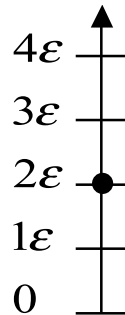


Thermodynamic processes almost always involve the exchange of energy between subsystems. To begin our study of energy exchange we consider the simplest possible system: a **harmonic oscillator** such as a vibrating diatomic molecule. The molecule can absorb or emit only an integral number of energy packets. We write the energy levels of an oscillator as a ladder of equally spaced energy levels, where epsilon ( $\epsilon$ ) is the quantum of energy:

$$\text{Energy} = n\epsilon$$

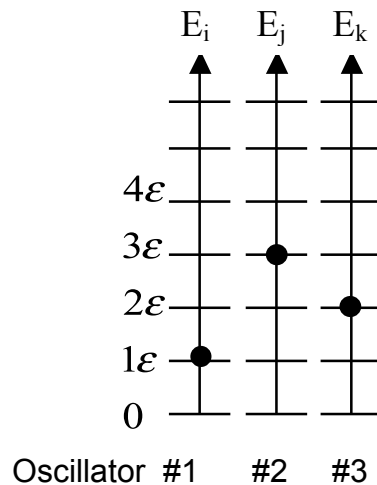


Convention:

$E$  = single-particle energy

$U$  = energy of a many-particle system

The oscillator shown has an energy of  $2\epsilon$ ; that is, it possesses two quanta (or packets) of energy. Now imagine three such oscillators which are able to exchange energy. One possible **microstate** of the 3-particle system is the following:

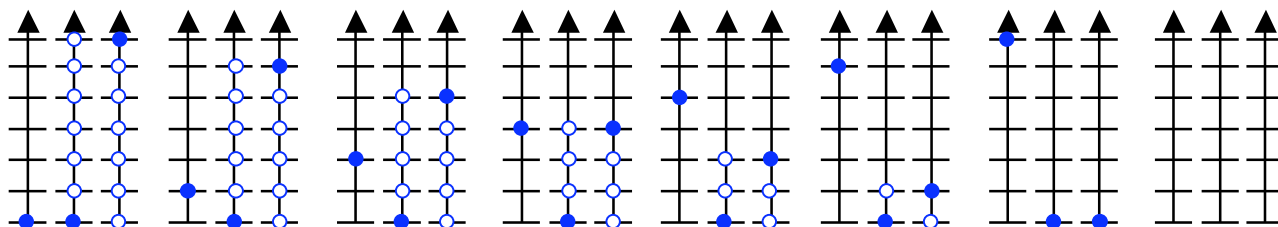


- a) What is the total energy  $U$  of the situation depicted above, in terms of  $\epsilon$ ?  $U = \underline{6\epsilon}$ . In general we will write  $U = q\epsilon$ , where  $q$  is the total number of energy quanta in the system. From the convention above,  $U$  represents the total energy of all of the oscillators. So all we need to do is add the number of quanta that are in each of the oscillators. Thus  $U = 6\epsilon$ .

- b) Determine the number of microstates for the three-oscillator system with a total energy of  $U = 6\epsilon$ .

[Hint: Start with  $E_i = 0$  and count the # of ways you could arrange  $6\epsilon$  in oscillators 2 and 3.

Next set  $E_i = \epsilon$  and count the # of ways you could arrange  $5\epsilon$  in oscillators 2 and 3. Continue, and add up all the possibilities.  $\Omega = \text{total \# microstates}$ . Here is some workspace:



Each diagram shows all the combinations for one choice of  $E_i$ . For example, when  $E_i = 0$ ,  $E_j$  can take any value between 0 and 6 (i.e., 7 values). For each value of  $E_j$ ,  $E_k$  is constrained by energy conservation. So, there are a total of  $7+6+5+4+3+2+1 = 28$  total microstates for  $U = 6\epsilon$ .

- c) Now fill out the table below by counting the number of microstates  $\Omega_q$  for the three-oscillator system for the total energies shown. To save time, notice how the  $6\epsilon$  answer is changed when you go to  $5\epsilon$ , etc. Check your answer with the formula given below, a derivation of which will be given in the posted solutions.

$U = q\epsilon$	$\Omega_q$
0	1
$1\epsilon$	3
$2\epsilon$	6
$3\epsilon$	10
$4\epsilon$	15
$5\epsilon$	21
$6\epsilon$	28

$$\Omega_q = \frac{(q + N - 1)!}{(N - 1)! q!}$$

$q = \# \text{ quanta}$        $N = \# \text{ oscillators}$

The calculation for each  $q$  is similar to the  $q = 6$  case.

This equation can be derived from the binomial coefficient,  $C_b(N,n)$ .

We want to distribute  $q$  quanta into  $N$  bins. Putting the  $N$  bins in a row, we see that there are  $N+1$  walls. So, a typical microstate will look like this: (for  $q = 7$  and  $N = 9$ )

| \* || \* \* | \* ||| \* \* | \* |

where we've denoted energy quanta by \* and walls by |.

Two adjacent \* indicates multiple occupancy. Two adjacent | indicates an empty bin. In the diagram, the occupancies of the 9 bins are: (1,0,2,1,0,0,2,1)

Note that the diagram must begin and end with |, because all energy must be in bins. We want to count the number of permutations of the  $N-1$  remaining walls and the  $q$  quanta (i.e.,  $q+N-1$  objects), remembering that the walls are indistinguishable, as are the quanta. Thus, the answer is:

$$\Omega_q = \frac{(q + N - 1)!}{(N - 1)! q!}$$

