > L \end{cases}$$

This  $U(x,y,z)$  can be “separated”:  
 $U(x,y,z) = U(x) + U(y) + U(z)$

$U = \infty$  if any of the three terms =  $\infty$ .

# Particle in a 3D Box (2)

Whenever  $U(x,y,z)$  can be written as the sum of functions of the individual coordinates, we can write some wave functions as products of functions of the individual coordinates: (see the supplementary slides)

$$\psi(x,y,z) = f(x)g(y)h(z)$$

For the 3D square well, each function is simply the solution to the 1D square well problem:

$$f_{n_x}(x) = N \sin\left(\frac{n_x \pi}{L} x\right) \quad E_{n_x} = \frac{h^2}{2m} \cdot \left(\frac{n_x}{2L}\right)^2$$

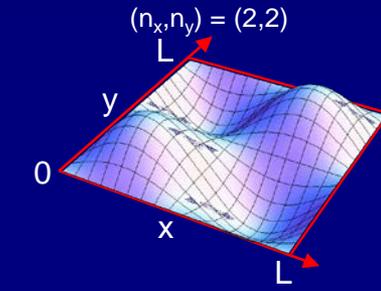
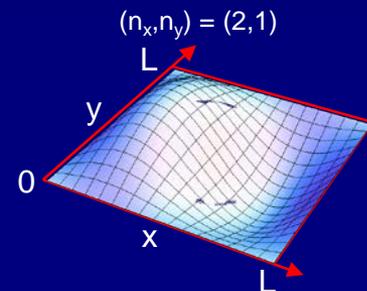
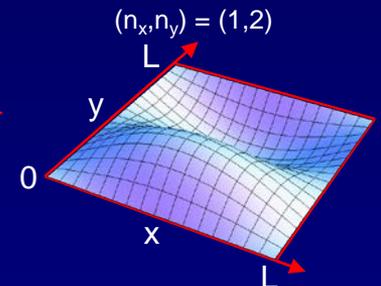
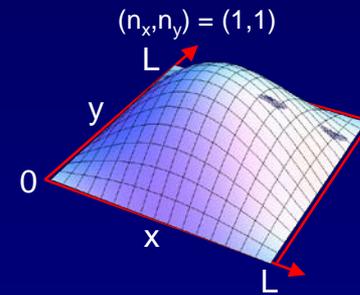
Similarly for  $y$  and  $z$ .

Each function contributes to the energy. The total energy is the sum:

$$E_{\text{total}} = E_x + E_y + E_z$$

2D wave functions:

$$\sin\left(\frac{n_x \pi}{L} x\right) \sin\left(\frac{n_y \pi}{L} y\right)$$



# Supplement: Separation of Variables (1)

In the 3D box, the SEQ is:

$$-\frac{\hbar^2}{2m} \left( \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} + \frac{\partial^2 \psi}{\partial z^2} \right) + (U(x) + U(y) + U(z))\psi = E\psi$$

NOTE:  
Partial derivatives.

Let's see if separation of variables works.  
Substitute this expression for  $\psi$  into the SEQ:

$$\psi(x, y, z) = f(x)g(y)h(z)$$

$$-\frac{\hbar^2}{2m} \left( gh \frac{d^2 f}{dx^2} + fh \frac{d^2 g}{dy^2} + fg \frac{d^2 h}{dz^2} \right) + (U(x) + U(y) + U(z)) fgh = E fgh$$

NOTE:  
Total derivatives.

Divide by fgh:

$$-\frac{\hbar^2}{2m} \left( \frac{1}{f} \frac{d^2 f}{dx^2} + \frac{1}{g} \frac{d^2 g}{dy^2} + \frac{1}{h} \frac{d^2 h}{dz^2} \right) + (U(x) + U(y) + U(z)) = E$$



# Supplement: Separation of Variables (2)

Regroup:

$$\left[ -\frac{\hbar^2}{2m} \frac{1}{f} \frac{d^2 f}{dx^2} + U(x) \right] + \left[ -\frac{\hbar^2}{2m} \frac{1}{g} \frac{d^2 g}{dy^2} + U(y) \right] + \left[ -\frac{\hbar^2}{2m} \frac{1}{h} \frac{d^2 h}{dz^2} + U(z) \right] = E$$

A function of x

A function of y

A function of z

We have three functions, each depending on a different variable, that must sum to a constant.

Therefore, each function must be a constant:

$$-\frac{\hbar^2}{2m} \frac{1}{f} \frac{d^2 f}{dx^2} + U(x) = E_x$$

$$-\frac{\hbar^2}{2m} \frac{1}{g} \frac{d^2 g}{dy^2} + U(y) = E_y$$

$$-\frac{\hbar^2}{2m} \frac{1}{h} \frac{d^2 h}{dz^2} + U(z) = E_z$$

$$E_x + E_y + E_z = E$$

Each function,  $f(x)$ ,  $g(y)$ , and  $h(z)$  satisfies its own 1D SEQ.

# Particle in a 3D Box (2)

Whenever  $U(x,y,z)$  can be written as the sum of functions of the individual coordinates, we can write some wave functions as products of functions of the individual coordinates: (see the supplementary slides)

$$\psi(x,y,z) = f(x)g(y)h(z)$$

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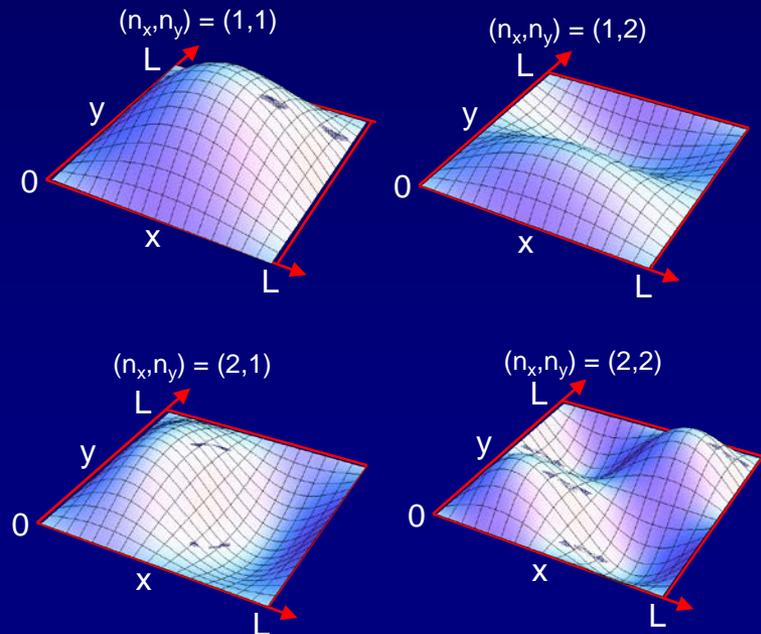
Similarly for  $y$  and  $z$ .

Each function contributes to the energy. The total energy is the sum:

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2D wave functions:

$$\sin\left(\frac{n_x \pi}{L} x\right) \sin\left(\frac{n_y \pi}{L} y\right)$$

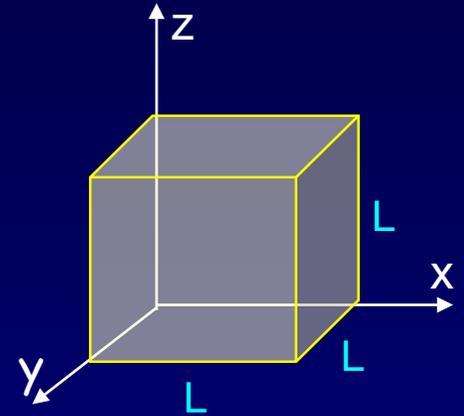


# Particle in a 3D Box (3)

The energy eigenstates and energy values in a 3D cubical box are:

$$\psi = N \sin\left(\frac{n_x \pi}{L} x\right) \sin\left(\frac{n_y \pi}{L} y\right) \sin\left(\frac{n_z \pi}{L} z\right)$$
$$E_{n_x n_y n_z} = \frac{h^2}{8mL^2} (n_x^2 + n_y^2 + n_z^2)$$

where  $n_x, n_y,$  and  $n_z$  can each have values 1,2,3,....



This problem illustrates two important points:

- Three quantum numbers ( $n_x, n_y, n_z$ ) are needed to identify the state of this three-dimensional system. That is true for every 3D system.
- More than one state can have the same energy: “Degeneracy”. Degeneracy reflects an underlying symmetry in the problem. 3 equivalent directions, because it’s a cube, not a rectangle.



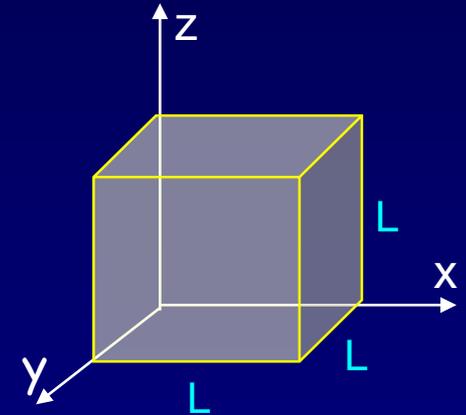
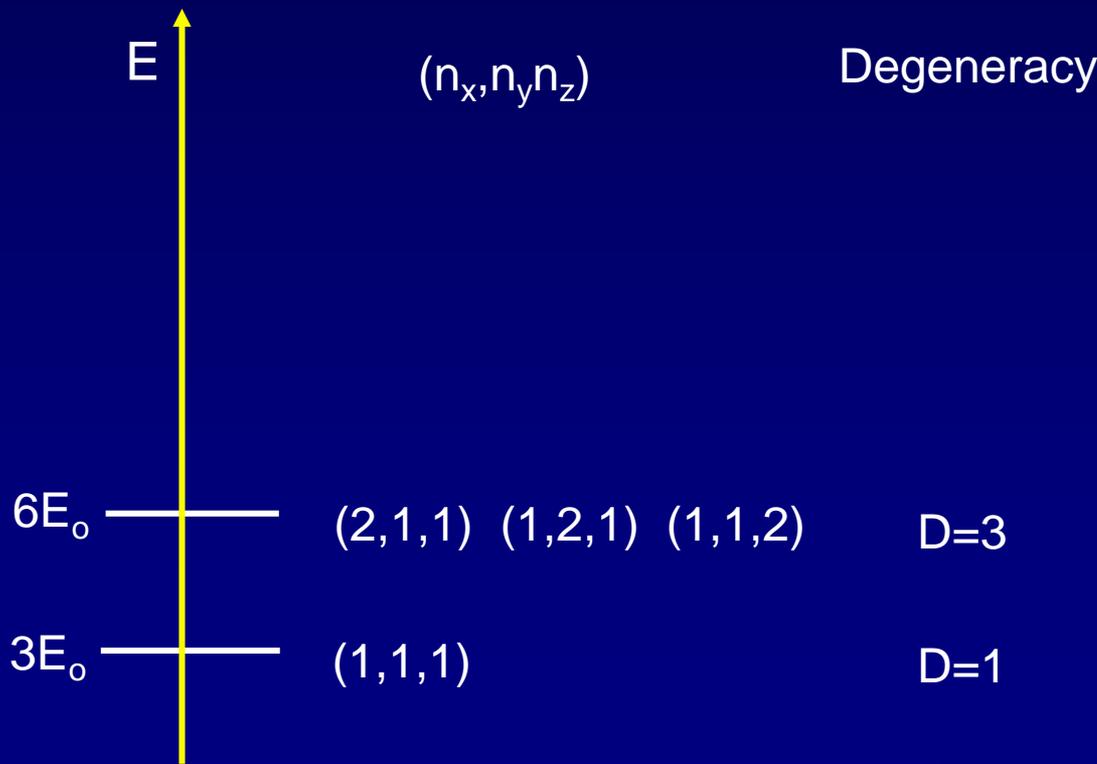
# Cubical Box Exercise

Consider a 3D cubic box:

Show energies and label  $(n_x, n_y, n_z)$  for the first 11 states of the particle in the 3D box, and write the degeneracy,  $D$ , for each allowed energy. Define  $E_0 = h^2/8mL^2$ .

$$\psi = N \sin\left(\frac{n_x \pi}{L} x\right) \sin\left(\frac{n_y \pi}{L} y\right) \sin\left(\frac{n_z \pi}{L} z\right)$$

$$E_{n_x n_y n_z} = \frac{h^2}{8mL^2} (n_x^2 + n_y^2 + n_z^2)$$



# Act 1

For a cubical box, we just saw that the 5<sup>th</sup> energy level is at  $12 E_0$ , with a degeneracy of 1 and quantum numbers (2,2,2).

1. What is the energy of the next energy level?

- a.  $13E_0$       b.  $14E_0$       c.  $15E_0$

2. What is the degeneracy of this energy level?

- a. 2              b. 4              c. 6

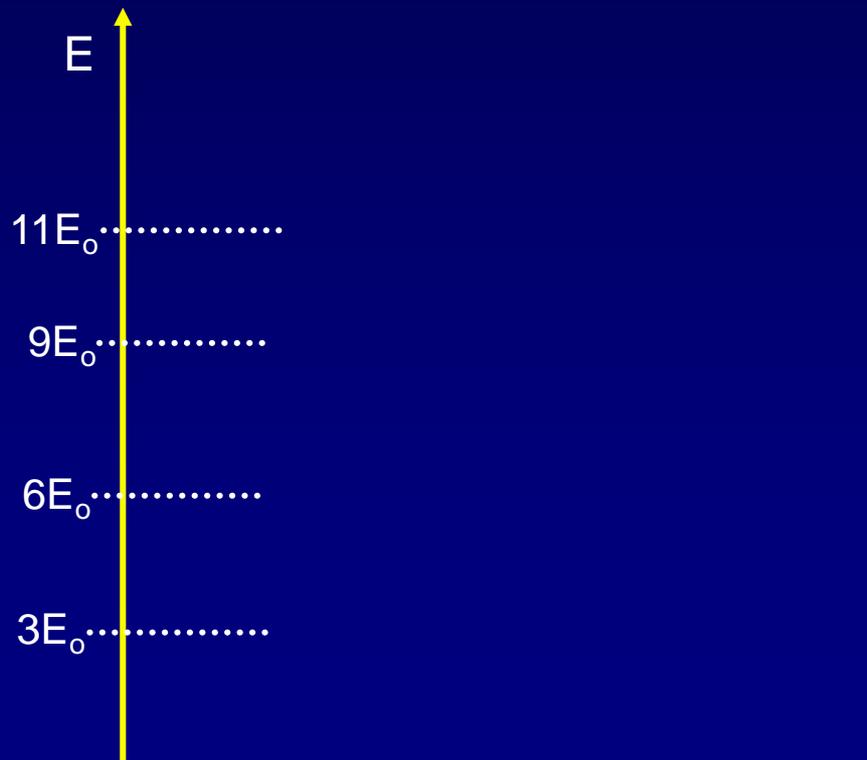
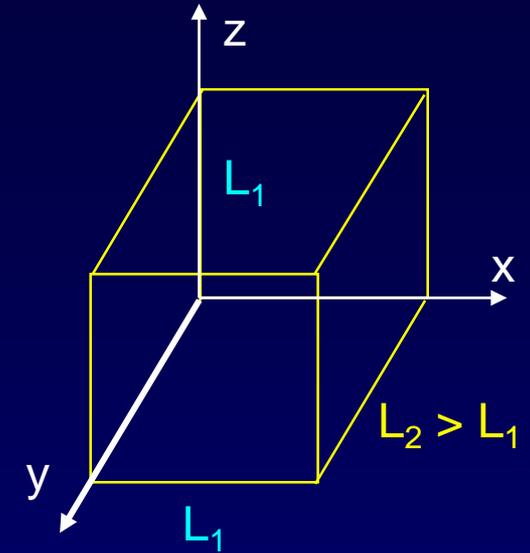
# Non-cubic Box

Consider a non-cubic box:

The box is stretched along the y-direction.

What will happen to the energy levels?

Define  $E_0 = h^2/8mL_1^2$





# Act 2

Consider a particle in a 2D well, with  $L_x = L_y = L$ .

1. Compare the energies of the (2,2), (1,3), and (3,1) states?

a.  $E_{(2,2)} > E_{(1,3)} = E_{(3,1)}$

b.  $E_{(2,2)} = E_{(1,3)} = E_{(3,1)}$

c.  $E_{(2,2)} < E_{(1,3)} = E_{(3,1)}$

2. If we squeeze the box in the x-direction (i.e.,  $L_x < L_y$ ) compare  $E_{(1,3)}$  with  $E_{(3,1)}$ .

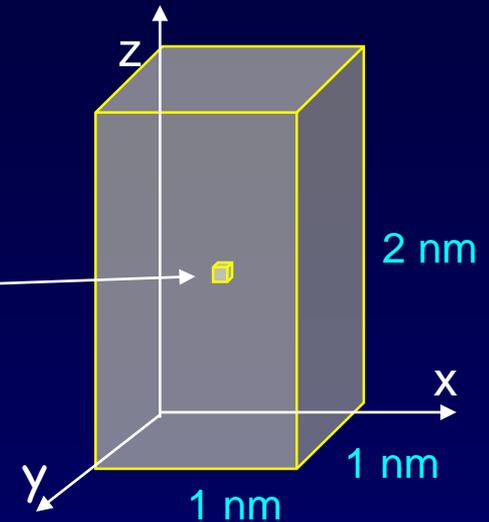
a.  $E_{(1,3)} < E_{(3,1)}$

b.  $E_{(1,3)} = E_{(3,1)}$

c.  $E_{(1,3)} > E_{(3,1)}$

# Probability Exercise

Consider an electron in a 3D rectangular box of size 1 x 1 x 2 nm. Assuming the electron is in the lowest energy state, what is the (approximate) probability to find it at the center of the box, within a region 0.1 x 0.1 x 0.1 nm?



Just as in 1D problems, the probability to find a particle is given by the integral of the probability density ( $\equiv |\psi|^2$ ) over the region of interest.

We properly normalize  $\psi$  by forcing the integral of  $|\psi|^2$  over all space = 1.

For an infinite 3D well, we have

$$\psi = \sqrt{\frac{2}{L_x}} \sqrt{\frac{2}{L_y}} \sqrt{\frac{2}{L_z}} \sin(k_x x) \sin(k_y y) \sin(k_z z)$$

# Another 3D System: The Atom

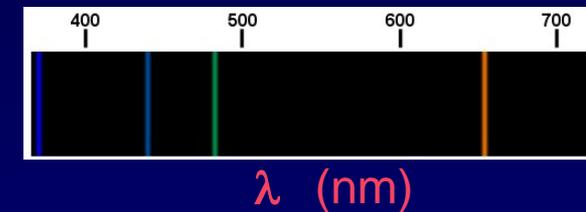
-electrons confined in Coulomb field of a nucleus

Early hints of the quantum nature of atoms:

Discrete Emission and Absorption spectra

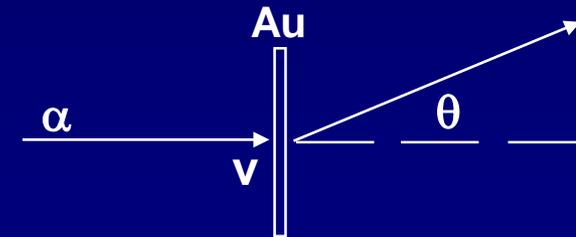
- When excited in an electrical discharge, atoms emit radiation only at discrete wavelengths
- Different emission spectra for different atoms

Atomic hydrogen



Geiger-Marsden (Rutherford) Experiment (1911):

- Measured angular dependence of a particles (He ions) scattered from gold foil.
- Mostly scattering at small angles  $\rightarrow$  supported the "plum pudding" model. But...
- Occasional scatterings at large angles  $\rightarrow$  Something massive in there!



- **Conclusion:** Most of atomic mass is concentrated in a small region of the atom

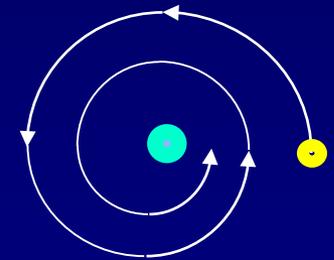
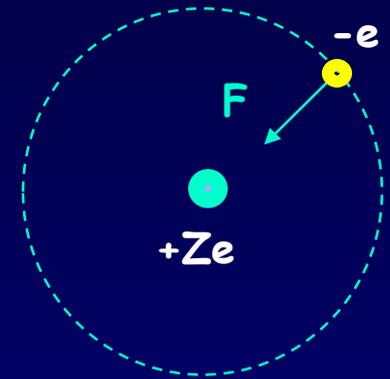
**$\rightarrow$  a nucleus!**



# Atoms: Classical Planetary Model

(An early model of the atom)

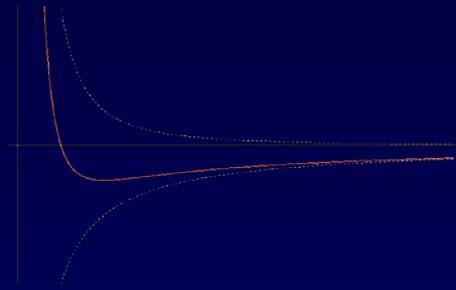
- **Classical picture:** negatively charged objects (electrons) orbit positively charged nucleus due to Coulomb force.
- There is a **BIG PROBLEM** with this:
  - As the electron moves in its circular orbit, it is **ACCELERATING**.
  - As you learned in Physics 212, accelerating charges radiate electromagnetic energy.
  - Consequently, an electron would continuously lose energy and spiral into the nucleus in about  $10^{-9}$  sec.



➔ The planetary model doesn't lead to stable atoms.

# Hydrogen Atom - Qualitative

Why doesn't the electron collapse into the nucleus, where its potential energy is lowest?



We must balance two effects:

- As the electron moves closer to the nucleus, its potential energy decreases (more negative):
- However, as it becomes more and more confined, its kinetic energy increases:

$$U = -\frac{\kappa e^2}{r}$$

$$p \approx \frac{\hbar}{r} \Rightarrow KE \approx \frac{\hbar^2}{2mr^2}$$

Therefore, the total energy is:

$$E = KE + PE \approx \frac{\hbar^2}{2mr^2} - \frac{\kappa e^2}{r}$$

$E$  has a minimum at:

$$r \approx \frac{\hbar^2}{m\kappa e^2} \equiv a_0 = 0.053 \text{ nm}$$

The "Bohr radius" of the H atom.

At this radius,

$$E \approx -\frac{m\kappa^2 e^4}{2\hbar^2} = -13.6 \text{ eV}$$

The ground state energy of the hydrogen atom.

Heisenberg's uncertainty principle prevents the atom's collapse.

One factor of  $e$  or  $e^2$  comes from the proton charge, and one from the electron.

# Act 3

Consider an electron around a nucleus that has two protons, like an ionized Helium atom.

1. Compare the “effective Bohr radius”  $a_{0,\text{He}}$  with the usual Bohr radius for hydrogen,  $a_0$ :

a.  $a_{0,\text{He}} > a_0$

b.  $a_{0,\text{He}} = a_0$

c.  $a_{0,\text{He}} < a_0$

$$r \approx \frac{\hbar^2}{m\kappa e^2} \equiv a_0 = 0.053 \text{ nm}$$

The “Bohr radius”  
of the H atom.

2. What is the ratio of ground state energies  $E_{0,\text{He}}/E_{0,\text{H}}$ ?

a.  $E_{0,\text{He}}/E_{0,\text{H}} = 1$

b.  $E_{0,\text{He}}/E_{0,\text{H}} = 2$

c.  $E_{0,\text{He}}/E_{0,\text{H}} = 4$

# Next Lectures

Angular momentum → atomic orbitals

“Spin” → Pauli Exclusion Principle